

The peculiar motions of early-type galaxies in two distant regions – II. The spectroscopic data

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ABSTRACT

We present the spectroscopic data for the galaxies studied in the EFAR project, which is designed to measure the properties and peculiar motions of early-type galaxies in two distant regions. We have obtained 1319 spectra of 714 early-type galaxies over 33 observing runs on 10 different telescopes. We describe the observations and data reductions used to measure redshifts, velocity dispersions and the Mgb and Mg_2 Lick linestrength indices. Detailed simulations and intercomparison of the large number of repeat observations lead to reliable error estimates for all quantities. The measurements from different observing runs are calibrated to a common zero-point or scale before being combined, yielding a total of 706 redshifts, 676 velocity dispersions, 676 Mgb linestrengths and 582 Mg_2 linestrengths. The median estimated errors in the combined measurements are $\Delta cz = 20 \text{ km s}^{-1}$, $\Delta\sigma/\sigma = 9.1$ per cent, $\Delta Mgb/Mgb = 7.2$ per cent and $\Delta Mg_2 = 0.015$ mag. Comparison of our measurements with published data sets shows no systematic errors in the redshifts or velocity dispersions, and only small zero-point corrections to bring our linestrengths on to the standard Lick system. We have assigned galaxies to physical clusters by examining the line-of-sight velocity distributions based on EFAR and ZCAT redshifts, together with the projected distributions on the sky. We derive mean redshifts and velocity dispersions for these clusters, which will be used in estimating distances and peculiar velocities and to test for trends in the galaxy population with cluster mass. The spectroscopic parameters presented here for 706 galaxies combine high-quality data, uniform reduction and measurement procedures, and detailed error analysis. They form the largest single set of velocity dispersions and linestrengths for early-type galaxies published to date.

Key words: surveys – galaxies: clusters: general – galaxies: distances and redshifts – galaxies: elliptical and lenticular, cD – large-scale structure of Universe.

1 INTRODUCTION

We are measuring the peculiar motions of galaxy clusters in the Hercules–Corona Borealis (HCB) and Perseus–Pisces–Cetus (PPC) regions at distances between 6000 and 15 000 km s^{-1} using the global properties of elliptical galaxies. This study (the EFAR project) has as primary goals: (i) characterizing the intrinsic properties of elliptical galaxies in clusters by compiling a large and homogeneous sample with high-quality photometric and spectroscopic data; (ii) testing possible systematic errors, such as

environmental dependence, in existing elliptical galaxy distance estimators; (iii) deriving improved distance estimators based on a more comprehensive understanding of the properties of ellipticals and how these are affected by the cluster environment; and (iv) determining the peculiar velocity field in regions that are dynamically independent of the mass distribution within 5000 km s^{-1} of our Galaxy in order to test whether the large-amplitude coherent flows seen locally are typical of bulk motions in the Universe.

The background and motivation of this work are discussed in

Paper I of this series (Wegner et al. 1996), which also describes in detail the choice of regions to study, the sample of clusters and groups, and the selection procedure and selection functions of the programme galaxies. In earlier papers we reported the photoelectric photometry for 352 programme galaxies which underpins the transformation of our CCD data to the standard R magnitude system (Colless et al. 1993), and described our technique for correcting for the effects of seeing on our estimates of length-scales and surface brightnesses (Saglia et al. 1993). This paper (Paper II) describes the spectroscopic observations and gives redshifts, velocity dispersions and linestrength indices for the programme galaxies. The CCD imaging observations of these galaxies, and their photometric parameters, are described in Paper III (Saglia et al. 1997a), while descriptions of the profile fitting techniques used to determine these parameters (along with detailed simulations establishing the uncertainties and characterizing the systematic errors) are given in Paper IV (Saglia et al. 1997b). The $Mg-\sigma$ relation and its implications are discussed in Paper V (Colless et al. 1999). Subsequent papers in the series will explore other intrinsic properties of the galaxies and their dependence on environment, derive an optimal distance estimator, and discuss the peculiar motions of the clusters in each of our survey regions and their significance for models of the large-scale structure of the Universe.

The structure of the present paper is as follows. In Section 2 we describe the observations and reductions used in obtaining the 1319 spectra in our data set (1250 spectra for 666 programme galaxies, and 69 spectra for 48 calibration galaxies), and discuss the quality of the data. We explain the techniques by which redshifts, velocity dispersions and linestrength indices were estimated from the spectra in Section 3, including the various corrections applied to the raw values. In Section 4 we describe the method used to combine data from different runs and evaluate the internal precision of our results using the large number of repeat measurements in our data set. We then give the final values of the spectroscopic parameters for each galaxy in our sample: we have redshifts for 706 galaxies, dispersions and Mgb linestrengths for 676 galaxies and Mg_2 linestrengths for 582 galaxies. We compare our results with previous studies in the literature to obtain external estimates of our random and systematic errors. In Section 5 we combine our redshifts with those from ZCAT in order to assign sample galaxies to physical clusters, and to estimate the mean redshifts and velocity dispersions of these clusters. Our conclusions are summarized in Section 6.

This paper presents the largest and most homogeneous sample of velocity dispersions and linestrengths for elliptical galaxies ever obtained. The precision of our measurements is sufficiently good to achieve the goal of measuring distances via the Fundamental Plane out to $15\,000\text{ km s}^{-1}$.

2 OBSERVATIONS

The spectroscopic observations for the EFAR project were obtained over a period of seven years from 1986 to 1993 in a total of 33 observing runs on 10 different telescopes. In this section we describe the spectroscopic setups, the observing procedures, the quality of the spectra and the data reduction techniques. Further detail on these points is given by Baggle (1996).

2.1 Spectroscopic setups

Table 1 gives the spectroscopic setup for each run, including the run number, date, telescope, spectrograph and detector, wavelength

range, spectral dispersion (in \AA pixel^{-1}), effective resolution (in km s^{-1}), effective aperture size and number of spectra taken in the run. Note that two runs (116 and 130) produced no useful data and are included in Table 1 only for completeness. Three runs utilized fibre spectrographs: runs 127 and 133 used Argus on the Cerro Tololo Inter-american Observatory (CTIO) 4-m, and run 131 used MEFOS on the European Southern Observatory (ESO) 3.6-m. All the other runs employed long-slit spectrographs, mostly on 2-m-class telescopes [Michigan–Dartmouth–MIT (MDM) Hiltner 2.4-m, Isaac Newton 2.5-m, Kitt Peak 2.1-m, Siding Spring 2.3-m, Calar Alto 2.2-m], although some 4-m-class telescopes were also used [Kitt Peak 4-m, William Herschel 4.2 m, the Multiple Mirror Telescope (MMT)].

The spectra from almost all runs span at least the wavelength range $5126\text{--}5603\text{ \AA}$, encompassing the $Mg\text{ I } b\ 5174\text{-\AA}$ band and the $Fe\text{ I } 5207\text{-}$ and 5269-\AA features in the rest frame for galaxies over the sample redshift range $cz \approx 6000\text{--}15\,000\text{ km s}^{-1}$. The exceptions are the spectra from runs 115 and 131. Run 115 comprises eight spectra obtained at the WHT with the blue channel of the ISIS spectrograph which have a red wavelength limit of 4970 \AA (i.e. including $H\beta$ but not Mgb). Since we have spectra for all these galaxies from other runs, we do *not* use the redshifts and dispersions from run 115. Run 131 comprises 128 spectra obtained at the ESO 3.6-m with the MEFOS fibre spectrograph to a red limit of 5468 \AA , including Mgb and $Fe\text{ I } 5207\text{ \AA}$ over the redshift range of interest, but not $Fe\text{ I } 5269\text{ \AA}$ beyond $cz \approx 11\,000\text{ km s}^{-1}$. For most of the runs the spectra also encompass $H\beta$, and several span the whole range from $Ca\text{ I } H+K\ 3933+3969\text{ \AA}$ to $Na\text{ I } D\ 5892\text{ \AA}$.

The effective instrumental resolution of the spectra, σ_i , was measured from the autocorrelation of stellar template spectra (see Section 3.1 below), and ranged from 80 to 170 km s^{-1} , with a median value of 125 km s^{-1} . Both long-slit and circular entrance apertures were used. Slits were typically $1.7\text{--}2.0$ arcsec wide and the spectra were extracted to the point where the galaxy fell to about 10 per cent of its peak value. Circular apertures (in the fibre spectrographs and the MMT Big Blue spectrograph) were between 1.9 and 2.6 arcsec in diameter. Further details of the observing setup for each telescope/instrument combination are given in Appendix A.

2.2 Observing procedures

The total integration times on programme galaxies varied considerably, depending on telescope aperture, observing conditions and the magnitude and surface brightness of the target (our programme galaxies have R -band total magnitudes in the range $10\text{--}16$). On 2-m-class telescopes (with which the bulk of the spectroscopy was done), exposure times were usually in the range $30\text{--}60$ min, with a median of 40 min; on 4-m-class telescopes, exposure times were generally $15\text{--}20$ min (up to 60 min for the faintest galaxies) with single-object slit spectrographs, but 60 or 120 min with the fibre spectrographs [where the aim was high signal-to-noise ratio (S/N) and completeness]. Slit orientations were not generally aligned with galaxy axes. The nominal goal in all cases was to obtain around 500 photon \AA^{-1} at Mgb , corresponding to a S/N per 100 km s^{-1} resolution element of about 30. In fact our spectra have a median of 370 photon \AA^{-1} at Mgb , corresponding to a S/N per 100 km s^{-1} of 26 (see Section 2.4).

In each run several G8 to K5 giant stars with known heliocentric velocities were observed. These ‘velocity standard stars’ are used as spectral templates for determining redshifts and velocity dispersions. In observing these standards, care was taken to ensure that the illumination across the slit was uniform, in order both to remove

Table 1. The spectroscopic observing runs.

Run #	Date	Telescope ^a	Spectrograph + grating	Detector	$\lambda\lambda$ (Å)	$\Delta\lambda^b$ (Å pixel ⁻¹)	σ_i^c (km s ⁻¹)	Aperture ^d (arcsec)	N^e
101	86/12	MG24	MarkIIIa + 600B	GEC	4912–6219	2.27	145	1.9×10%	22
102	87/03	MG24	MarkIIIa + 600B	RCA	4787–6360	3.10	145	1.9×10%	58
103	87/05	MG24	MarkIIIa + 600B	RCA	4809–6364	3.07	145	1.9×10%	37
104	88/04	MG24	MarkIIIa + 600V	RCA	5025–6500	2.90	125	1.9×10%	12
105	88/06	MG24	MarkIIIa + 600V	RCA	5055–6529	2.90	125	1.9×10%	37
106	88/09	MG24	MarkIIIa + 600V	Thompson	5041–6303	2.21	130	1.9×10%	23
107	88/10	MG24	MarkIIIa + 600V	RCA	5048–6522	2.90	130	1.9×10%	27
108	88/07	MMTB	BigBlue + 300B	Reticon	3700–7200	1.14	135	2.5	10
109	88/11	KP4M	RC + UV-Fast + 17B	TI2	4890–5738	1.07	100	2.0×3.9	104
110	88/11	KP2M	GoldCam + #240	TI5	4760–5879	1.52	105	2.0×3.9	72
111	88/11	MMTB	BigBlue + 300B	Reticon	3890–7500	1.34	135	2.5	20
112	89/04	MG24	MarkIIIb + 600V	RCA	5066–6534	2.91	130	1.7×10%	34
113	89/06	MG24	MarkIIIb + 600V	TI4849	5126–6393	2.18	150	2.4×10%	23
114	89/06	MMTB	BigBlue + 300B	Reticon	3890–7500	1.34	135	2.5	12
115	89/08	WHT4	ISIS-Blue + R600B	CCD-IPCS	4330–4970	0.45	095	2.0×3.9	8
116 ^f	89/10	MMTR	RedChannel + 600B	TI	4750–5950	1.50	125	1.5×10%	7
117	89/10	MG24	MarkIIIb + 600V	Thompson	5031–6300	2.21	130	1.7×10%	61
118	89/11	MG24	MarkIIIb + 600V	RCA	5018–6499	2.89	170	1.7×10%	14
119	90/05	MMTR	RedChannel + 600B	TI	4750–5950	1.50	125	1.5×10%	17
120	90/10	IN25	IDS + 235mm + R632V	GEC6	4806–5606	1.46	100	1.9×7.2	40
121	91/05	IN25	IDS + 235mm + R632V	GEC3	4806–5603	1.46	100	1.9×7.2	87
122	91/10	MG24	MarkIIIb + 600V	Thompson	5018–6278	2.20	125	1.7×10%	43
123	91/11	IN25	IDS + 235mm + R632V	GEC3	4806–5603	1.46	100	1.9×7.2	29
124	92/01	MG24	MarkIIIb + 600V	Thompson	5038–6267	2.15	125	1.7×10%	35
125	92/06	MG24	MarkIIIb + 600V	Loral	4358–7033	1.31	125	1.7×10%	57
126	92/06	CA22	B&C spec + #7	TEK6	4800–6150	1.40	100	5.0×4.2	39
127 ^g	92/09	CT4M	Argus + KPGL#3	Reticon II	3877–6493	2.19	145	1.9	199
128	93/05	MG24	MarkIIIb + 600V	Loral	5090–7050	1.40	105	1.7×10%	24
129	93/06	MG24	MarkIIIb + 600V	TEK	4358–6717	2.31	135	1.7×10%	3
130 ^f	93/06	SS23	DBS-blue + 600B	PCA	5015–5555	0.80	085	2.0×10%	0
131 ^g	93/10	ES36	MEFOS + B&C + #26	TEK512CB	4850–5468	1.22	105	2.6	128
132	93/10	SS23	DBS-blue + 600B	Loral	4820–5910	1.10	080	2.0×10%	14
133 ^g	93/10	CT4M	Argus + KPGL#3	Reticon II	3879–6485	2.19	145	1.9	193

^a Telescopes: MG24 = MDM 2.4-m; KP4M/KP2M = KPNO 4-m/2.1-m; WHT4 = William Herschel 4.2-m (La Palma Observatory); IN25 = Isaac Newton 2.5-m (LPO); MMTB/MMTR = MMT (blue/red); CA22 = Calar Alto 2.2-m; CT4M = Cerro Tololo 4-m; SS23 = Siding Spring 2.3-m; ES36 = ESO 3.6-m.

^b Spectral dispersion in Å pixel⁻¹.

^c Instrumental resolution (σ , not FWHM) in km s⁻¹, as determined from the cross-correlation analysis calibration curves (see Section 3.1).

^d The aperture over which the galaxy spectrum was extracted: diameter for circular apertures and fibres, width×length for rectangular slits (10 per cent means that the spectrum was extracted out to the point where the luminosity had fallen to 10 per cent of its peak value).

^e The number of spectra taken in the run.

^f These runs produced no useful data.

^g These runs used fibre spectrographs.

redshift zero-point errors and to mimic the illumination produced by a galaxy, thereby minimizing systematic errors in velocity dispersion estimates. This was achieved in various ways: by defocusing the telescope slightly, by moving the star back and forth across the slit several times, or by trailing it up and down the slit. Such procedures were not necessary for standards obtained with fibre spectrographs, as internal reflections in the fibres ensure even illumination of the spectrograph for all sources. Very high S/Ns (typically >10 000 photon Å⁻¹) were obtained in order that the stellar templates did not contribute to the noise in the subsequent analysis.

The normal calibration exposures were also obtained: bias frames, flat-fields (using continuum lamps internal to the spectrographs or illuminating the dome) and spectra of wavelength calibration lamps before and/or after each galaxy or star exposure. In general we did not make use of spectrophotometric standards, as

fluxed spectra were not necessary and we wished to minimize overheads as much as possible.

The calibration procedures were slightly different for the three large data sets taken using fibre-fed spectrographs at CTIO (runs 127 and 133) and ESO (run 131). Because of the need to calibrate the relative throughput of the fibres in order to perform sky subtraction, fibre observations always included several twilight sky flat-field exposures. Each velocity standard star was observed through several fibres by moving the fibres sequentially to accept the starlight.

2.3 Reductions

The reductions of both the long-slit and fibre observations followed standard procedures as implemented in the IRAF, MIDAS and Starlink FIGARO software packages. We briefly summarize the main steps in

the reduction of our long-slit and fibre data below; further details can be found in Bagley (1996).

