

COSMIC PARALLAX

A NEW TOOL IN OBSERVATIONAL COSMOLOGY

Luca Amendola¹, Miguel Quartin^{1,2}, Claudia Quercellini³

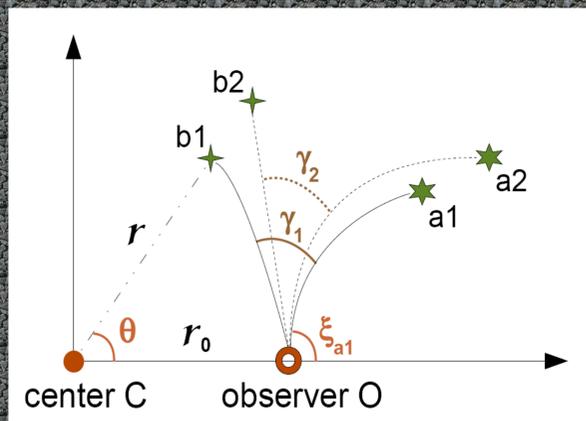
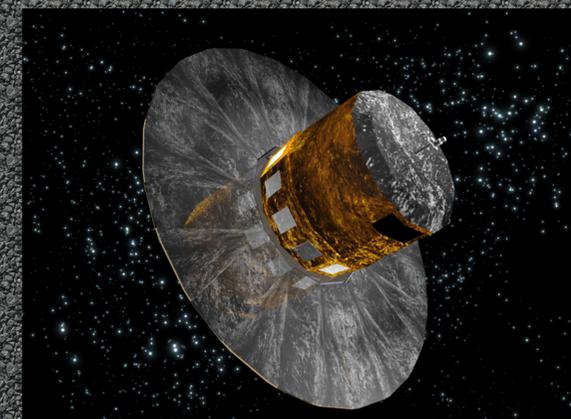


Figure 1: Overview, notation and conventions. Note that for clarity purposes we assumed here that the points C; O; a1; b1 all lie on the same plane. By symmetry, points a2; b2 remain on this plane as well. Comoving coordinates r and r_0 correspond to physical coordinates X and X_0 .

Future satellite missions such as GAIA will achieve astrometry measurements with an accuracy of about $10 \mu\text{as}$ (microarcseconds) for bright sources; other satellite proposals aim at $1 \mu\text{as}$. We show in this paper that such refined measurements allow us to detect large-scale deviations from isotropy through real-time observations of changes in the angular separation between sources at cosmic distances. We also show that this cosmic parallax effect is a powerful consistency test of Friedmann-Robertson-Walker metric and may set very strong constraints on alternative anisotropic models like Lemaître-Tolman-Bondi cosmologies with off-center observers.

We will refer to (t, r, θ, ϕ) as the comoving coordinates with origin on the center of a spherically symmetric model (see Figure 1). Peculiar velocities apart, the symmetry of such a model forces objects to expand radially outwards, keeping r, θ and ϕ constant.

In a FRW metric, $\Delta\gamma \equiv \gamma_2 - \gamma_1 = 0$ (see Figure 1). In any anisotropic metric, however, $\Delta\gamma \neq 0$, and we have **cosmic parallax**. In such cases, FRW-like estimates can give a reasonable prediction of the magnitude of this effect.



Artist conception of the ESA GAIA mission satellite. Expected launch date: Dec 2011. GAIA expects to perform astrometric measurements of some **500,000** distant quasars.

FRW-Like estimates for pair of sources at (i) the same redshift and (ii) same position in the sky,

$$X \equiv a(r, t_0) r$$

$$s \equiv X_0 / X$$

$$\Delta\gamma_\theta = 2 s \cos\theta (H_0 - H_r) \Delta\theta \Delta t$$

$$\Delta\gamma_z = s \sin\theta \left[\frac{dH_r}{dX} \frac{1}{H(r, z)} + \frac{H_0 - H_r}{X H(r, z)} \right] \Delta z \Delta t$$

Peculiar velocity noise is **not** overwhelming.

$$\Delta\gamma_{\text{pec}} = \left(\frac{v_{\text{pec}}}{1000 \frac{\text{km}}{\text{s}}} \right) \left(\frac{D_A}{1 \text{ Gpc}} \right)^{-1} \left(\frac{\Delta t}{10 \text{ years}} \right) 10^{-11} \text{ rad}$$

The general LTB metric

$$ds^2 = -dt^2 + \frac{[R'(t, r)]^2}{1 + \beta(r)} dr^2 + R^2(t, r) d\Omega^2$$

LTB metric for a matter dominated universe

$$R = (\cosh \eta - 1) \frac{\alpha}{2\beta} + R_{\text{ISS}} \left[\cosh \eta + \sqrt{\frac{\alpha + \beta R_{\text{ISS}}}{\beta R_{\text{ISS}}}} \sinh \eta \right]$$

$$\sqrt{\beta} t = (\sinh \eta - \eta) \frac{\alpha}{2\beta} + R_{\text{ISS}} \left[\sinh \eta + \sqrt{\frac{\alpha + \beta R_{\text{ISS}}}{\beta R_{\text{ISS}}}} (\cosh \eta - 1) \right]$$

