

Massive Stars :

their

Birth Sites

and

Distribution

Pavel Kroupa

*Institut für Theoretische Physik und Astrophysik
(Universität Kiel)*

Massive Stars:
their *Birth Sites* and
Distribution

Outline

- Mass distribution (the IMF)
- Spatial distribution (birth sites and outliers)
- Cluster dynamics (mass segregation)
- HII regions (gas expulsion)

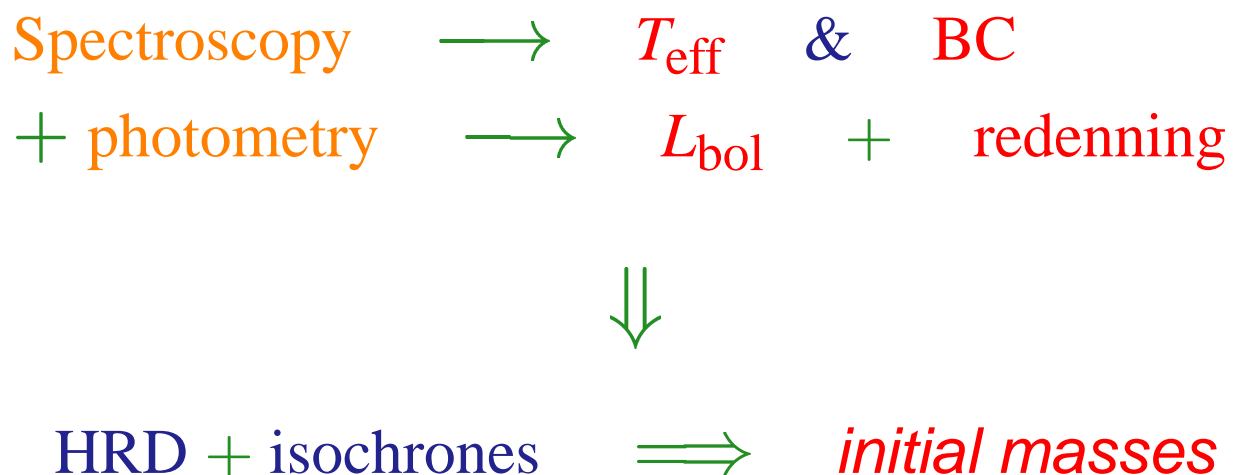
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The Initial Stellar Mass Function

Three steps:

1. Infer *present-day stellar masses* (e.g. in an OB association);
2. Determine ages and SFH;
3. Calculate *initial masses* \longrightarrow IMF

Philip Massey (various papers):



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Def. $lm \equiv \log_{10} m$

$$\begin{aligned} dN &= \xi(m) dm && = \text{number of stars with} \\ & && \text{mass} \in [m, m + dm] \\ &= \xi_L(lm) dlm && = \text{number of stars with} \\ & && \text{log-mass} \in [lm, lm + dlm] \end{aligned}$$

$$\implies \xi_L(m) = (m \ln 10) \xi(m)$$

$$\text{slope : } \Gamma(m) \equiv \frac{d}{dlm} \log_{10} [\xi_L(lm)]$$

Example: power-law form

$$\begin{aligned} \xi_L &= A m^\Gamma = A m^{-x} \\ \xi &= A' m^\alpha = A' m^{-\gamma} \quad (A' = \frac{A}{\ln 10}) \end{aligned}$$

$$\Gamma = -x = 1 + \gamma = 1 - \alpha$$

$$\alpha(m) = 1 - \Gamma(m)$$

Construct **alpha plot** (α vs m) to search for variations in IMF shape.

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Massey et al. (1995) find for OB associations:
(IMF only determined for stars with $\tau_{\text{ms}} \geq \tau = \text{age of pop.}$)

1)	SMC	LMC	MW
	$Z = 0.002$	0.008	0.02
	$\alpha = 2.3 \pm 0.1$	2.3 ± 0.1	2.1 ± 0.1
	$\xi(m)$	$= \xi(m)$	$= \xi(m)$

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2) Maximum stellar mass $\approx 100 - 120 M_{\odot}$
in all cases.

3) Independence of density:

Example: R136

central density $\approx 10^5$ stars/pc³;

39 O3 stars (or more) !

$$\alpha = -2.35 \pm 0.15$$

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Multiplicity

Most massive stars are in

binary (\mathcal{B}),

triple (\mathcal{T}) or

quadruple (\mathcal{Q}) systems.