The first stage of the reductions, common to all observations, was to remove the CCD bias using a series of bias frames taken at the start or end of the night. These frames were median-filtered and the result, scaled to the mean level of the CCD overscan strip, was subtracted from each frame in order to remove both the spatial structure in the bias pedestal and temporal variations in its overall level. We also took long dark exposures to check for dark current, but in no case did it prove significant. Subsequent reductions differed somewhat for long-slit and fibre observations.

For long-slit data, the next step was the removal of pixel-to-pixel sensitivity variations in the CCD by dividing by a sensitivity map. This map was produced by median-filtering the flat-field exposures (of an internal calibration lamp or dome lamp) and dividing this by a smoothed version of itself (achieved by direct smoothing or two-dimensional surface fitting) in order to remove illumination variations in the ‘flat’-field. If necessary (because of a long exposure time or a susceptible CCD), cosmic ray events were identified and interpolated over in the two-dimensional image using either algorithmic or manual methods (or both).

The transformation between wavelength and pixel position in long-slit data was mapped using the emission lines in the comparison lamp spectra. The typical precision achieved in wavelength calibration, as indicated by the residuals of the fit to the calibration line positions, was $\lesssim 0.1$ pixel, corresponding to $0.1\text{--}0.3 \text{ \AA}$ or $5\text{--}15 \text{ km s}^{-1}$, depending on the spectrograph setup (see Table 1). The spectra were then rebinned into equal intervals of $\log \lambda$ so that each pixel corresponded to a fixed velocity interval, $\Delta v \equiv c\Delta z = c(10^{\Delta \log \lambda} - 1)$, chosen to preserve the full velocity resolution of the data.

The final steps in obtaining long-slit spectra are sky subtraction and extraction. The sky level was measured from two or more

regions along the slit sufficiently far from the target object to be uncontaminated by its light. To account for variations in transmission along the slit, the sky under the object was interpolated using a low-order curve fitted to the slit illumination profile. A galaxy spectrum was then extracted by summing along the profile, usually over the range where the luminosity of the object was greater than ~ 10 per cent of its peak value, but sometimes over a fixed width in arcsec (see Table 1). Standard star spectra were simply summed over the range along the slit that they had been trailed or defocused to cover.

For the fibre runs the individual object and sky spectra were extracted first, using a centroid-following algorithm to map the position of the spectrum along the CCD. The extraction algorithm fitted the spatial profile of the fibre, in order to remove cosmic ray events and pixel defects, and then performed a weighted sum over this fit out to the points where the flux fell to ~ 5 per cent of the peak value. Next, the dome illumination flat-field spectra were median-combined and a sensitivity map for each fibre constructed by dividing the flat-field spectrum of each fibre by the average over all fibres and normalizing the mean of the result to unity. The pixel-to-pixel variations in the CCD response were then removed by dividing all other spectra from that fibre by this sensitivity map. Wavelength calibration was accomplished using the extracted comparison lamp spectra, giving similar precision to the long-slit calibrations, and the object spectra were rebinned to a $\log \lambda$ scale. Using the total counts through each fibre from the twilight sky flat-field to give the relative throughputs, the several sky spectra obtained in each fibre exposure were median-combined (after manually removing ‘sky’ fibres which were inadvertently placed on faint objects). The resulting high-S/N sky spectrum, suitably normalized to the throughput of each fibre, was then subtracted from each galaxy or standard star spectrum.

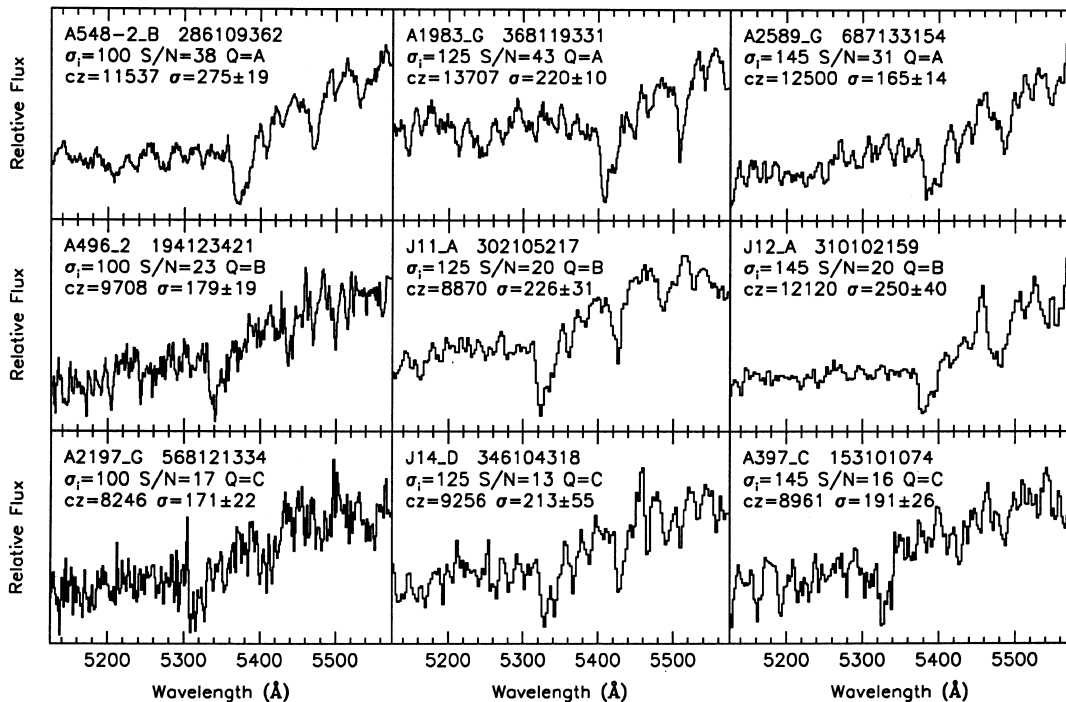


Figure 1. Examples of spectra covering a range of quality classes and instrumental resolutions: the top, middle and bottom rows are spectra with quality classes A, B and C respectively; the left, central and right columns are spectra with resolutions 100, 125 and 145 km s^{-1} respectively. The label for each spectrum gives the galaxy name, the GINRUNSEQ spectrum identifier, the instrumental resolution, the S/N and quality class of the spectrum, the redshift, the dispersion and its estimated error. Note that the panels show relative flux and have a false zero for viewing convenience.

The final step in the reductions for both long-slit and fibre data was manually to clean all the one-dimensional spectra of remaining cosmic ray events or residual sky lines (usually only the 5577-Å line) by linearly interpolating over affected wavelengths.

2.4 Spectrum quality

We have two methods for characterizing the quality of our spectra. One is a classification of the spectra into five quality classes, based on our experience in reducing and analysing such data. Classes A and B indicate that both the redshift and the velocity dispersion are reliable (with class A giving smaller errors than class B); class C spectra have reliable redshifts and marginally reliable dispersions; class D spectra have marginally reliable redshifts but unreliable dispersions; class E spectra have neither redshifts nor dispersions. The second method is based on the S/N per 100 km s⁻¹ bin, estimated approximately from the mean flux over the rest-frame wavelength range used to determine the redshifts and dispersions (see Section 3.1) under the assumption that the spectrum is shot-noise-dominated. These two measures of spectral quality are complementary: the S/N estimate is objective but cannot take into account qualitative problems which are readily incorporated in the

subjective classifications. Fig. 1 shows example spectra covering a range of quality classes and instrumental resolutions.

Fig. 2 shows the S/N distribution for the whole sample and for each quality class individually, and gives the total number of objects, the fraction of the sample and the median S/N in each class. For the whole sample, 39 per cent of the spectra have S/N > 30, 70 per cent have S/N > 20, and 96 per cent have S/N > 10. The two quality measures are clearly correlated, in the sense that better-classed spectra tend to have higher S/N. However there is also considerable overlap in the S/N range spanned by the different classes. This overlap has various sources: (i) factors other than S/N that affect the quality of the redshift and dispersion estimates, notably the available rest-frame spectral range (which depends on both the spectrograph setup and the redshift of the target) and whether the object has emission lines; (ii) errors in estimating the S/N [e.g. owing to sky subtraction errors, the neglect of the sky contribution in computing the S/N for fainter galaxies, or uncertainties in the CCD gain (affecting the conversion from counts to photons)]; (iii) subjective uncertainties in the quality classification, particularly in determining the reliability of dispersion estimates (i.e. between classes B and C). Both ways of determining spectral quality are therefore needed in order to estimate the

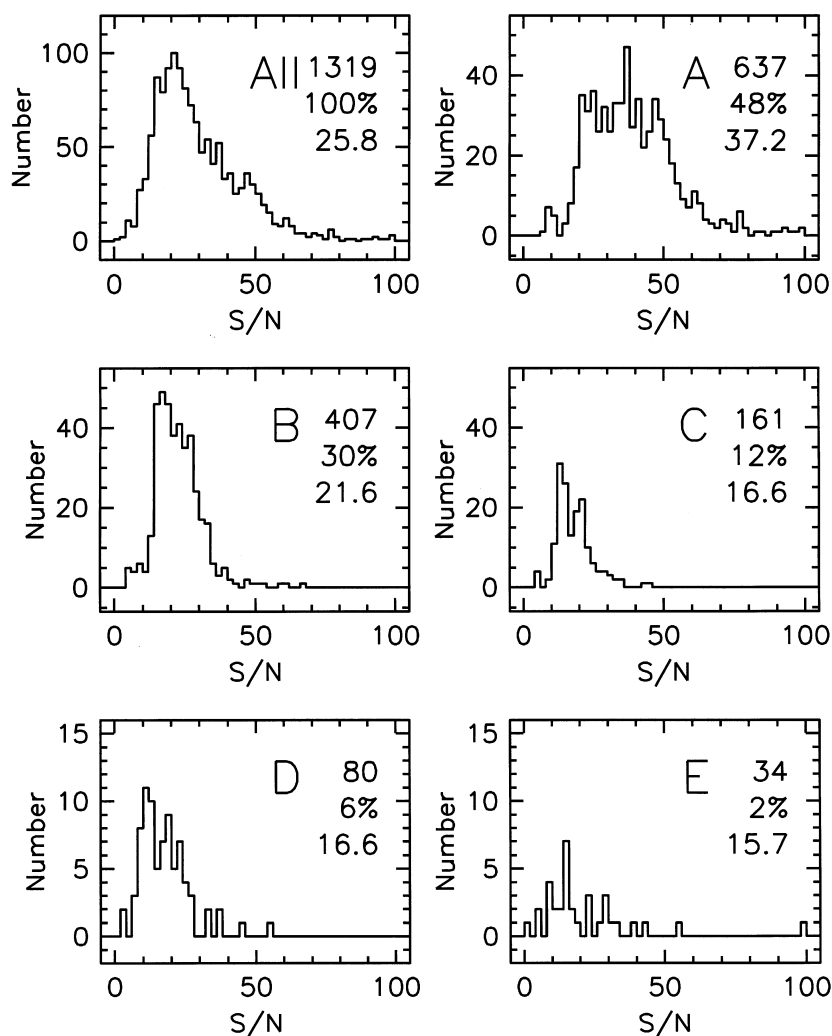


Figure 2. The distribution of S/N with quality class. For each class the panels give the total number of spectra, the percentage of the whole sample and the median S/N.

reliability and precision of the spectroscopic parameters that we measure.

3 ANALYSIS

3.1 Redshifts and dispersions

We derived redshifts and velocity dispersions from our spectra using the `FXCOR` utility in `IRAF`, which is based on the cross-correlation method of Tonry & Davis (1979). We preferred this straightforward and robust method to more elaborate techniques since it is well-suited to the relatively modest S/N of our spectra. We used a two-step procedure, obtaining an initial estimate of the redshift using the whole available spectrum and then using a fixed rest-frame wavelength range for the final estimates of redshift and velocity dispersion. The procedure was applied in a completely uniform manner to all the spectra in our sample as far as differences in wavelength range and resolution would allow.

The first step in the cross-correlation analysis is to fit and subtract the continuum of each spectrum in order to avoid the numerical difficulties associated with a dominant low-frequency spike in the Fourier transform. In the first pass through `FXCOR` the continuum shape was fitted with a cubic spline, with the number of segments along the spectrum chosen so that each segment corresponded to about 8000 km s^{-1} . Each iteration of the fit excluded points more than 1.5σ below or 3σ above the previous fit. In this way we achieved a good continuum fit without following broad spectral features. We then apodized 10 per cent of the spectrum at each end with a cosine bell before padding the spectrum to 2048 pixels with zeros.

This continuum-subtracted, apodized spectrum was then Fourier transformed and a standard ‘ramp’ filter applied. This filter is described by four wavenumbers (k_1, k_2, k_3, k_4) , rising linearly from 0 to 1 between k_1 and k_2 and then falling linearly from 1 to 0 between k_3 and k_4 . In the first pass these wavenumbers were chosen to be $k_1 = 4\text{--}8$ and $k_2 = 9\text{--}12$ (tailored to remove residual power from the continuum without affecting broad spectral features), and $k_3 = N_{\text{pix}}/3$ and $k_4 = N_{\text{pix}}/2$ (N_{pix} is the number of pixels in original spectrum before it is padded to 2048 pixels; these choices attenuate high-frequency noise and eliminate power beyond the Nyquist limit at $N_{\text{pix}}/2$). The same procedures were also applied to the spectrum of the stellar velocity standard to be used as a template. The cross-correlation of the galaxy and stellar template was then computed, and the top 90 per cent of the highest cross-correlation peak fitted with a Gaussian in order to obtain a redshift estimate.

This procedure was repeated for every template from that run, and the redshifts corrected to the heliocentric frame. Offsets in the velocity zero-point between templates, measured as the mean difference in the redshifts measured with different templates for all the galaxies in the run, were typically found to be $\lesssim 30 \text{ km s}^{-1}$. These were brought into relative agreement within each run by choosing the best-observed K0 template as defining the fiducial velocity zero-point. Applying these offsets brought the galaxy redshifts estimated from different templates into agreement to within $\lesssim 3 \text{ km s}^{-1}$. (The removal of run-to-run velocity offsets is described below.) The mean over all templates then gave the initial redshift estimate for the galaxy.

This initial redshift was then used to determine the wavelength range corresponding to the rest-frame range $\lambda_{\text{min}} = 4770 \text{ \AA}$ to $\lambda_{\text{max}} = 5770 \text{ \AA}$. This range was chosen for use in the second pass through `FXCOR` because: (i) it contains the `Mg I b` 5174- \AA band, `H β` 4861 \AA and the `Fe I` 5207- and 5269- \AA lines, but excludes the `Na I D` line at 5892 \AA , which gives larger velocity dispersions than the lines

in the region of `Mg b` (Faber & Jackson 1976); (ii) for redshifts up to our sample limit of $cz = 15000 \text{ km s}^{-1}$ this rest-frame wavelength range is included in the great majority of our spectra. The input for the second pass was thus the available spectrum within the range corresponding to rest-frame 4770–5770 \AA . All but two of our runs cover the rest frame out to at least 5330 \AA for $cz = 15000 \text{ km s}^{-1}$; the exceptions are run 115 (which is not used for measuring dispersions) and run 131 (which reaches rest-frame 5207 \AA).

In the second pass through `FXCOR` we employed only minimal continuum subtraction based on a one- or two-segment cubic spline fit, preferring the better control over continuum suppression afforded by more stringent filtering at low wavenumbers. After considerable experimentation and simulation, we found that the best filter for recovering velocity dispersions was a ramp with the same k_3 and k_4 values as in the first pass, but with $k_2 = 0.01(N_{\text{max}} - N_{\text{min}})$, where N_{min} and N_{max} are the pixels corresponding to λ_{min} and λ_{max} , and $k_1 = 0.75k_2$. Again, the top 90 per cent of the highest cross-correlation peak was fitted with a Gaussian. The position of this peak, corrected for the motion of the template star and the heliocentric motion of the Earth relative to both the template and the galaxy, gave the final redshift estimate.