The Alnes et al. [1] class of LTB models. We considered 2 such models, basically differing by the value of Δr .

$$\alpha(r) = (H_0^{\text{out}})^2 r^3 \left[1 - \frac{\Delta\alpha}{2} \left(1 - \tanh \frac{r - r_{\text{vo}}}{2\Delta r} \right) \right]$$

$$\beta(r) = (H_0^{\text{out}})^2 r^2 \frac{\Delta\alpha}{2} \left(1 - \tanh \frac{r - r_{\text{vo}}}{2\Delta r} \right)$$

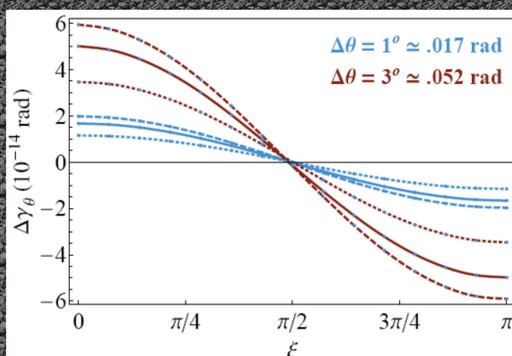


Figure 2: $\Delta\gamma$ for three sources at the same shell, at $z = 3$, for both models I (full lines) and II (dashed), and the FRW-like estimate (dotted). The dark, red lines correspond to a separation $\Delta\theta$ of 3° , while the light, blue lines represent a separation of 1° . As expected, $\Delta\gamma$ is linear in $\Delta\theta$.

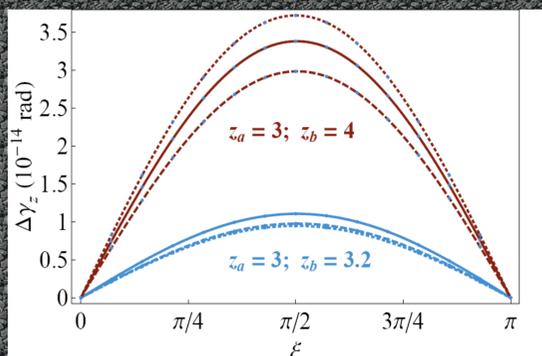


Figure 3: Same as Figure 2 but for $\Delta\theta = 0$ and redshift pairs $\{3, 3.2\}$ (lighter, blue lines) and $\{3, 4\}$ (darker, red lines).

Plots for an off-center distance of 15 Mpc

The cosmic parallax effect can be combined with measurements of the time-drift of the redshift of the sources [3], to fully reconstruct the 3D cosmic flow of distant sources [4].

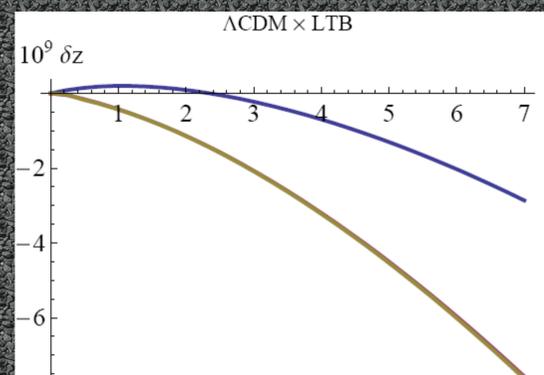


Figure 4: Time-drift of redshift in ΛCDM (upper, blue curve) and in both LTB models (lower, superimposed curves), for a time interval of 10 years. The drift is **much more prominent in an LTB model** [4]. The angle dependence of δz in ξ is only marginal, so the effect is almost isotropic. Note the **insensitivity to the specifics of the LTB model**.

Conclusions:

- * Parallax effects, i.e. the apparent change in position of an object relative to the observer, are among the simplest and historically more important method to measure astronomical distances;
- * The cosmic parallax of quasars and other distance sources in a LTB model is **withing observable reach** of some planned near-future mission such as GAIA;
- * Cosmic parallax is a tool to measure cosmic anisotropy, and as such can distinguish FRW from other metrics as well, such as some Bianchi ones.
- * In an off-center LTB model, cosmic parallax can be an **1+ order of magnitude better probe** of off-center distance than the current best: the CMB dipole [1];
- * Cosmic parallax **avoids 2 intrinsic limitations** of the CMB dipole effect:
 1. It cannot be completely countered by the observer's own peculiar velocity;
 2. Limitation by cosmic variance.
- * Combined with a measurement of the time-drift of redshift [3], cosmic parallax can in principle fully reconstruct the 3D cosmic flow of distant sources.

References:

- [1] Havard Alnes & Morad Amarzguoui, Phys. Rev. D74 (2006) 103520.
- [2] Claudia Quercellini, Miguel Quartin & Luca Amendola, arXiv: 0809.3675v1
- [3] Jean-Philippe Uzan, Chris Clarkson & George F.R. Ellis, Phys. Rev. Lett. **100** 191303 (2008).
- [4] Miguel Quartin and Luca Amendola, in preparation.