Def: The Companion Star Fraction:

$$CSF = \frac{\mathcal{B} + 2\mathcal{T} + 3\mathcal{Q}}{\mathcal{S} + \mathcal{B} + \mathcal{T} + \mathcal{Q}}$$

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Example: Orion Nebula Cluster (≈ 1 Myr old)

The multiplicity of the 8 most massive stars
(Preibisch et al. '99):

θ^1 Ori A	\mathcal{T}	
θ^1 Ori B	\mathcal{Q}	
θ^1 Ori C	\mathcal{B}	– <u>the</u> exciting star
θ^1 Ori D	\mathcal{S}	
θ^2 Ori A	\mathcal{T}	
θ^2 Ori B	\mathcal{S}	
LP Ori	\mathcal{T}	
ν Ori	\mathcal{T}	

$$\implies CSF = \frac{1 + 8 + 3}{2 + 1 + 4 + 1} = 1.5$$

separations: $\lesssim 1$ AU – few 1000 AU

mass ratios: $0.1 \lesssim q \leq 1$

Compare: low-mass field stars $CSF \approx 0.5$
low-mass pm stars $CSF \approx 1$

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Surveys: (review: Zinnecker 2003)



- massive stars have > 1 companions on average
- $CSF \uparrow$ as $\rho_{cl} \uparrow$



- Probable *main origin* of massive stars:
accretion-induced protostellar collisions in dense clusters
- High binarity \implies *revision of IMF!*
(Sagar & Richtler 1991)

$$\alpha_{app} = 2.3:$$

(Salpeter)

$$CSF \geq 1 \implies \alpha_{true} \geq 2.7$$

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The Spatial Distribution

Some statistics from the literature:

- 50 % of known O stars within 3 kpc of the Sun are in associations, the others not (Garmany et al.'82).
- 10–25 % of O stars are runaways ($v \gtrsim 40$ km/s)
 ≈ 2 % of B stars are runaways (Gies & Bolton'86).

Massey also notes *isolated populations* of OB stars in SMC, LMC & MW with $\alpha \approx 4.5$!

Is this a different mode of star formation?

– a small cloud gives birth to “a single O star plus a little change” (J. Gallagher).

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A Rare Mode of Star Formation ?

The balance of heating and cooling in a **molecular cloud** can be approximately described by a **polytropic equation of state**:

$$P = k \rho^\gamma$$

N.B. ideal gas: $P = \mathcal{R} T \rho$

isothermal ideal gas: $T = \text{constant}, \gamma = 1$

Galactic molecular clouds: $0.2 < \gamma < 1.4$
(Spaans & Silk 2000)

A region in the cloud collapses under self gravity if its mass surpasses the **Jeans Mass**:

$$M_J = \left(\frac{k' \pi}{G} \right)^{\frac{3}{2}} \gamma^{\frac{3}{2}} \rho^{\frac{3}{2}} \left(\gamma - \frac{4}{3} \right)$$

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(Li, Klessen & Mac Low 2003):

$$\frac{\partial M_J}{\partial \rho} \propto \left(\frac{3\gamma - 4}{2} \right) \rho^{\frac{3\gamma - 6}{2}}$$

< 0 for $\gamma < \frac{4}{3}$
 $M_J \downarrow$ as $\rho \uparrow$ during collapse

$\Rightarrow \frac{\partial M_J}{\partial \rho} = 0$ for $\gamma = \frac{4}{3}$
 $M_J = \text{const}$ as $\rho \uparrow$

> 0 for $\gamma > \frac{4}{3}$
 $M_J \uparrow$ as $\rho \uparrow$: **choke-off**, no collapse !



- **star clusters** form when $\gamma < \frac{4}{3}$;
- **single massive stars** form when $1 < \gamma < 1.4$;
such γ 's are perhaps possible for metal-free gas (inefficient cooling).

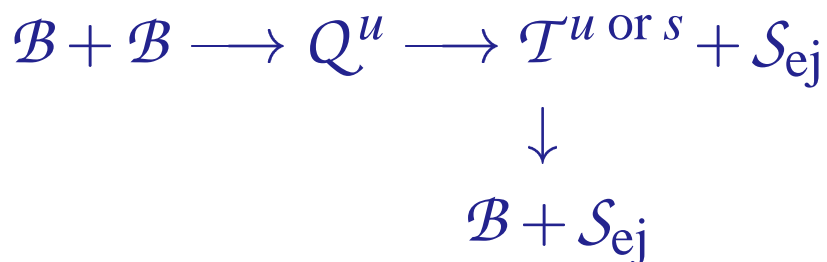
Li et al. **verify this** numerically using SPH calculations.

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Isolated Massive Stars

Two other possibilities:

1. **Supernova explosion** in **massive close binary**:
 $v_{\text{ej}} \approx v_{\text{orb}} \approx 300 \text{ km/s}$ (Tauris & Takens'98).
2. **Dynamical ejection** from **short-lived multiple system** formed through encounters in **cluster core**:



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Leonard 1991: Many $\mathcal{B} + \mathcal{B}$ experiments

$$\begin{aligned} \implies v_{ej} &\leq 1 v_{esc} && \text{for low-mass star} \\ v_{ej} &\leq 0.5 v_{esc} && \text{for massive star} \end{aligned}$$

NB: The asymptotic $\max v_{ej}$ are reached after ∞ number of $\mathcal{B} + \mathcal{B}$ encounters.