The velocity dispersion of the galaxy, σ_g , is in principle related to the dispersion of the Gaussian fitted to the cross-correlation peak, σ_x , by $\sigma_x^2 = \sigma_g^2 + 2\sigma_i^2$ (where σ_i is the instrumental resolution: Tonry & Davis 1979). In practice this relationship needs to be calibrated empirically because of the imperfect match between the spectra of a broadened stellar template and a galaxy and the effects of the filter applied to both spectra. The calibration relation between σ_x and σ_g for a typical case is shown in Fig. 3 (see the caption for more details). We estimate the instrumental resolution for a given run from the mean value of the calibration curve intercepts for all the templates in the run ($\sigma_i \approx \sigma_x/\sqrt{2}$ when $\sigma_g = 0$); these are the values listed in Table 1.

The values of heliocentric radial velocity and velocity dispersion were determined in this second pass through `FXCOR` for each galaxy spectrum using all the templates in the same run. The final step is then to combine the redshift and dispersion estimates from each template, as summarized below.

For the redshifts the steps involved were as follows. (i) Cases in which the ratio of cross-correlation function peak height to noise [the R parameter defined by Tonry & Davis (1979)] was less than 2 were rejected, as were cases that differed from the median by more than a few hundred km s^{-1} . (ii) The mean offset between the redshifts from a fiducial K0 template and each other template was used to shift all the redshifts from the other template to the velocity zero-point of the fiducial. These offsets were typically $\lesssim 50 \text{ km s}^{-1}$. (iii) A mean redshift for each galaxy was then computed from all the unrejected cases using two-pass 2σ clipping. (iv) Any template that gave consistently discrepant results was rejected and the entire procedure repeated. The scatter in the redshift estimates from different templates after this procedure was typically a few km s^{-1} .

A very similar procedure was followed in combining velocity dispersions, except that a scalefactor rather than an offset was applied between templates. (i) Cases with $R < 4$ were rejected. (ii) The mean ratio between the dispersions from a fiducial K0 template and each other template was used to scale all the dispersions from the other template to the dispersion scale of the fiducial. These dispersion scales differed by less than 5 per cent for 90 per cent of the templates. (iii) A mean dispersion for each galaxy was then computed from all the unrejected cases using two-pass 2σ clipping. (iv) Any template with a scale differing by more than 10 per cent from the mean was rejected as being a poor match to the

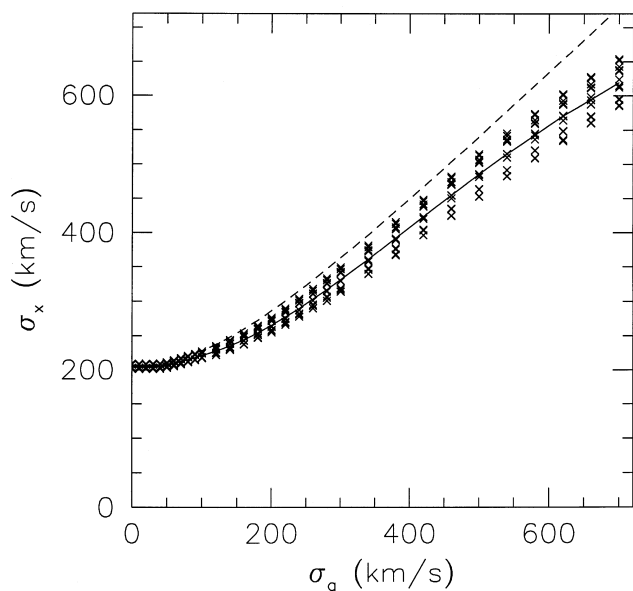


Figure 3. A typical calibration curve showing the relation between the width of the cross-correlation peak, σ_x , and the true velocity dispersion of the galaxy, σ_g . The crosses are the individual calibrations obtained by broadening each of the other templates in the run and cross-correlating with the template being calibrated. The solid curve is the calibration curve used, a series of linear segments joining the median value of σ_x at each calibrated value of σ_g . The dashed curve is the theoretical relation when no filtering is applied, $\sigma_x^2 = \sigma_g^2 + 2\sigma_i^2$, where σ_i is the instrumental resolution, in this case 145 km s^{-1} . Note that the calibration curve flattens for $\sigma_g < \sigma_i$, indicating that the true dispersion becomes increasingly difficult to recover as it drops below the instrumental resolution.

programme galaxies, and the entire procedure was then repeated. (Note that no significant correlation was found between scalefactor and spectral type over the range G8 to K5 spanned by our templates.) The scatter in the dispersion estimates from different templates after this procedure was typically 3–4 per cent.

Two corrections need to be applied to the velocity dispersions before they are fully calibrated: (i) an aperture correction to account for different effective apertures sampling different parts of the galaxy velocity dispersion profile; and (ii) a run correction to remove systematic scale errors between different observing setups. The latter type of correction is also applied to the redshifts to give them a common zero-point. These two corrections are discussed below in Sections 3.4 and 3.5 respectively.

3.2 Linestrength indices

Once redshifts and velocity dispersions were determined, linestrength indices could also be measured using the prescription given by González (1993). This is a refinement of the original ‘Lick’ system in which a standard set of bands was defined for measuring linestrength indices for 11 features in the spectra of spheroidal systems (Burstein et al. 1984). González (1993), Worthey (1993) and Worthey et al. (1994) describe how this system has been updated and expanded to a set of 21 indices. Here we measure both the *Mgb* and *Mg₂* indices.

The feature bandpass for *Mgb* index is $5160.1\text{--}5192.6 \text{ \AA}$, encompassing the Mg I triplet with components at 5166.6 , 5172.0 and 5183.2 \AA . The continuum on either side of the absorption feature is

defined in bands covering $5142.6\text{--}5161.4$ and $5191.4\text{--}5206.4 \text{ \AA}$. *Mgb* is an *atomic* index, and so is defined as the equivalent width of the feature in ångstroms,

$$Mgb = \int \left[1 - \frac{S(\lambda)}{C(\lambda)} \right] d\lambda, \quad (1)$$

where the integral is over the feature bandpass, $S(\lambda)$ is the object spectrum and $C(\lambda)$ is the linear pseudo-continuum defined by interpolating between two continuum estimates, taken at the mid-points of the blue and red continuum bands to be the mean values of the observed spectrum in those bands.

Closely related to *Mgb* is the *Mg₂* index, for which the feature bandpass is $5154.1\text{--}5196.6 \text{ \AA}$ and the continuum bands are $4895.1\text{--}4957.6$ and $5301.1\text{--}5366.1 \text{ \AA}$. This index measures both the Mg I atomic absorption and the broader MgH molecular absorption feature. *Mg₂* is a *molecular* index, and so is defined as the mean ratio of flux to local continuum in magnitudes,

$$Mg_2 = -2.5 \log_{10} \left[\frac{\int S(\lambda)/C(\lambda) d\lambda}{\Delta\lambda} \right], \quad (2)$$

where the integral is over the *Mg₂* feature bandpass, $\Delta\lambda = 42.5 \text{ \AA}$ is the width of that bandpass, and the pseudo-continuum is interpolated from the *Mg₂* continuum bands.

In fact we will often find it convenient to express the *Mgb* index in magnitudes rather than as an equivalent width. By analogy with the *Mg₂* index, we therefore define *Mgb'* to be

$$Mgb' = -2.5 \log_{10} \left(1 - \frac{Mgb}{\Delta\lambda} \right), \quad (3)$$

where in this case $\Delta\lambda = 32.5 \text{ \AA}$, the width of the *Mgb* feature bandpass.

In passing it should be noted that a different definition of linestrength indices has sometimes been used (e.g. Worthey 1994, equations 4 and 5), in which the integral of the ratio of the object spectrum and the continuum in equations (1) and (2) is replaced by the ratio of the integrals. This alternative definition has merits (such as simplifying the error properties of measured indices), but it is not mathematically equivalent to the standard definition. In practice, however, the two definitions generally give linestrengths with negligibly different numerical values.

It is usual in studies of this sort to employ the *Mg₂* index as the main indicator of metallicity and star formation history. However, we find it useful for operational reasons also to measure the *Mgb* index. One problem is that the limited wavelength coverage of the spectra from some runs means that in a number of cases we cannot measure the *Mg₂* index (requiring as it does a wider wavelength range), although we can measure the *Mgb* index. We obtain *Mgb* for 676 objects (with 299 having repeat measurements) and *Mg₂* for 582 objects (with 206 having repeat measurements). Another problem with *Mg₂* is that the widely separated continuum bands make it more susceptible than *Mgb* to variations in the non-linear continuum shape of our unfluxed spectra, which result from using a variety of different instruments and observing galaxies over a wide range in redshift. We therefore present measurements of both *Mgb* and *Mg₂*: the former because it is better determined and available for more sample galaxies; the latter for comparison with previous work. As previously demonstrated (Gorgas, Efstathiou & Salamanca 1990; Jørgensen 1997) and confirmed here, *Mgb* and *Mg₂* are strongly correlated, and so can to some extent be used interchangeably.

Several corrections must be applied to obtain a linestrength measurement that is calibrated to the standard Lick system. The

first correction allows for the fact that the measured linestrength depends on the instrumental resolution. Since all our spectra were obtained at higher resolution than the spectra on which the Lick system was defined, we simply convolve our spectra with a Gaussian of dispersion $(\sigma_{\text{Lick}}^2 - \sigma_i^2)^{1/2}$ in order to broaden our instrumental resolution σ_i (see Table 1) to the Lick resolution of 200 km s^{-1} .

The second correction allows for the fact that the measured linestrength depends on the internal velocity dispersion of the galaxy – a galaxy with high enough velocity dispersion σ_g will have features broadened to the point that they extend outside their index bandpasses, and so their linestrengths will be underestimated. Moreover, if an absorption feature is broadened into the neighbouring continuum bands then the estimated continuum will be depressed and the linestrength will be further reduced. The ‘ σ -correction’ needed to calibrate out this effect can be obtained either by measuring linestrength as a function of velocity broadening for a set of suitable stellar spectra (such as the templates obtained for measuring redshifts and dispersions) or by modelling the feature in question.

Although most previous studies have adopted the former approach, we prefer to use a model to calibrate our indices, since we observe a dependence of the *Mgb* profile shape on σ that is not taken into account by simply broadening stellar templates. Our simple model assumes *Mgb* to be composed of three Gaussians centred on the three *Mg I* lines at $\lambda_b = 5166.6 \text{ \AA}$, $\lambda_c = 5172.0 \text{ \AA}$ and $\lambda_r = 5183.2 \text{ \AA}$ with corresponding relative strengths varying linearly with dispersion from 1.0:1.0:1.0 at $\sigma = 100 \text{ km s}^{-1}$ to 0.2:1.0:0.7 at $\sigma = 300 \text{ km s}^{-1}$. This dependence on dispersion is empirically determined and approximate (the relative strengths of the individual lines are not tightly constrained), but it does significantly improve the profile fits compared with assuming any fixed set of relative weights. Such variation of the *Mgb* profile shape reflects changes, as a function of velocity dispersion, in the stellar population mix and relative abundances (particularly of *Mg*, *C*, *Fe* and *Cr*), which each affect the profile in complex ways (Tripicco & Bell 1995).

Using the estimated value of the index to normalize the model profile and the effective dispersion $(\sigma_g^2 + \sigma_{\text{Lick}}^2)^{1/2}$ to give the broadening, we can estimate both the profile flux that is broadened out of the feature bandpass and the resulting depression of the continuum. Correcting for both these effects gives an improved estimate for the linestrength. Iterating leads rapidly to convergence and an accurate σ -correction for the *Mgb* and *Mg₂* indices. We find that the *Mgb* σ -correction is typically +4 per cent at 100 km s^{-1} and increases approximately linearly to +16 per cent at 400 km s^{-1} ; the *Mg₂* σ -correction is typically 0.000 mag up to 200 km s^{-1} and increases approximately linearly to 0.004 mag at 400 km s^{-1} .

Note that the usual method of determining the σ -correction by broadening standard stars ignores the dependence of profile shape on changes in the stellar population mix as a function of luminosity or velocity dispersion. Our tests indicate that, by doing so, the usual method tends to overestimate *Mgb* for galaxies with large dispersions: by 2 per cent at 200 km s^{-1} , 6 per cent at 300 km s^{-1} and 14 per cent at 400 km s^{-1} . The two methods give essentially identical results for *Mg₂*, since it has much smaller σ -corrections owing to its wider feature bandpass and well-separated continuum bands.

The other corrections that need to be applied to the linestrength estimates are: (i) an aperture correction to account for different effective apertures sampling different parts of the galaxy (Section 3.4); (ii) a run correction to remove systematic scale errors between different observing setups (Section 3.5); and

(iii) an overall calibration to the Lick system determined by comparisons with literature data (Section 4.3).

3.3 Error estimates

Error estimates for our redshifts, velocity dispersions and linestrengths come from detailed Monte Carlo simulations of the measurement process for each observing run. By calibrating the errors estimated from these simulations against the rms errors obtained from the repeat measurements that are available for many of the objects (see Section 3.6), we can obtain precise and reliable error estimates for each measurement of every object in our sample.

The procedure for estimating the uncertainties in our redshifts and velocity dispersions was as follows. For each stellar template in each observing run, we constructed a grid of simulated spectra with Gaussian broadenings of $100\text{--}300 \text{ km s}^{-1}$ in 20 km s^{-1} steps and continuum counts corresponding to S/N ratios of $10\text{--}90$ in steps of 10. For each spectrum in this grid we generated 16 realizations assuming Poisson noise. These simulated spectra were then cross-correlated against all the other templates from the run in order to derive redshifts and velocity dispersions in the standard manner. The simulations do not account for spectral mismatch between the galaxy spectra and the stellar templates, but for well-chosen templates this effect is only significant at higher S/N than is typically found in our data.

Fig. 4 shows the random error in redshift and the systematic and random errors in dispersion as functions of input dispersion and S/N for four of the larger runs. The systematic errors in redshift are not shown as they are negligibly small ($\sim 1 \text{ km s}^{-1}$), although the simulations do not include possible zero-point errors. The systematic errors in dispersion are generally small (a few per cent or less) for $\text{S/N} > 20$, but become rapidly larger at lower S/N. The random errors in redshift increase for lower S/N and higher dispersion, while the random errors in dispersion increase for lower S/N but have a broad minimum at around twice the instrumental dispersion. These curves have the general form predicted for the random errors from the cross-correlation method (Tonry & Davis 1979; Colless 1987).

Given the dispersion and S/N measured for a spectrum, we interpolated the error estimates from the simulation for that particular observing run to obtain the systematic and random errors in each measured quantity. We used the results of these simulations to correct the systematic errors in the velocity dispersions and to estimate the uncertainties in individual measurements of redshift and dispersion. For quality class D measurements of redshifts, where the spectra are too poor to estimate a dispersion and hence a reliable redshift error, we take a conservative redshift error of 50 km s^{-1} .

The linestrength error estimates were obtained by generating 50 Monte Carlo realizations of the object spectrum with Poisson noise appropriate to the S/N level of the spectrum. The *Mgb* and *Mg₂* linestrengths were then measured for each of these realizations and the error estimated as the rms error of these measurements about the observed value. The error estimate obtained in this fashion thus takes into account the noise level of the spectrum, but does not account for errors in the linestrength owing to errors in the redshift and dispersion estimates, nor for systematic run-to-run differences in the underlying continuum shape.

The estimated errors in the spectroscopic parameters are compared with, and calibrated to, the rms errors derived from repeat observations in Section 3.6.

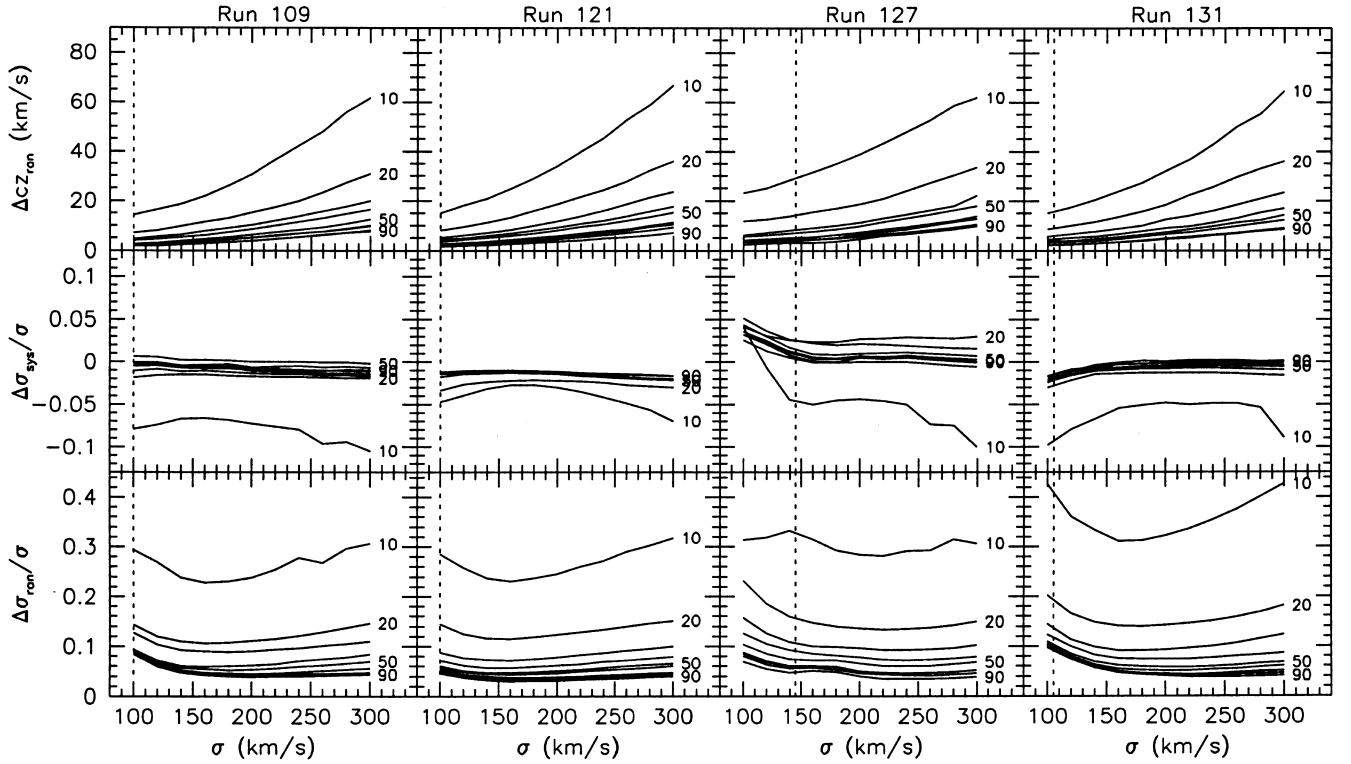


Figure 4. Redshift and dispersion errors as functions of input dispersion and S/N (labelling the curves) from the simulations of four of the larger runs. The top panels show the random error in the redshift, and the centre and bottom panels show the systematic and random errors in the dispersion. The vertical dotted lines indicate the instrumental dispersion of each run.

3.4 Aperture corrections

The velocity dispersion measured for a galaxy is the luminosity-weighted velocity dispersion integrated over the region of the galaxy covered by the spectrograph aperture. It therefore depends on: (i) the velocity dispersion profile; (ii) the luminosity profile; (iii) the distance of the galaxy; (iv) the size and shape of the spectrograph aperture; and (v) the seeing in which the observations were made. In order to intercompare dispersion measurements it is therefore necessary to convert them to a standard scale. The ‘aperture correction’ that this requires has often been neglected because it depends in a complex manner on a variety of quantities, some of which are poorly known. The neglect of such corrections may account in part for the difficulties often found in reconciling dispersion measurements from different sources.

The aperture correction applied by Davies et al. (1987) was derived by measuring dispersions for a set of nearby galaxies through apertures of 4×4 arcsec² and 16×16 arcsec². In this way they used their nearby galaxies to define the velocity dispersion profile and obtained a relation between the corrected value, σ_{cor} , and the observed one, σ_{obs} . This turned out to be an approximately linear relation amounting to a 5 per cent correction over the distance range between Virgo and Coma.

More recently Jørgensen, Franx & Kjægaard (1995) have derived an aperture correction from kinematic models based on data in the literature. Published photometry and kinematics for 51 galaxies were used to construct two-dimensional models of the surface brightness, velocity dispersion, and rotational velocity projected on the sky. They found that the position angle only gave rise to 0.5 per cent variations in the derived dispersions and could thus be ignored. They converted rectangular apertures into an

‘equivalent circular aperture’ of radius r_{ap} which the models predicted would give the same dispersion as the rectangular slit. They found that to an accuracy of 4 per cent one could take $r_{\text{ap}} = 1.025(xy/\pi)^{1/2}$, where x and y are the width and length of the slit.

From their models they then calculated the correction factor from the observed dispersion to the dispersion in some standard aperture. For a standard *metric* aperture, they found this aperture correction to be well approximated by a power law of the form

$$\frac{\sigma_{\text{cor}}}{\sigma_{\text{obs}}} = \left[\left(\frac{r_{\text{ap}}}{r_0} \right) \left(\frac{cz}{cz_0} \right) \right]^{0.04}, \quad (4)$$

where σ_{obs} and σ_{cor} are the observed and corrected dispersions, r_0 is a standard aperture radius, defined to be 1.7 arcsec, and cz_0 is a standard redshift, defined as the redshift of Coma. The standard metric aperture is thus $0.54 h^{-1}$ kpc in radius. Alternatively, one can correct to a standard *relative* aperture (defined to be $R_c/8$) using the same power-law relation,

$$\frac{\sigma_{\text{cor}}}{\sigma_{\text{obs}}} = \left(\frac{r_{\text{ap}}}{R_c/8} \right)^{0.04}. \quad (5)$$

This power law approximates the true relation to within 1 per cent over the observed range of effective apertures [compare the distribution metric aperture sizes in Fig. 5(a) with fig. 4(c) of Jørgensen et al.].

We also apply an aperture correction to our linestrengths. Jørgensen et al. noted that the radial gradient in the Mg_2 index is similar to the radial gradient in $\log \sigma$, and so applied the same aperture correction for Mg_2 as for $\log \sigma$. We adopt this procedure for Mg_2 . For Mgb we convert to Mgb' (equation 3) and, assuming that the radial profile of Mgb' is similar to that of Mg_2 (and hence $\log \sigma$),

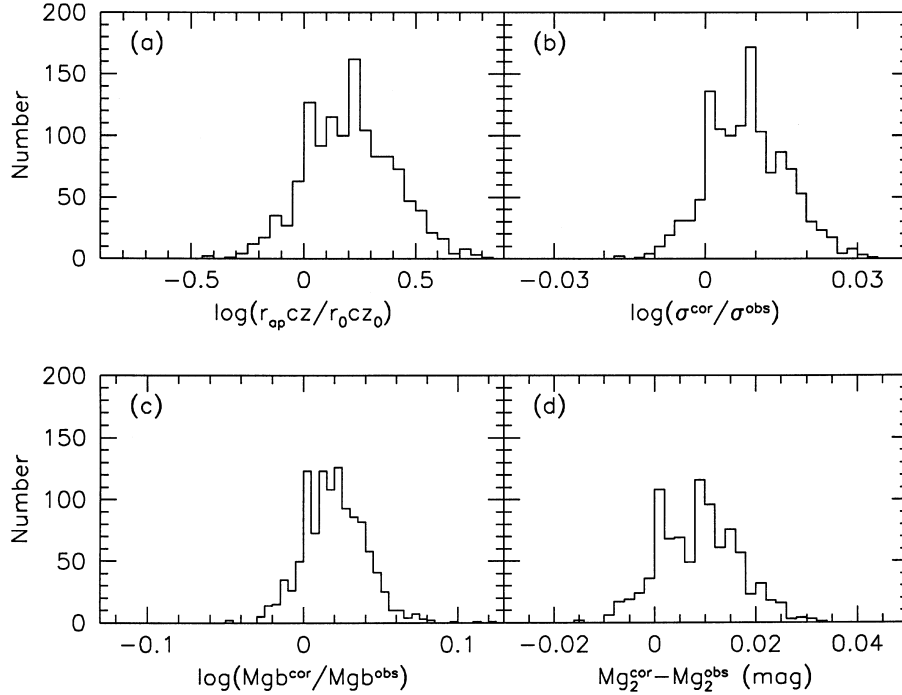


Figure 5. The distribution of aperture corrections to a standard metric aperture. (a) The distribution of the ratio of the observed metric apertures to the standard metric aperture (corresponding to 1.7 arcsec at the redshift of Coma). (b)–(d) The aperture corrections to this standard metric aperture for the dispersion, Mgb and Mg_2 measurements. Note that σ^{cor} is the aperture-corrected dispersion and σ^{obs} is the raw observed dispersion; likewise for the linestrengths.

we apply the $\log \sigma$ aperture correction to Mgb' before converting back to Mgb .

The distributions of corrections to the standard metric aperture for the dispersions and linestrengths are shown in Figs 5(b)–(d). These corrections are generally positive, as most objects in our sample are observed through larger effective apertures and are further away than Jørgensen et al.’s standard aperture and redshift. The corrections to standard relative apertures are quite similar, although having slightly greater amplitude and range. We choose to adopt the correction to a standard metric aperture in order to minimize the size and range of the corrections and to facilitate comparisons with dispersions and linestrengths in the literature.

3.5 Combining different runs

In comparing the redshifts, dispersions and linestrengths obtained from different runs, we found some significant systematic offsets. The origin of these run-to-run offsets is not fully understood. For the redshifts, the use of different velocity standard stars as the fiducials in different runs clearly contributes some systematic errors. For the dispersions, the calibration procedure that we use should in principle remove instrumental systematics; in practice, scale differences are common, as is shown by the range of scalefactors needed to reconcile velocity dispersions from various sources in the compilation by McElroy (1995; see Table 2).

We cannot directly calibrate the measurements from each run to the system defined by a chosen fiducial run, as there is no run with objects in common with all other runs to serve as the fiducial. Instead, we use the mean offset, Δ , between the measurements from any particular run and *all* the other runs. To compute this offset we separately compute, for each galaxy i , the error-weighted mean value of the measurements obtained from the run in question, x_{ij} ,

and from all other runs, y_{ik} :

$$\langle x_i \rangle = \frac{\sum_j x_{ij} / \delta_{ij}^2}{\sum_j 1 / \delta_{ij}^2}, \quad \langle y_i \rangle = \frac{\sum_k y_{ik} / \delta_{ik}^2}{\sum_k 1 / \delta_{ik}^2}. \quad (6)$$

Here j runs over the m_i observations of galaxy i in the target run and k runs over the n_i observations of galaxy i in all other runs; δ_{ij} and δ_{ik} are the estimated errors in x_{ij} and y_{ik} . We then take the average over all galaxies, weighting by the number of comparison pairs, to arrive at an estimate for the offset of the target run:

$$\Delta = \frac{\sum_i m_i n_i (\langle x_i \rangle - \langle y_i \rangle)}{\sum_i m_i n_i}. \quad (7)$$

Here i runs over the l galaxies in the sample. We can reject outliers at this point by excluding galaxies for which the difference $\langle x_i \rangle - \langle y_i \rangle$ is larger than some cut-off: for cz , $\log \sigma$, Mgb' and Mg_2 we required differences of less than 300 km s^{-1} , 0.2 dex, 0.1 mag and 0.1 mag respectively. The uncertainty, ϵ , in this estimate of the run offset is given by

$$\epsilon^2 = \frac{\sum_i (m_i n_i)^2 (\delta \langle x_i \rangle^2 + \delta \langle y_i \rangle^2)}{(\sum_i m_i n_i)^2}, \quad (8)$$

where $\delta \langle x_i \rangle$ and $\delta \langle y_i \rangle$ are the error-weighted uncertainties in $\langle x_i \rangle$ and $\langle y_i \rangle$ given by

$$\delta \langle x_i \rangle^2 = (\sum_j \delta_{ij}^{-2})^{-1}, \quad \delta \langle y_i \rangle^2 = (\sum_k \delta_{ik}^{-2})^{-1}. \quad (9)$$

We subtract the offset determined in this manner from each run and then iterate the whole procedure until there are no runs with residual offsets larger than 0.5ϵ . As a final step, we place the entire data set (now corrected to a common zero-point) on to a fiducial system by subtracting from all runs the offset of the fiducial system. Note that the run corrections for dispersion and Mgb are determined in terms of offsets in $\log \sigma$ and Mgb' .

In order to maximize the number of objects with multiple

measurements, we included the data set from the ‘Streaming Motions of Abell Clusters’ project (SMAC: Hudson, private communication; see also Smith et al. 1997) in this analysis. There is a considerable overlap between the SMAC and EFAR samples which significantly increases the number of comparison observations and reduces the uncertainties in the run offsets. We chose to use the ‘Lick’ system of Davies et al. (1987; included in the SMAC data set) as our fiducial, in order to bring the 7 Samurai, EFAR and SMAC data sets on to a single common system. This is not possible with Mgb , which is not measured in most previous work or by SMAC. We therefore chose run 109 (the Kitt Peak 4-m run of 1988 November) as the Mgb fiducial because it had a large number of high-quality observations and the systematics of the slit spectrograph are believed to be well understood.

We checked that this procedure gives relative run corrections consistent with those obtained by directly comparing runs in those cases where there *are* sufficient objects in common. We have also compared our method with a slightly different method used by the SMAC collaboration to determine the run corrections for their own data, and found good agreement (Hudson, private communication). We carried out Monte Carlo simulations of the whole procedure in order to check the uncertainties in the offsets computed according to equation (8). We found that this equation in general provides a good estimate of the uncertainties, although when the number of

comparisons is small or they involve a small number of other runs it can underestimate the uncertainties by up to 30 per cent. Our final estimates of the uncertainties are therefore derived as the rms of the offsets from 100 Monte Carlo simulations.

Table 2 lists the offsets for each run computed according to the above procedure, their uncertainties based on Monte Carlo simulations, the number of individual measurements (N) and the number of comparison pairs (N^c). Note that to correct our observed measurements to the fiducial system we *subtract* the appropriate run offset in Table 2 from each individual measurement. Of the 31 spectroscopic runs with usable data, only runs 104 and 129 have no objects in common with other runs and hence no run corrections; run 118 has no Mg_2 measurements in common and so no run correction for Mg_2 .

Weighting by the number of individual measurements in each run, the mean amplitude of the corrections and their uncertainties are $28 \pm 8 \text{ km s}^{-1}$ in cz , 0.023 ± 0.015 dex in $\log \sigma$, 0.008 ± 0.006 mag in Mgb' and 0.015 ± 0.006 mag in Mg_2 . The significance of the individual run corrections (in terms of the ratio of the amplitude of the offset to its uncertainty) varies; however, over all runs the reduced χ^2 is highly significant: 15.7, 4.0, 3.3 and 11.4 for the corrections to the redshifts, dispersions, Mgb and Mg_2 respectively. Application of the run corrections reduces the median rms error amongst those objects with repeat measurements from 18 to

Table 2. Calibration of observing runs to a common system.