Example:

$v_{esc} = 1400 \text{ km/s}$ from surface of a $60 M_{\odot}$ star

\implies maximum $v_{ej} = 700 \text{ km/s}$
for another $60 M_{\odot}$ star

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Example:

PG 0832 + 676: $m \approx 13 M_{\odot}$

18 kpc above Galactic plain

→ need $v_{ej} \gtrsim 1150$ km/s.

This is possible after $10^4 - 10^5$ $\mathcal{B} + \mathcal{B}$ encounters involving at least one $60 M_{\odot}$ star.

However, this many $60 M_{\odot}$ stars are **not available** in Galactic disk !

This many encounters could also not have occurred in $\approx 10^7$ yr.

⇒ dynamical ejection **very unlikely !**

But:

$\mathcal{B} + \text{bh} \longrightarrow (\mathcal{S} + \text{bh}) + \mathcal{S}_{ej}$

(v_{ej} up to **4000 km/s**) (Hills'88).

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Dynamics of Massive Stars in Clusters

Massive stars loose orbital angular momentum through dynamical friction.



$$t_{\text{msegr}} \approx 2 \left(\frac{m_{\text{av}}}{m_1} \right) t_{\text{relax}}$$

$$t_{\text{relax}} = \frac{21}{\ln(0.4N)} \left(\frac{M_{\text{cl}}}{100M_{\odot}} \right)^{\frac{1}{2}} \left(\frac{1M_{\odot}}{m_{\text{av}}} \right) \left(\frac{R_{0.5}}{1\text{pc}} \right)^{\frac{3}{2}}$$

Example:

$t_{\text{relax}} \approx 0.6$ Myr for pre-exposed ONC

$\implies t_{\text{msegr}} \approx 0.012$ Myr \ll age of ONC

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Self-consistent N -body models:

Bonnell & Davies'98: **No** for ONC

Kroupa'02: **Maybe** for ONC

→ Issue **unresolved** but very important for **formation scenarios** of massive stars:

Two possibilities:

- Always in cores of clusters; **coagulation** of intermediate-mass proto-stars plus **accretion of low-angular momentum gas onto core** → **core shrinkage** & more coagulation ... (Bonnell et al.'98).
- Like less-massive stars through **rapid accretion from massive circumstellar disk**; **e.g.:** NGC 7538 – a high-mass protostar with a massive rotating disk (Sandell, Wright & Foster'03).

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Dynamically Unstable Cores

- Mass-segregation very rapid, $t_{\text{msegr}} \approx t_{\text{cross}}$
- ⇒ quick formation of core consisting of the N_m most massive stars.

The core decays within a time

$$t_{\text{decay}} \approx N_m \times t_{\text{cross}}^{\text{core}}$$
$$t_{\text{cross}}^{\text{core}} \approx 5 \left(\frac{M^{\text{core}}}{100 M_{\odot}} \right)^{-\frac{1}{2}} \left(\frac{R_{0.5}^{\text{core}}}{1 \text{ pc}} \right)^{\frac{3}{2}}$$

Example: ONC:

$$R_{0.5}^{\text{core}} \approx 0.02 \text{ pc}, \quad M^{\text{core}} \approx 150 M_{\odot} \quad (N_m \approx 5)$$

$$\rightarrow t_{\text{cross}}^{\text{core}} \approx 1.2 \times 10^4 \text{ yr}$$

$$\rightarrow t_{\text{decay}} \approx 10^4 - 10^5 \text{ yr} \ll \tau_{\text{ONC}}$$

- ⇒ Why is there a Trapezium today ?
Or, was the initial core more massive ?

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Summary

- $\alpha = 2.3$ (Salpeter) is **universal** for $m \gtrsim 0.5 M_{\odot}$.
- $CSF > 1$ for massive stars
($CSF \leq 1$ for low-mass stars).
- $CSF \geq 1 \implies \alpha \geq 2.7$
 - Salpeter wrong after all **?!**
 - Massive stars form through coagulation in cluster cores **?**
- **Fraction** of **isolated** massive stars **large**.
Why **?**
 - different mode of SF **?**
 - supernova explosion in \mathcal{B} **?**
 - $\mathcal{B} + \mathcal{B} \longrightarrow (\mathcal{B} + 2S_{ej})$ or $(\mathcal{T}^s + S_{ej})$ **?**
 - $\mathcal{B} + \text{bh} \longrightarrow \mathcal{B}_{\text{bh}} + S_{ej}$ **?**
- **Mass-segregation** theoretically **very rapid**
 \implies **unstable core** of massive stars.
- **Gas expulsion** \implies **cluster expansion & dissolution**, probably very rapid.

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