Run	Δcz (km s^{-1})	N_z	N_z^c	$\Delta \log \sigma$ (dex)	N_σ	N_σ^c	$\Delta Mgb'$ (mag)	N_b	N_b^c	ΔMg_2 (mag)	N_2	N_2^c
101	-9 ± 9	19	62	$+0.014 \pm 0.016$	18	65	$+0.009 \pm 0.009$	18	32	-0.008 ± 0.007	16	30
102	$+62 \pm 10$	56	57	$+0.034 \pm 0.023$	56	56	$+0.012 \pm 0.011$	56	24	-0.013 ± 0.010	56	12
103	-43 ± 10	36	53	-0.034 ± 0.028	34	48	-0.002 ± 0.011	34	27	-0.025 ± 0.009	34	20
104 ^b	+0	11	0	+0.000	11	0	+0.000	11	0	+0.000	11	0
105	$+16 \pm 9$	36	54	$+0.018 \pm 0.019$	36	55	$+0.020 \pm 0.007$	36	37	$+0.008 \pm 0.009$	21	14
106	-17 ± 14	23	61	$+0.024 \pm 0.029$	22	54	$+0.009 \pm 0.007$	22	41	$+0.001 \pm 0.014$	14	13
107	-19 ± 10	27	61	$+0.060 \pm 0.033$	27	64	$+0.000 \pm 0.009$	27	33	-0.003 ± 0.008	21	24
108	-71 ± 23	9	35	-0.055 ± 0.035	9	34	-0.005 ± 0.008	9	12	$+0.032 \pm 0.008$	9	11
109 ^d	$+1 \pm 4$	93	222	-0.015 ± 0.005	92	220	$+0.000 \pm 0.000$	92	167	$+0.008 \pm 0.002$	92	126
110	$+3 \pm 6$	71	186	-0.010 ± 0.008	61	171	$+0.004 \pm 0.004$	61	78	-0.009 ± 0.004	61	50
111	-10 ± 14	19	72	-0.024 ± 0.024	19	76	-0.005 ± 0.006	19	33	$+0.046 \pm 0.006$	19	25
112	$+45 \pm 8$	31	82	-0.006 ± 0.008	31	103	$+0.001 \pm 0.009$	31	16	-0.027 ± 0.013	16	7
113	$+154 \pm 15$	20	9	$+0.041 \pm 0.038$	20	9	$+0.025 \pm 0.012$	20	7	-0.015 ± 0.026	2	1
114	-9 ± 11	12	22	-0.059 ± 0.024	12	22	-0.012 ± 0.009	12	20	$+0.032 \pm 0.007$	12	15
115	$+9 \pm 10$	8	24	$+0.024 \pm 0.022$	8	23	-0.069 ± 0.026	1	2	-0.087 ± 0.034	1	2
116 ^d	—	—	—	—	—	—	—	—	—	—	—	—
117	-2 ± 7	59	132	$+0.005 \pm 0.021$	55	121	$+0.016 \pm 0.007$	55	90	$+0.028 \pm 0.007$	41	51
118 ^c	$+120 \pm 22$	14	4	-0.018 ± 0.066	13	4	-0.011 ± 0.029	13	3	+0.000	5	0
119	-20 ± 9	17	20	-0.004 ± 0.015	17	19	$+0.003 \pm 0.007$	17	19	-0.013 ± 0.009	17	8
120	-39 ± 7	38	47	$+0.009 \pm 0.011$	34	44	$+0.005 \pm 0.005$	33	25	$+0.059 \pm 0.008$	26	14
121	-66 ± 8	86	177	$+0.038 \pm 0.010$	82	181	$+0.002 \pm 0.008$	82	23	$+0.008 \pm 0.008$	54	17
122	-28 ± 8	41	70	$+0.001 \pm 0.012$	37	58	$+0.020 \pm 0.006$	37	32	$+0.033 \pm 0.008$	31	24
123	-22 ± 8	22	41	$+0.022 \pm 0.017$	17	34	$+0.005 \pm 0.006$	16	26	$+0.010 \pm 0.006$	13	16
124	$+14 \pm 17$	22	49	$+0.020 \pm 0.044$	14	40	$+0.029 \pm 0.012$	14	8	$+0.012 \pm 0.022$	11	1
125	-48 ± 14	57	62	$+0.037 \pm 0.010$	55	64	$+0.007 \pm 0.008$	55	8	-0.018 ± 0.011	55	7
126	$+57 \pm 14$	36	43	-0.004 ± 0.039	33	40	$+0.029 \pm 0.023$	33	9	$+0.008 \pm 0.023$	33	6
127	-3 ± 4	131	187	$+0.002 \pm 0.007$	127	167	-0.007 ± 0.003	127	136	-0.007 ± 0.003	127	83
128	$+6 \pm 12$	24	29	-0.042 ± 0.010	23	29	$+0.010 \pm 0.015$	23	7	-0.055 ± 0.013	9	4
129 ^b	+0	3	0	+0.000	3	0	+0.000	3	0	+0.000	3	0
130 ^d	—	—	—	—	—	—	—	—	—	—	—	—
131 ^e	$+35 \pm 5$	123	174	-0.078 ± 0.014	99	152	$+0.002 \pm 0.004$	98	136	—	—	—
132	-24 ± 19	12	7	-0.025 ± 0.026	12	6	-0.032 ± 0.008	12	4	$+0.005 \pm 0.016$	12	4
133	-11 ± 4	128	247	-0.009 ± 0.006	128	241	-0.009 ± 0.002	128	183	-0.014 ± 0.002	128	151

^a Run 109 is the fiducial run for Mgb' , defined to have zero offset. ^b Runs 104 and 129 have no objects in common with other runs. ^c Run 118 has no Mg_2 measurements in common with other runs. ^d Runs 116 and 130 have no usable data. ^e Run 131 has no Mg_2 measurements.

14 km s⁻¹ in redshift, 6.3 to 5.6 per cent in dispersion, 4.9 to 4.4 per cent in *Mgb* and 0.012 to 0.009 mag in *Mg*₂. We also checked to see whether applying the run corrections reduced the scatter in external comparisons between our data and measurements in the literature (see Section 4.3). We found that, although the scatter is dominated by the combined random errors, the corrections did reduce the scatter slightly in all cases.

As another test of the run corrections for *Mgb'* and *Mg*₂ (and also, more weakly, for log σ), we compared the *Mgb'*- σ and *Mg*₂- σ distributions for each run (after applying the run corrections) with the global *Mgb'*- σ and *Mg*₂- σ relations derived in Paper V. Using the χ^2 goodness-of-fit statistic to account both for measurement errors in the dispersions and linestrengths and for the intrinsic scatter about the *Mg*- σ relations, we find that for *Mgb'*- σ there were two runs (113 and 132) with reduced χ^2 greater than 3, while for *Mg*₂- σ there was one such run (122). In all three cases the removal of one or two obvious outliers decreased the reduced χ^2 to a non-significant level.

3.6 Calibrating the estimated errors

Obtaining precise error estimates is particularly important because we will make extensive use of them in applying maximum likelihood methods to deriving the Fundamental Plane and relative cluster distances for our sample. Although we have estimated the measurement errors as carefully as possible, simulating the noise in the observations and the measurement procedures, some sources of error are likely to remain unaccounted for and we may be systematically mis-estimating the errors. We therefore autocalibrate our errors by scaling the estimated errors in the combined measurements (the *internal* error estimate, based on the individual measurement errors derived from simulations; see Sections 3.3 and 4.2) to match the rms errors from objects with repeat measurements (an *external* error estimate).

Fig. 6 shows the differential and cumulative distributions of the ratio of the rms error to the estimated error for each galaxy with repeat measurements of redshift, dispersion, *Mgb* and *Mg*₂. The smooth curves are the predicted differential and cumulative distributions of this ratio assuming that the estimated errors are the true

errors. The top panel shows the comparison using the estimated errors (including all the corrections discussed above). For the redshifts and linestrengths, the estimated errors are generally underestimates of the true errors, since the ratio of rms and estimated errors tends to be larger than predicted. For the dispersions the estimated errors are generally overestimates of the true errors, since this ratio tends to be smaller than predicted. For all quantities the assumption that the estimated errors are consistent with the true errors is ruled out with high confidence by a Kolmogorov–Smirnov (KS) test applied to the observed and predicted cumulative distributions. These differences between the estimated errors from the simulations and the rms errors from repeat measurements reflect the approximate nature of the S/N estimates and systematic measurement errors not accounted for in the simulations.

In order to bring our estimated errors into line with the rms errors from the repeat measurements, we found it necessary to add 15 km s⁻¹ in quadrature to the estimated redshift errors, to scale the dispersion and *Mgb* errors by factors of 0.85 and 1.15 respectively, and to add 0.005 mag to the *Mg*₂ errors. These corrections were determined by maximizing the agreement of the observed and predicted distributions of the ratio of rms and estimated errors under a KS test (excluding outliers with values of this ratio >3.5). The corrections are quite well determined: to within a couple of km s⁻¹ for the redshift correction, a few per cent for the dispersion and *Mgb* corrections, and 0.001 mag for the *Mg*₂ correction. Applying these corrections, and repeating the comparison of rms and estimated errors, gives the lower panels of Fig. 6, which show the good agreement between the rms errors from repeat measurements and the calibrated error estimates for the redshifts, dispersions and *Mg* linestrengths.

The need for such a correction to the redshift errors may be due in part to the residual zero-point uncertainties in the redshifts, and in part to a tendency for the simulations to underestimate the errors for high S/N spectra. The origin of the overestimation of the dispersion errors is uncertain, although it may result from slightly different prescriptions for estimating the S/N in the observations and the simulations. The underestimation of the linestrength errors may be due to neglecting the effects of errors in the redshift and dispersion

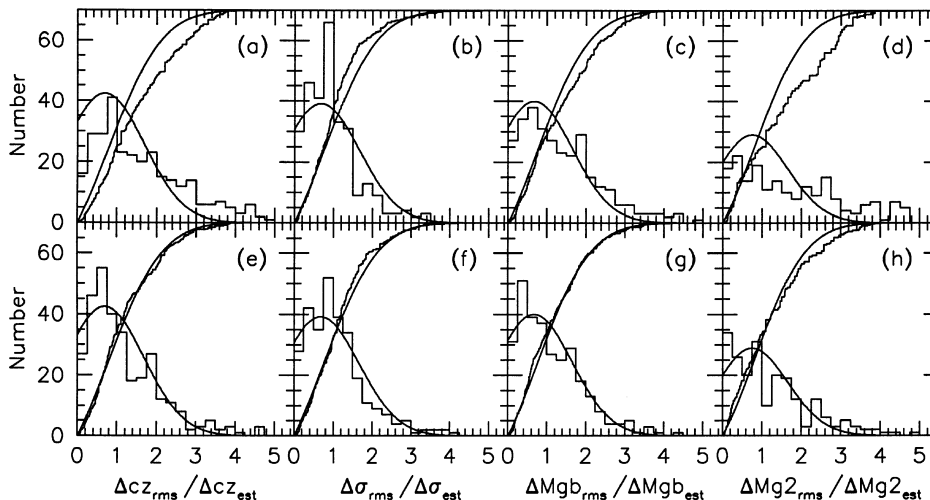


Figure 6. Comparison of the estimated errors derived from simulations and the rms errors for galaxies with repeat measurements of redshift, dispersion, *Mgb* and *Mg*₂. Each panel shows the differential and cumulative distributions of the ratio of rms error to estimated error. The stepped curves are the observed distributions, while the smooth curves are the predicted distributions. The upper panels show the comparisons using the original estimated errors; the lower panels show the comparisons after correcting the estimated errors as described in the text.

Table 3. Individual spectroscopic measurements (abridged). The full table is available electronically.

(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)	(11)	(12)	(13)	(14)	(15)
GINRUNSEQ	Galaxy Name	Telescope	Obsvn Date	Q	S/N	cz (km s ⁻¹)		σ (km s ⁻¹)		Mgb (Å)		Mg ₂ (mag)		Notes
001106215	A76 A	MG24	880914	A	21.4	11296	51	262	50	5.20	0.54	0.353	0.020	
002106221	A76 B	MG24	880914	B	18.8	11317	47	225	47	4.76	0.59	0.326	0.019	
003106218	A76 C	MG24	880914	A	17.8	11973	49	217	48	5.17	0.68	0.392	0.021	
004107095	A76 D	MG24	881011	C	14.8	12184	31	316	103	6.18	0.83	0.326	0.028	
005120504	A76 E	IN25	901016	B	33.2	12241	19	217	16	6.35	0.22	0.312	0.011	
005120505	A76 E	IN25	901017	B	32.8	12147	17	169	11	6.16	0.23	0.275	0.012	
006107098	A76 F	MG24	881011	B	17.6	12371	28	307	82	5.29	0.83	0.331	0.027	
008122534	A85 A	MG24	911016	B	18.7	16577	39	290	52	5.14	0.59	0.344	0.021	
008123319	A85 A	IN25	911201	D	21.4	16692	50	—	—	—	—	—	—	
009120626	A85 B	IN25	901017	A	37.6	17349	34	436	46	5.74	0.22	—	—	
010120628	A85 C	IN25	901017	D	33.9	22837	50	—	—	—	—	—	—	Mgb at sky
011122540	A85 1	MG24	911016	B	25.8	15112	20	195	18	3.82	0.39	0.253	0.016	
012132001	A85 2	SS2B	931022	B	28.3	16264	25	294	32	5.86	0.31	0.337	0.011	
013101059	A119 A	MG24	861202	B	16.5	11516	49	299	48	6.06	0.51	0.294	0.022	
013109339	A119 A	KP4M	881107	A	44.2	11457	19	289	18	5.00	0.20	0.320	0.010	
013131330	A119 A	ES36	931008	A	30.9	11451	24	320	29	4.86	0.26	—	—	
014101063	A119 B	MG24	861202	C	14.4	13205	63	323	65	5.90	0.72	0.320	0.025	
014109339	A119 B	KP4M	881107	A	35.0	13345	21	276	21	4.86	0.22	0.361	0.010	
015109343	A119 C	KP4M	881107	A	33.2	13508	19	250	19	5.51	0.22	0.295	0.011	
015131330	A119 C	ES36	931008	B	19.5	13484	28	265	35	5.88	0.48	—	—	
016109346	A119 D	KP4M	881107	C*	20.8	14980	16	104	13	3.30	0.37	0.151	0.016	H β
016110611	A119 D	KP2M	881114	D*	15.6	15022	50	—	—	—	—	—	—	H β
016131330	A119 D	ES36	931008	D*	16.1	14996	50	—	—	—	—	—	—	H β
017109347	A119 E	KP4M	881107	A	36.1	12807	19	251	17	5.18	0.18	0.326	0.010	
017131330	A119 E	ES36	931008	A	27.0	12788	21	243	22	6.01	0.32	—	—	
018109342	A119 F	KP4M	881107	B	28.4	13034	18	193	15	4.95	0.32	0.245	0.013	
018131330	A119 F	ES36	931008	C	21.1	13006	16	112	20	4.69	0.37	—	—	
019109342	A119 G	KP4M	881107	E	32.7	—	—	—	—	—	—	—	—	mis-ID
019131330	A119 G	ES36	931008	A	36.7	13457	19	267	19	4.78	0.20	—	—	
021122654	A119 I	MG24	911018	A	31.0	13271	20	225	19	4.85	0.38	0.326	0.014	
022131330	A119 J	ES36	931008	B	22.3	13520	21	219	24	4.55	0.33	—	—	
023131330	A119 1	ES36	931008	C*	11.5	4127	18	92	38	0.50	0.78	—	—	H β
024122657	A119 2	MG24	911018	E	28.9	—	—	—	—	—	—	—	—	mis-ID
024131330	A119 2	ES36	931008	B	20.4	12346	16	100	20	5.13	0.38	—	—	
025107103	J3 A	MG24	881011	B	15.2	14453	32	333	111	6.25	0.74	0.329	0.025	
026120714	J3 B	IN25	901018	A	39.5	13520	18	231	15	4.93	0.23	—	—	
027107106	J3 C	MG24	881011	C	15.8	13771	21	163	38	4.33	0.86	0.224	0.027	
028120712	J3 D	IN25	901017	A	36.3	14316	22	287	23	5.45	0.24	—	—	double
028120713	J3 D	IN25	901017	B	30.8	14837	19	207	15	4.87	0.26	—	—	double
031107190	J4 A	MG24	881012	C	13.2	12074	24	185	51	5.35	0.75	0.387	0.029	
032107189	J4 B	MG24	881012	C	13.3	12090	31	302	102	4.11	1.08	0.261	0.032	
033132002	J4 C	SS2B	931021	A	30.0	17154	30	358	43	6.20	0.26	0.355	0.010	
036107195	A147 A	MG24	881012	C	12.2	12811	26	208	62	3.96	1.16	0.288	0.034	
036117298	A147 A	MG24	891015	B	19.0	12741	35	281	49	4.84	0.60	0.314	0.022	
036133157	A147 A	CT4M	931019	A	61.4	12760	17	235	12	4.77	0.17	0.285	0.010	
036133158	A147 A	CT4M	931019	A	62.3	12776	17	253	13	5.01	0.18	0.289	0.010	

The columns of this table give: (1) observation identifier (GINRUNSEQ); (2) galaxy name; (3) telescope used; (4) date of observation; (5) quality parameter; (6) signal-to-noise ratio; (7)–(8) redshift and estimated error; (9)–(10) velocity dispersion and estimated error; (11)–(12) Mgb linestrength and estimated error; (13)–(14) Mg₂ linestrength and estimated error; and (15) notes on each observation. In the notes, ‘double’ means that the EFAR galaxy is double; ‘star’ means that the EFAR object is a star, not a galaxy; ‘mis-ID’ means that the spectrum is for some galaxy other than the nominated EFAR object; ‘mis-ID*’ means that the spectrum is for a nearby star rather than the EFAR object; ‘Mgb at sky’ means that the object is at a redshift that puts Mgb on the 5577-Å sky line; ‘number = number’ notes the duplicated pairs in the EFAR sample (see Paper I; only the first of the two GINs is used); emission-line objects (with an asterisk on Q) have the emission features listed; ‘H β abs’ or ‘H β abs, [OIII]’ means that the redshift is based on the H β absorption feature (and [OIII] if present), as the spectrum stops short of Mgb (no dispersion or Mgb index is given for these objects). The objects for which we have no spectrum have GINs 7, 20, 29, 30, 34, 35, 55, 62, 64, 67, 82, 83, 91, 104, 121, 131, 133, 134, 161, 181, 191, 214, 225, 228, 231, 234, 256, 265, 309, 327, 391, 405, 407, 417, 434, 435, 442, 450, 451, 452, 458, 463, 464, 465, 470, 475, 477, 483, 484, 486, 494, 516, 520, 521, 522, 523, 526, 544, 551, 553, 567, 569, 570, 575, 576, 577, 587, 594, 597, 603, 605, 624, 625, 644, 671, 727, 760, 793, 798, 801, 901. This is an abridged version of this table; the full table is available on Synergy, from NASA’s Astrophysical Data Center (ADC); (URL <http://adc.gsfc.nasa.gov>, table ID number J/MNRAS/305/259) and the Centre de Données astronomiques de Strasbourg (CDS); (URL <http://cdsweb.u-strasbg.fr>, table ID number J/MNRAS/305/259).

estimates and the different continuum shapes of spectra from different runs when measuring linestrengths.

4 RESULTS

4.1 Individual measurements

The previous two sections describe the observations and analysis of our spectroscopic data. Table 3 (full version available electronically) lists the observational details for each spectrum and the fully corrected measurements of redshift, dispersion, *Mgb* and *Mg₂*, together with their calibrated error estimates. Note that these error estimates are the individual measurement errors, and must be combined in quadrature with the run correction uncertainties given in Table 2 to give the total error estimate. We list the measurement errors rather than the total errors because the total errors are not independent, being correlated for objects in the same run. The version of the table presented here is abridged; the full table is available from NASA’s Astrophysical Data Center (ADC) and from the Centre de Données astronomiques de Strasbourg (CDS).

The entries in Table 3 are as follows: column 1 gives GINRUN-SEQ, a unique nine-digit identifier for each spectrum, composed of the galaxy identification number (GIN) as given in the master list of EFAR sample galaxies (table 3 of Paper I), the run number (RUN) as given in Table 1, and a sequence number (SEQ) which uniquely specifies the observation within the run; column 2 gives the galaxy name, as in the master list of Paper I; column 3 is the telescope code, as in Table 1; column 4 is the UT date of the observation; columns 5

and 6 are the quality parameter (with an asterisk if the spectrum shows emission features) and S/N of the spectrum (see Section 2.4); columns 7 and 8 are the fully corrected heliocentric redshift cz (in km s^{-1}) and its measurement error; columns 9 and 10 are the fully corrected velocity dispersion σ (in km s^{-1}) and its measurement error; columns 11 and 12 are the fully corrected *Mgb* linestrength index and its measurement error (in \AA); columns 13 and 14 are the fully corrected *Mg₂* linestrength index and its measurement error (in mag); column 15 provides comments, the full meanings of which are described in the notes to the table.

There are 1319 spectra in this table. Note that 81 objects from our original sample do not have spectroscopic observations and do not appear in the table (see the list of missing GINs in the table notes). Three of these are the duplicate objects (GINs 55, 435, 476) and three are known stars (GINs 131, 133, 191). Most of the others are objects that our imaging showed are not early-type galaxies, although there are a few early-type galaxies for which we did not get a spectrum. There are 34 spectra that are unusable ($Q = E$) either because the spectrum is too poor (13 cases) or because the object was mis-identified (20 cases) or is a known star (one case, GIN 123). Of the 1285 usable spectra (for 706 different galaxies), there are 637 spectra with $Q = A$, 407 with $Q = B$, 161 with $Q = C$ and 80 with $Q = D$.

The cumulative distributions of the total estimated errors in the individual measurements (combining measurement errors and run correction uncertainties in quadrature) are shown in Fig. 7 for quality classes A, B and C, and for all three classes together. The error distributions can be quantitatively characterized by their 50 and 90 per cent points, which are listed in Table 4. The overall median error in a single redshift measurement is 22 km s^{-1} , the median relative errors in single measurements of dispersion and *Mgb* are 10.5 and 8.2 per cent, and the median error in a single measurement of *Mg₂* is 0.015 mag.

4.2 Combining measurements

We use a weighting scheme to combine the individual measurements of each quantity to obtain a best estimate (and its uncertainty) for each galaxy in our sample. The weighting has three components.

(i) *Error weighting.* For multiple measurements X_i having estimated total errors Δ_i (the measurement errors and run correction uncertainties added in quadrature), we weight the values inversely with their variances, i.e. by Δ_i^{-2} .

(ii) *Quality weighting.* We apply a weighting W_Q which quantifies our degree of belief (over and above the estimated errors) in measurements obtained from spectra with different quality parameters. Following the discussion in Section 2.4, for spectra with $Q = A, B, C, D, E$ we use $W_Q = 1, 1, 1, 0.5, 0$ in computing redshifts, $W_Q = 1, 1, 0.5, 0, 0$ in computing dispersions, and $W_Q = 1, 1, 0.5, 0, 0$ in computing linestrengths.

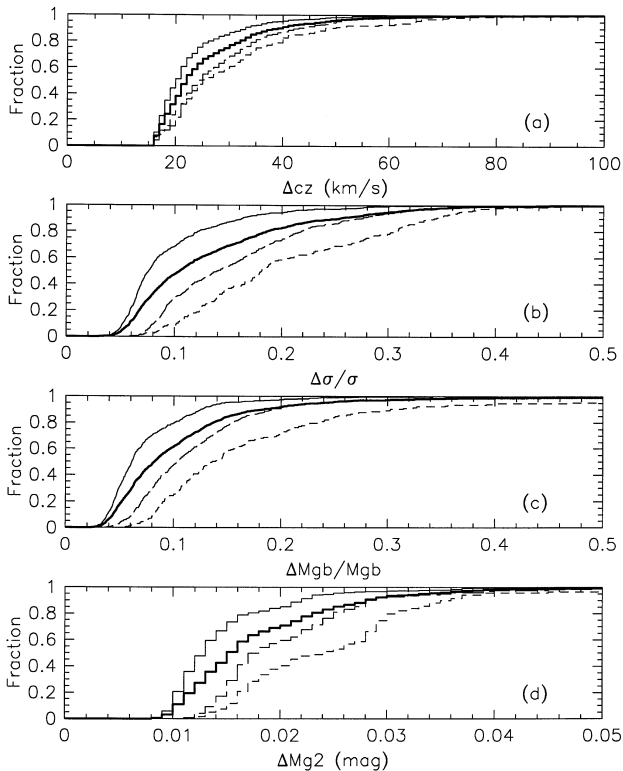


Figure 7. The cumulative distributions of estimated errors for individual measurements of redshift, velocity dispersion and *Mgb* linestrength. The distributions for quality classes A, B and C are shown as full, long-dashed and short-dashed lines respectively; the overall distribution is the thick full line. (a) The distribution of estimated errors in cz ; (b) estimated relative errors in σ ; (c) estimated relative errors in *Mgb*; (d) estimated errors in *Mg₂*.

Table 4. The distribution of estimated errors per measurement.

Q	Δcz (km s^{-1})		$\Delta\sigma/\sigma$		$\Delta Mgb/Mgb$		ΔMg_2 (mag)	
	50%	90%	50%	90%	50%	90%	50%	90%
	All	22	40	0.105	0.255	0.082	0.184	0.015
A	20	33	0.076	0.163	0.061	0.125	0.013	0.022
B	24	43	0.140	0.275	0.104	0.192	0.018	0.028
C	25	48	0.181	0.343	0.136	0.303	0.024	0.036

(iii) *Run weighting.* We also apply a run weighting $W_R = 0$ to exclude run 115, for reasons explained in Section 2.1; all other runs are given $W_R = 1$.

The combined estimate X is thus computed from the individual measurements X_i as the weighted mean

$$X = \frac{\sum_i W_i X_i}{\sum_i W_i}, \quad (10)$$

where $W_i = \Delta_i^{-2} W_{Qi} W_{Ri}$. The uncertainty in this weighted mean is computed as

$$\Delta = \left(\sum_i W_i \right)^{-1/2}. \quad (11)$$

This procedure is used to obtain combined estimates of the redshift, dispersion and linestrengths for each galaxy. We estimate the overall quality Q as the highest quality amongst the individual measurements, and obtain a combined estimate of the S/N as

$$S/N = \left[\sum_i (S/N)_i^2 W_{Qi} W_{Ri} \right]^{1/2}, \quad (12)$$

using the same weightings as for the dispersions (except when the overall quality is $Q = D$, when these weightings are omitted).

Table 5 gives the combined estimates of the spectroscopic parameters for each galaxy in the EFAR sample. The table lists galaxy identification number (GIN), galaxy name, cluster assignment number (CAN; see Section 5), and the number of spectra, redshifts, dispersions and Mgb and Mg_2 linestrengths obtained for this object; then, for each of redshift, dispersion, Mgb and Mg_2 , the combined estimate, its estimated total error (Δ) and the weighted rms error from any repeat observations (δ); finally, the combined S/N estimate and the overall quality parameter (with an asterisk if the galaxy possesses emission lines). Note that only objects with useful measurements are included; hence the lowest quality class present in this table is $Q = D$, and the seven galaxies with only $Q = E$ spectra (GINs 123, 284, 389, 448, 599, 637 and 679) in Table 3 are omitted.

The cumulative distributions of uncertainties in the combined results are shown in Fig. 8, both for the entire data set and for quality classes A, B and C separately. The error distributions can be quantitatively characterized by their 50 and 90 per cent points, which are listed in Table 6. The overall median error in redshift is 20 km s^{-1} , the median relative errors in dispersion and Mgb are 9.1 and 7.2 per cent, and the median error in Mg_2 is 0.015 mag. For the whole sample, and for quality classes A and B, the median errors in the combined measurements are smaller than the median errors in the individual measurements, as one expects. However, for dispersion, Mgb and Mg_2 the errors are larger for quality class C and at the 90th percentile; this results from assigning a quality weighting of 0.5 to $Q = C$ when combining the individual measurements of these quantities.

The uncertainties listed in Table 5 represent the best estimates of the total errors in the parameters for each galaxy. However, it must be emphasized that they are *not* independent of each other, as the run correction errors are correlated across all measurements from a run. To simulate the joint distribution of some parameter properly for the whole data set, one must first generate realisations of the run correction errors (drawn from Gaussians with standard deviations given by the uncertainties listed in Table 2) and the individual measurement errors (drawn from Gaussians with standard deviations given by the uncertainties listed Table 3). For each individual measurement, one must add the realization of its measurement error and the realization of the appropriate run correction error (the same for all measurements in a given run) to the measured value of the parameter. The resulting realizations of the individual

measurements are finally combined using the recipe described above to yield a realization of the value of the parameter for each galaxy in the data set.

The distributions of redshift, velocity dispersion, Mgb and Mg_2 for the galaxies in the EFAR sample are displayed in Fig. 9. The galaxies for which we measured velocity dispersions are only a subset of our sample of programme galaxies (629/743), and represent a refinement of the sample selection criteria. Fig. 10 shows the fraction of programme galaxies with measured dispersions as a function of the galaxy diameter D_W on which the selection function of the programme galaxy sample is defined. There is a steady decline in the fraction of the sample for which usable dispersions were measured, from 100 per cent for the largest galaxies (with $D_W \geq 40$ arcsec) to about 75 per cent for the smallest (with $8 \leq D_W \leq 15$ arcsec; there are only three programme galaxies with $D_W < 8$ arcsec). This additional selection effect must be allowed for when determining Fundamental Plane distances.

4.3 Internal and external comparisons

One of the strengths of our spectroscopic sample is the high fraction of objects with repeat observations: there are 375 galaxies with a single dispersion measurement, 160 with two measurements and 141 with three or more measurements. Fig. 11 shows the cumulative distributions of rms errors in redshift, dispersion, Mgb and Mg_2 obtained from these repeat observations. The detailed internal comparisons made possible by these repeat measurements have been used to establish the run corrections (Section 3.5) and to calibrate the estimated errors (Section 3.6). The latter process ensured that the estimated errors were statistically consistent with the rms errors of the repeat measurements.

We also make external comparisons of our measurements with the work of other authors. The EFAR redshifts are compared in Fig. 12 with redshifts given in the literature by the 7 Samurai (Davies et al. 1987), Dressler & Shectman (1988), Beers et al. (1991), Malumuth et al. (1992), Zabludoff et al. (1993), Colless & Dunn (1996) and Lucey et al. (1997). Only 11 of the 256 comparisons give redshift differences greater than 300 km s^{-1} : in six cases the EFAR redshift is confirmed either by repeat measurements or other published measurements; in the remaining five cases the identification of the galaxy in question is uncertain in the literature. For the 245 cases in which the redshift difference is less than 300 km s^{-1} , there is no significant velocity zero-point error and the rms scatter is 85 km s^{-1} . Since our repeat measurements show much smaller errors (90 per cent are less than 36 km s^{-1}), most of this scatter must arise in the literature data, some of which were taken at lower resolution or S/N than our data.

Fig. 13 compares the EFAR dispersions with published dispersions from the work of the 7 Samurai (Davies et al. 1987), Guzmán (1993), Jørgensen (1997) and Lucey et al. (1997), and the compilation of earlier measurements by Whitmore, McElroy & Tonry (1985). Note that we do not compare with the more recent compilation by McElroy (1995), since its overlap with our sample is essentially just the sum of above sources. The mean differences, $\Delta = \log \sigma_{\text{EFAR}} - \log \sigma_{\text{lit}}$, and their standard errors are indicated in the figure; none of these scale differences is larger than 6 per cent and in fact all five comparisons are consistent with zero scale error at the 2σ level or better. The rms scatter in these comparisons is significantly greater than the errors in our dispersion measurements, implying that in general the literature measurements have larger errors and/or that there are unaccounted for uncertainties in the comparison.

Table 5. Spectroscopic parameters for the EFAR galaxies.

(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)	(11)	(12)	(13)	(14)	(15)	(16)	(17)	(18)
GIN	Galaxy Name	CAN	N $s z \sigma b 2$	cz (km s^{-1})	Δcz	δcz	σ (km s^{-1})	$\Delta \sigma$	$\delta \sigma$	Mgb	ΔM_{gb} (\AA)	δM_{gb}	M_{g_2}	ΔM_{g_2} (mag)	δM_{g_2}	S/N	Q
1	A76 A	1	1 1 1 1 1	11296	53	—	262	53	—	5.20	0.57	—	0.353	0.024	—	21.4	A
2	A76 B	1	1 1 1 1 1	11317	49	—	224	49	—	4.76	0.62	—	0.326	0.024	—	18.8	B
3	A76 C	1	1 1 1 1 1	11973	51	—	217	50	—	5.17	0.70	—	0.392	0.025	—	17.8	A
4	A76 D	1	1 1 1 1 1	12184	33	—	316	150	—	6.18	1.21	—	0.326	0.041	—	10.5	C
5	A76 E	1	2 2 2 2 2	12189	14	47	184	10	23	6.26	0.18	0.09	0.295	0.010	0.018	46.7	B
6	A76 F	1	1 1 1 1 1	12371	30	—	307	85	—	5.29	0.86	—	0.331	0.028	—	17.6	B
8	A85 A	2	2 2 1 1 1	16604	35	49	290	53	—	5.14	0.61	—	0.344	0.022	—	18.7	B
9	A85 B	2	1 1 1 1 0	17349	35	—	436	47	—	5.74	0.25	—	—	—	—	37.6	A
10	A85 C	102	1 1 0 0 0	22837	71	—	—	—	—	—	—	—	—	—	—	33.9	D
11	A85 1	2	1 1 1 1 1	15112	22	—	195	19	—	3.82	0.42	—	0.253	0.018	—	25.8	B
12	A85 2	2	1 1 1 1 1	16264	31	—	294	37	—	5.86	0.37	—	0.337	0.019	—	28.3	B
13	A119 A	3	3 3 3 3 2	11459	15	17	297	15	13	5.04	0.16	0.31	0.316	0.009	0.010	56.4	A
14	A119 B	3	2 2 2 2 2	13330	20	42	278	21	10	4.90	0.22	0.21	0.358	0.010	0.011	36.5	A
15	A119 C	3	2 2 2 2 1	13500	16	11	253	17	6	5.57	0.20	0.14	0.295	0.011	—	38.5	A
16	A119 D	3	3 3 1 1 1	14982	16	9	104	18	—	3.30	0.52	—	0.151	0.023	—	14.7	C*
17	A119 E	3	2 2 2 2 1	12798	14	9	248	14	4	5.37	0.16	0.35	0.326	0.010	—	45.1	A
18	A119 F	3	2 2 2 2 1	13018	12	14	175	13	33	4.88	0.28	0.11	0.245	0.013	—	32.1	B
19	A119 G	3	2 1 1 1 0	13457	20	—	267	21	—	4.78	0.22	—	—	—	—	36.7	A
21	A119 I	3	1 1 1 1 1	13271	22	—	225	20	—	4.85	0.41	—	0.326	0.016	—	31.0	A
22	A119 J	3	1 1 1 1 0	13520	22	—	219	25	—	4.55	0.35	—	—	—	—	22.3	B
23	A119 1	103	1 1 1 1 0	4127	19	—	92	54	—	0.50	1.12	—	—	—	—	8.1	C*
24	A119 2	3	2 1 1 1 0	12346	17	—	100	20	—	5.13	0.39	—	—	—	—	20.4	B
25	J3 A	4	1 1 1 1 1	14453	34	—	333	114	—	6.25	0.77	—	0.329	0.026	—	15.2	B
26	J3 B	4	1 1 1 1 0	13519	19	—	231	16	—	4.93	0.26	—	—	—	—	39.5	A
27	J3 C	4	1 1 1 1 1	13770	23	—	163	57	—	4.33	1.26	—	0.224	0.040	—	11.2	C
28	J3 D	4	2 2 2 2 0	14610	15	258	231	13	37	5.18	0.20	0.29	—	—	—	47.6	A
31	J4 A	5	1 1 1 1 1	12074	26	—	185	75	—	5.35	1.11	—	0.387	0.043	—	9.3	C
32	J4 B	5	1 1 1 1 1	12090	33	—	302	148	—	4.11	1.56	—	0.261	0.047	—	9.4	C
33	J4 C	105	1 1 1 1 1	17154	36	—	358	48	—	6.20	0.32	—	0.355	0.019	—	30.0	A
36	A147 A	6	4 4 4 4 4	12771	11	19	244	9	12	4.88	0.13	0.14	0.289	0.007	0.008	89.9	A
37	A147 B	6	4 4 4 4 4	13119	11	15	316	10	20	4.68	0.11	0.30	0.304	0.006	0.017	104.2	A
38	A147 C	6	3 3 3 3 3	13156	13	8	247	12	7	5.22	0.15	0.22	0.305	0.008	0.007	66.8	A
39	A147 D	6	3 3 3 3 3	13444	12	2	185	9	15	4.98	0.18	0.28	0.294	0.008	0.004	67.8	A
40	A147 E	6	3 3 3 3 3	13049	10	10	176	8	7	4.58	0.14	0.12	0.267	0.007	0.002	76.1	A
41	A147 F	6	2 2 2 2 2	11922	12	5	87	10	4	3.65	0.27	0.02	0.195	0.012	0.002	37.9	B
42	A147 1	6	3 3 3 3 3	12832	11	29	148	9	3	4.35	0.16	0.09	0.252	0.008	0.013	65.8	A
43	A160 A	7	2 2 2 2 1	11401	15	3	181	17	19	3.86	0.26	0.04	0.250	0.017	—	35.1	A
44	A160 B	107	3 3 1 1 1	18258	24	17	192	21	—	6.61	0.30	—	0.289	0.022	—	26.3	B
45	A160 C	7	1 1 1 1 0	12380	34	—	412	58	—	4.61	0.37	—	—	—	—	27.8	A
46	A160 D	107	3 3 0 0 0	18271	41	84	—	—	—	—	—	—	—	—	—	36.1	D
47	A160 E	7	4 4 4 4 2	14056	13	24	226	16	12	5.01	0.23	0.26	0.293	0.017	0.014	39.2	A
48	A160 F	7	3 3 3 3 2	13657	14	16	176	18	27	5.18	0.24	0.43	0.307	0.013	0.009	33.9	A
49	A160 G	7	4 4 4 4 2	13137	13	25	196	20	22	4.82	0.25	0.39	0.293	0.015	0.056	32.4	B
50	A160 H	7	1 1 1 1 0	13589	21	—	195	23	—	5.18	0.30	—	—	—	—	21.6	A
51	A160 I	107	3 3 0 0 0	18643	41	30	—	—	—	—	—	—	—	—	—	43.9	D
52	A160 J	7	4 4 4 4 1	11254	9	13	145	10	15	2.74	0.20	0.32	0.217	0.041	—	46.1	A*
53	A160 1	107	2 2 0 0 0	18108	50	103	—	—	—	—	—	—	—	—	—	27.3	D*
54	A160 2	107	1 1 0 0 0	18201	71	—	—	—	—	—	—	—	—	—	—	21.7	D
56	A168 A	108	1 1 1 1 0	5299	19	—	265	16	—	4.35	0.26	—	—	—	—	51.7	A
57	A168 B	108	1 1 1 1 0	5253	23	—	310	25	—	5.29	0.30	—	—	—	—	43.6	A
58	A168 C	8	1 1 1 1 1	13313	22	—	216	22	—	4.79	0.40	—	0.273	0.016	—	27.1	A
59	A168 D	8	1 1 1 1 1	13545	48	—	345	58	—	5.30	0.50	—	0.396	0.021	—	21.7	B
60	A168 E	8	1 1 1 1 1	13402	44	—	340	54	—	4.93	0.48	—	0.333	0.020	—	23.6	B
61	A168 F	8	1 1 1 1 1	13582	33	—	307	41	—	5.32	0.50	—	0.345	0.017	—	26.3	A
63	A168 1	8	1 1 1 1 1	13488	25	—	230	26	—	3.76	0.47	—	0.294	0.020	—	23.5	B
65	A189 A	9	1 1 1 1 1	5588	24	—	250	24	—	4.55	0.32	—	0.271	0.015	—	27.2	A
66	A189 B	9	1 1 1 1 1	4837	25	—	181	15	—	5.61	0.27	—	0.315	0.019	—	35.6	A
68	A189 D	109	1 1 1 1 1	9599	29	—	296	31	—	5.93	0.30	—	0.364	0.019	—	35.7	A
69	J30 A	10	2 2 2 2 2	15390	16	3	303	12	26	5.77	0.14	0.07	0.328	0.008	0.001	77.7	A
70	J30 B	10	2 2 2 2 2	14807	12	8	206	9	2	4.95	0.15	0.22	0.293	0.008	0.011	65.2	A
71	J30 C	10	2 2 2 2 2	14704	13	9	204	9	20	5.17	0.18	0.09	0.303	0.009	0.007	65.6	A
72	J30 D	10	2 2 2 2 2	16451	13	19	234	9	4	5.50	0.14	0.06	0.313	0.008	0.001	73.4	A

Table 5 – continued

(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)	(11)	(12)	(13)	(14)	(15)	(16)	(17)	(18)
GIN	Galaxy Name	CAN	N $s z \sigma b 2$	cz (km s^{-1})	Δcz	δcz	σ (km s^{-1})	$\Delta \sigma$	$\delta \sigma$	Mgb	ΔM_{gb} (\AA)	δM_{gb}	M_{g_2}	ΔM_{g_2} (mag)	δM_{g_2}	S/N	Q
73	J30 E	10	2 2 2 2 2	15253	13	3	223	9	0	4.81	0.18	0.03	0.284	0.008	0.000	66.4	A
74	J30 F	210	2 2 2 2 2	20957	12	6	234	7	9	5.43	0.14	0.12	0.317	0.007	0.001	84.4	A
75	J30 G	110	2 2 0 0 0	12206	50	4	—	—	—	—	—	—	—	—	—	52.4	D*
76	J30 1	10	2 2 2 2 2	15752	13	1	202	8	10	4.83	0.17	0.01	0.317	0.008	0.003	71.0	A
77	J30 2	10	2 2 2 2 2	16463	11	6	58	8	8	3.16	0.30	0.11	0.211	0.011	0.010	43.0	B
78	A193 A	11	1 1 1 1 1	14740	68	—	629	132	—	2.24	0.27	—	0.216	0.014	—	34.4	A
79	A193 B	11	1 1 1 1 1	14673	25	—	187	16	—	5.86	0.32	—	0.348	0.019	—	32.3	B
80	A193 C	11	1 1 1 1 1	14395	28	—	221	26	—	5.45	0.39	—	0.348	0.020	—	25.6	B
81	A193 D	11	1 1 1 1 1	15181	25	—	254	29	—	4.02	0.50	—	0.281	0.017	—	26.5	B
84	J32 A	12	3 3 3 3 3	12367	12	6	251	9	11	5.18	0.15	0.33	0.339	0.008	0.009	72.3	A
85	J32 B	12	3 3 3 3 3	11647	11	7	273	8	6	4.99	0.13	0.27	0.309	0.007	0.012	95.6	A
86	J32 C	112	3 3 0 0 0	20992	41	62	—	—	—	—	—	—	—	—	—	73.9	D*
87	J32 D	12	2 2 2 2 2	12252	13	1	196	8	11	5.24	0.17	0.17	0.329	0.008	0.011	70.6	A
88	J32 E	12	2 2 2 2 2	12345	13	3	247	8	3	5.03	0.15	0.22	0.322	0.008	0.005	81.3	A
89	J32 F	12	2 2 2 2 2	12192	11	22	116	8	1	2.75	0.25	0.15	0.124	0.009	0.000	50.7	A
90	J32 G	12	2 2 2 2 2	13338	11	1	101	7	3	3.82	0.20	0.10	0.240	0.010	0.002	57.1	A
92	J32 1	12	3 3 2 2 2	12694	12	5	109	11	20	3.78	0.26	0.19	0.182	0.011	0.011	36.7	A
93	J32 2	12	2 2 2 2 2	12251	12	6	202	7	4	4.84	0.17	0.47	0.291	0.008	0.001	74.7	A
94	J32 3	12	2 2 2 2 2	11938	12	10	144	8	19	4.54	0.19	0.03	0.301	0.009	0.015	60.0	A
95	J32 4	12	2 2 2 2 2	12398	12	8	127	8	9	3.83	0.24	0.10	0.209	0.009	0.002	51.1	A
96	A260 A	13	1 1 1 1 1	10651	33	—	346	106	—	5.90	0.73	—	0.415	0.024	—	17.8	B
97	A260 B	13	2 2 1 1 1	11391	19	4	182	14	—	5.64	0.31	—	0.335	0.015	—	30.3	A
98	A260 C	13	1 1 1 1 1	10296	30	—	295	122	—	4.09	1.23	—	0.305	0.034	—	11.5	C
99	A260 D	13	1 1 1 1 1	11721	17	—	139	11	—	4.36	0.29	—	0.215	0.015	—	28.6	A
100	A260 E	13	1 1 1 1 1	11789	31	—	283	88	—	6.51	0.77	—	0.368	0.033	—	14.5	B
101	A260 F	13	1 1 1 1 1	10613	23	—	154	60	—	4.50	1.41	—	0.282	0.037	—	10.1	C
102	A260 G	13	1 1 1 1 1	10557	20	—	257	18	—	5.22	0.25	—	0.319	0.013	—	40.2	A
103	A260 H	13	1 1 1 1 1	11573	19	—	223	16	—	5.14	0.25	—	0.304	0.014	—	37.6	A
105	A262 A	14	3 3 3 3 2	4845	16	5	194	23	18	6.71	0.41	0.33	0.392	0.020	0.000	23.3	B
106	A262 B	14	2 2 2 2 1	5227	18	15	216	17	3	4.53	0.27	0.05	0.315	0.015	—	30.4	B
107	A262 C	14	2 2 2 2 1	5086	15	10	224	11	1	4.36	0.21	0.19	0.317	0.012	—	50.4	A
108	A262 D	14	2 2 2 2 1	5036	15	3	229	11	0	4.38	0.18	0.04	0.328	0.011	—	51.0	A
109	A262 E	14	1 1 1 1 1	5346	18	—	178	11	—	4.12	0.26	—	0.254	0.012	—	35.7	A
110	A262 F	14	1 1 1 1 1	5330	20	—	221	16	—	4.40	0.28	—	0.301	0.014	—	31.1	A
111	A262 G	14	2 2 2 2 1	4644	17	7	241	14	17	4.43	0.22	0.08	0.255	0.011	—	43.3	A
112	A262 H	14	1 1 1 1 1	4808	17	—	101	10	—	3.71	0.32	—	0.206	0.014	—	26.4	B
113	A262 I	14	1 1 1 1 1	5584	22	—	180	19	—	4.97	0.44	—	0.293	0.016	—	19.4	B
114	A262 1	14	1 1 1 1 0	5542	18	—	102	12	—	1.27	0.34	—	—	—	—	34.3	A
115	J7 A	15	1 1 1 1 1	10801	50	—	231	49	—	4.29	0.63	—	0.289	0.025	—	19.0	B
116	J7 B	15	1 1 1 1 1	10009	61	—	285	61	—	4.91	0.48	—	0.307	0.023	—	21.4	A
117	J7 C	15	1 1 1 1 1	10683	41	—	201	40	—	3.85	0.56	—	0.286	0.024	—	19.9	B
118	J7 D	15	1 1 1 1 1	10538	50	—	183	54	—	4.85	0.81	—	0.315	0.028	—	14.5	B
119	J7 E	15	1 1 1 1 1	10765	23	—	295	24	—	5.71	0.27	—	0.333	0.014	—	36.9	A
120	J7 F	15	1 1 1 1 1	11553	17	—	161	10	—	4.89	0.26	—	0.300	0.014	—	37.7	A
122	J8 A	16	2 2 2 2 1	8528	27	8	294	29	25	4.75	0.33	0.55	0.306	0.015	—	32.3	A
124	J8 C	16	3 3 3 3 3	9430	23	15	246	26	23	4.62	0.31	0.45	0.321	0.012	0.006	30.0	B
125	J8 D	16	3 3 3 3 3	9492	15	2	108	17	25	4.37	0.47	0.48	0.310	0.017	0.031	18.7	C
126	J8 E	16	2 2 2 2 2	9610	25	7	373	27	11	4.76	0.24	0.04	0.343	0.010	0.013	48.7	A
127	J8 F	16	1 1 1 1 1	10208	20	—	198	16	—	4.89	0.33	—	0.312	0.015	—	26.5	A
128	J8 G	16	2 2 2 2 1	9043	14	4	176	12	2	4.57	0.26	0.03	0.265	0.013	—	35.9	A
129	J8 H	16	1 1 1 1 1	9072	16	—	71	9	—	3.23	0.36	—	0.143	0.014	—	23.3	B
130	J8 I	16	1 1 1 1 1	9218	20	—	176	15	—	4.75	0.30	—	0.288	0.014	—	24.8	B
132	J8 K	16	1 1 1 1 1	9782	21	—	196	18	—	4.55	0.48	—	0.283	0.016	—	23.3	B
135	A376 A	17	1 1 1 1 1	14549	38	—	254	51	—	6.08	0.64	—	0.440	0.022	—	17.2	B
136	A376 B	17	1 1 1 1 1	13619	32	—	212	39	—	4.53	0.81	—	0.248	0.022	—	17.0	B
137	A376 C	17	1 1 1 1 1	14142	49	—	304	73	—	4.40	0.62	—	0.306	0.022	—	16.7	B
138	A376 D	17	1 1 1 1 0	14826	22	—	212	20	—	5.48	0.35	—	—	—	—	25.8	B
139	A376 E	17	2 2 1 1 1	15003	21	5	205	22	—	4.45	0.47	—	0.266	0.018	—	23.9	B
140	A376 F	17	2 2 2 2 1	13955	17	10	166	25	23	3.33	0.49	0.49	0.267	0.034	—	18.3	C
141	A376 G	17	1 1 1 1 0	14823	23	—	236	23	—	5.95	0.31	—	—	—	—	25.7	B
142	J9 A	18	3 3 2 2 2	7763	17	18	276	17	41	4.99	0.22	0.24	0.341	0.009	0.001	48.9	A
143	J9 B	18	3 2 2 2 2	8754	19	3	232	18	17	5.30	0.27	0.62	0.328	0.012	0.011	31.0	A

Table 5 – *continued*

(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)	(11)	(12)	(13)	(14)	(15)	(16)	(17)	(18)
GIN	Galaxy Name	CAN	N $s z \sigma b 2$	cz (km s^{-1})	Δcz	δcz	σ (km s^{-1})	$\Delta \sigma$	$\delta \sigma$	Mgb	ΔMgb (\AA)	δMgb	Mg_2	ΔMg_2 (mag)	δMg_2	S/N	Q
144	J9 C	18	2 2 2 2 2	8666	21	13	328	24	18	4.92	0.22	0.08	0.341	0.009	0.018	45.1	A
145	J9 D	18	1 1 1 1 1	8960	21	—	157	18	—	3.93	0.42	—	0.278	0.015	—	18.6	B
146	J9 E	18	1 1 1 1 1	8197	20	—	187	16	—	4.36	0.32	—	0.281	0.015	—	25.3	A
147	J9 F	18	1 1 1 1 1	7927	18	—	166	11	—	4.60	0.28	—	0.279	0.012	—	32.5	A
148	J9 G	18	1 1 0 0 0	8207	71	—	—	—	—	—	—	—	—	—	—	15.4	D
149	J33 A	19	2 1 1 1 1	8745	24	—	266	27	—	5.62	0.34	—	0.301	0.016	—	31.0	A
150	J33 B	19	1 1 1 1 1	9767	20	—	162	15	—	4.48	0.34	—	0.257	0.013	—	27.8	B
151	A397 A	20	3 3 3 3 3	10251	13	14	215	10	11	4.77	0.18	0.18	0.296	0.008	0.016	57.9	A
152	A397 B	20	4 4 4 4 4	10371	12	15	289	11	22	5.37	0.14	0.22	0.318	0.006	0.015	84.6	A
153	A397 C	20	4 4 4 4 4	8970	11	9	202	9	11	4.30	0.18	0.39	0.243	0.008	0.011	63.6	A
154	A397 D	20	4 4 4 4 4	9361	12	19	203	10	12	5.08	0.20	0.40	0.276	0.008	0.008	57.6	A
155	A397 E	20	3 2 2 2 2	10017	12	2	137	9	9	4.45	0.19	0.03	0.246	0.009	0.005	49.1	A
156	A397 F	20	3 3 2 2 2	9418	15	39	246	13	17	4.87	0.23	0.05	0.256	0.009	0.004	50.3	A
157	A397 G	20	3 3 3 3 3	9462	12	19	201	10	15	4.43	0.17	0.07	0.273	0.009	0.003	65.0	A
158	A397 H	120	3 3 3 3 2	16841	15	17	278	14	19	5.36	0.21	0.33	0.337	0.009	0.002	57.1	A
159	A397 I	20	3 3 3 3 3	9561	11	26	134	9	7	4.12	0.26	0.35	0.234	0.009	0.013	48.9	A
160	A400 A	21	2 2 2 2 1	6826	18	53	307	17	17	4.83	0.18	0.10	0.342	0.013	—	56.0	A
162	A400 C	21	2 2 2 2 1	7559	16	15	222	15	13	4.06	0.23	0.32	0.290	0.014	—	39.2	A
163	A400 D	21	3 3 3 3 1	6758	12	7	185	10	7	4.45	0.18	0.20	0.283	0.014	—	44.3	A
164	A400 E	21	3 3 3 3 1	6748	14	10	279	15	13	4.93	0.18	0.21	0.350	0.013	—	55.3	A
165	A400 F	21	3 3 3 3 1	7385	11	9	128	9	7	3.72	0.27	0.35	0.215	0.015	—	35.2	A
166	A400 G	21	3 3 3 3 1	7225	11	10	153	9	5	4.74	0.22	0.24	0.285	0.014	—	40.8	A
167	A400 H	21	2 2 2 2 1	8205	12	6	103	14	23	3.41	0.40	0.05	0.133	0.027	—	20.7	C*
168	A400 I	21	2 2 2 2 0	6665	13	24	133	26	25	2.78	0.56	0.05	—	—	—	16.4	C*
169	A400 J	21	3 3 3 3 1	6287	11	22	207	11	7	4.71	0.17	0.16	0.306	0.012	—	51.8	A
170	A400 K	21	1 1 1 1 0	8569	20	—	157	28	—	2.49	0.53	—	—	—	—	17.3	B*
171	A400 L	21	2 2 0 0 0	6788	50	28	—	—	—	—	—	—	—	—	—	38.1	D
172	A400 1	21	1 1 1 1 0	6617	21	—	135	38	—	4.49	0.61	—	—	—	—	13.3	B
173	A400 2	21	1 1 0 0 0	8501	71	—	—	—	—	—	—	—	—	—	—	10.7	D
174	J28 A	22	2 2 2 2 2	8717	17	10	233	14	29	4.75	0.24	0.30	0.276	0.010	0.006	41.0	A
175	J28 B	22	1 1 1 1 1	8661	19	—	210	14	—	4.09	0.32	—	0.265	0.012	—	34.6	A
176	J28 C	22	1 1 1 1 1	8621	20	—	207	16	—	4.81	0.33	—	0.285	0.012	—	27.1	B
177	A419 A	23	4 4 4 4 4	20331	13	7	323	12	14	5.55	0.15	0.17	0.344	0.006	0.010	87.3	A
178	A419 B	23	4 4 3 3 3	21197	10	9	110	8	16	3.56	0.22	0.16	0.229	0.009	0.012	50.3	B
179	A419 C	23	3 3 3 3 3	21099	12	20	277	10	8	6.33	0.16	0.42	0.310	0.007	0.008	75.7	A
180	A419 D	23	3 3 3 3 3	20909	10	3	222	7	6	5.13	0.15	0.24	0.302	0.007	0.003	79.7	A
182	A419 F	23	3 3 3 3 3	19510	11	19	250	9	15	5.39	0.16	0.27	0.319	0.007	0.007	75.4	A
183	A419 G	123	3 3 3 3 3	11818	10	5	184	6	0	4.76	0.14	0.16	0.300	0.007	0.003	90.4	A
184	A419 H	23	4 4 4 4 4	20122	11	17	249	7	17	5.19	0.13	0.52	0.329	0.006	0.015	96.6	A
185	A419 I	23	4 4 4 4 4	20382	10	7	232	8	17	4.89	0.15	0.28	0.300	0.007	0.011	80.4	A
186	A419 J	123	3 3 3 3 3	11887	10	27	214	7	18	5.54	0.14	0.22	0.338	0.007	0.005	82.6	A
187	A419 1	23	3 3 3 3 3	19943	10	19	143	7	6	4.90	0.21	0.37	0.292	0.008	0.005	58.5	A
188	A419 2	23	3 3 3 3 3	20108	10	18	142	8	11	4.59	0.21	0.75	0.277	0.009	0.017	52.8	B
189	A496 A	24	5 5 5 5 5	9835	11	21	248	10	16	5.67	0.15	0.28	0.323	0.007	0.019	71.7	A
190	A496 B	24	3 3 3 3 3	10393	11	7	222	9	12	5.40	0.14	0.23	0.305	0.007	0.007	71.9	A
192	A496 D	24	3 3 3 3 3	9677	11	5	204	9	5	5.31	0.15	0.09	0.307	0.007	0.005	74.8	A
193	A496 1	24	3 3 3 3 3	9405	12	10	204	9	14	4.90	0.15	0.18	0.290	0.008	0.004	71.3	A
194	A496 2	24	4 4 4 4 4	9715	10	19	232	9	33	4.93	0.12	0.11	0.273	0.006	0.018	82.4	A
195	A496 3	24	4 4 4 4 4	8883	9	32	187	6	10	4.64	0.11	0.14	0.283	0.006	0.005	98.1	A
196	J34 A	25	2 2 2 2 1	10790	22	13	273	30	20	7.10	0.31	1.36	0.310	0.023	—	24.0	B
197	J34 B	25	2 2 2 2 1	12237	15	6	223	15	7	4.89	0.25	0.69	0.304	0.016	—	36.5	A
198	J34 C	25	1 1 1 1 1	10380	27	—	240	39	—	4.45	0.53	—	0.293	0.022	—	14.3	C
199	J34 D	25	2 2 2 2 1	10329	14	16	210	12	21	4.81	0.22	0.27	0.272	0.014	—	39.3	A
200	J34 E	25	2 2 2 2 1	11610	12	1	153	13	40	4.97	0.26	0.13	0.291	0.014	—	31.8	B
201	J34 F	25	1 1 1 1 1	11743	19	—	232	15	—	4.42	0.23	—	0.311	0.012	—	37.5	A
202	J34 1	25	2 2 1 1 0	10882	22	26	251	28	—	5.74	0.28	—	—	—	—	24.1	A
203	J34 2	25	2 2 1 1 0	10436	19	6	193	21	—	3.80	0.38	—	—	—	—	24.2	A
204	J34 3	25	1 1 1 1 0	10777	17	—	71	37	—	4.23	0.76	—	—	—	—	10.5	C*
205	J34 4	25	1 1 1 1 1	10808	35	—	238	42	—	6.24	0.50	—	0.322	0.028	—	23.2	B
206	J10 A	26	4 4 4 4 4	8918	17	32	293	17	25	5.10	0.20	0.13	0.313	0.009	0.017	46.7	A
207	J10 B	26	3 2 0 0 0	8087	50	30	—	—	—	—	—	—	—	—	—	13.2	D
208	J10 C	26	5 5 5 5 5	9026	17	7	124	22	10	4.88	0.66	1.13	0.241	0.016	0.043	27.4	A

Table 5 – continued

(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)	(11)	(12)	(13)	(14)	(15)	(16)	(17)	(18)
GIN	Galaxy Name	CAN	N $s z \sigma b 2$	cz (km s^{-1})	Δcz	δcz	σ (km s^{-1})	$\Delta \sigma$	$\delta \sigma$	Mgb	ΔM_{gb} (\AA)	δM_{gb}	M_{g_2}	ΔM_{g_2} (mag)	δM_{g_2}	S/N	Q
209	P597-1 A	27	2 2 2 2 0	4496	21	24	233	40	20	6.17	0.61	0.16	—	—	—	25.3	A
210	P597-1 B	27	3 3 3 3 0	4712	16	33	167	25	13	5.12	0.51	0.15	—	—	—	30.0	B
211	P597-1 C	27	2 2 2 2 0	4165	25	9	200	39	18	5.26	0.77	0.41	—	—	—	23.3	B
212	P597-1 D	27	2 2 2 2 0	4165	26	9	207	39	18	5.40	0.72	0.41	—	—	—	23.4	A
213	P597-1 E	27	2 2 1 1 0	5167	39	33	175	94	—	5.59	1.88	—	—	—	—	7.4	C
215	J35 A	28	2 2 1 1 1	9673	48	271	337	99	—	6.04	1.15	—	0.367	0.041	—	10.5	C
216	J35 B	28	1 1 1 1 1	9461	54	—	324	83	—	4.44	0.70	—	0.370	0.026	—	16.2	B
217	J35 C	28	1 1 1 1 1	9530	37	—	226	46	—	4.60	0.60	—	0.351	0.023	—	16.2	B
218	J35 D	28	1 1 1 1 1	10061	24	—	62	30	—	1.02	1.25	—	0.204	0.044	—	8.8	C
219	J35 E	128	1 1 1 1 0	16920	28	—	379	34	—	5.21	0.27	—	—	—	—	40.5	A
220	J35 I	28	1 1 1 1 1	11423	18	—	153	12	—	4.50	0.29	—	0.262	0.013	—	34.3	A
221	J34/35 A	29	1 1 1 1 1	9502	22	—	198	26	—	3.19	0.64	—	0.249	0.021	—	20.0	B
222	J34/35 B	29	1 1 1 1 1	10733	32	—	220	39	—	4.10	0.59	—	0.277	0.021	—	17.4	A
223	J34/35 C	29	1 1 1 1 1	8994	18	—	124	16	—	3.94	0.60	—	0.175	0.024	—	15.1	C
224	J34/35 D	29	1 1 1 1 1	8774	17	—	75	13	—	2.90	0.44	—	0.182	0.024	—	17.6	C
226	J34/35 2	129	1 1 1 1 1	16483	26	—	210	20	—	5.87	0.33	—	0.312	0.020	—	30.5	B
227	J34/35 3	29	1 1 1 1 1	9455	24	—	69	6	—	4.79	0.27	—	0.266	0.019	—	29.7	B
229	P777-1 B	130	2 2 2 2 2	16760	15	7	211	14	4	5.75	0.23	0.61	0.357	0.011	0.002	31.3	B
230	P777-1 C	30	2 2 2 2 2	12711	19	16	243	20	4	5.27	0.30	0.28	0.344	0.012	0.003	31.5	A
232	P777-2 A	31	2 2 2 2 1	7356	14	22	199	9	8	4.72	0.15	0.25	0.300	0.009	—	51.1	A
233	P777-2 B	131	2 2 2 2 2	9139	12	3	156	8	4	4.31	0.15	0.42	0.246	0.008	0.014	52.3	A
235	P777-2 D	131	1 1 1 1 1	9713	16	—	140	7	—	4.72	0.17	—	0.276	0.010	—	43.0	A
236	P777-2 E	31	1 1 1 1 1	7309	17	—	214	12	—	4.65	0.17	—	0.300	0.011	—	39.6	A
237	P777-2 F	31	2 2 2 2 2	7198	11	5	58	11	11	2.38	0.40	0.16	0.184	0.016	0.010	18.1	C
238	P777-2 I	131	1 1 1 1 1	9301	24	—	182	22	—	5.85	0.37	—	0.299	0.015	—	21.2	B
239	P777-3 A	32	3 3 3 3 3	16258	12	43	285	11	7	5.28	0.12	0.30	0.325	0.006	0.008	82.7	A
240	P777-3 B	32	3 3 3 3 3	16583	10	8	255	8	8	4.81	0.10	0.04	0.281	0.006	0.012	100.4	A
241	P777-3 C	232	4 4 4 4 4	8926	9	20	220	6	5	5.30	0.10	0.20	0.313	0.005	0.010	113.6	A
242	P777-3 D	332	3 3 3 3 3	4758	9	10	127	5	3	4.10	0.08	0.10	0.249	0.005	0.002	133.3	A
243	P777-3 E	132	3 3 3 3 3	13461	10	8	218	8	2	5.26	0.12	0.12	0.318	0.007	0.007	83.1	A
244	P777-3 F	32	3 3 3 3 3	16674	11	11	283	10	12	5.54	0.10	0.15	0.336	0.006	0.014	94.5	A
245	P777-3 G	32	4 4 4 4 4	16382	11	14	306	9	12	5.33	0.11	0.13	0.342	0.006	0.019	105.4	A
246	P777-3 H	132	3 3 3 3 3	13270	10	10	262	8	6	4.97	0.10	0.11	0.301	0.006	0.003	110.0	A
247	P777-3 I	32	3 3 3 3 3	16179	13	14	264	12	14	4.83	0.17	0.29	0.299	0.008	0.014	69.7	A
248	P777-3 2	132	2 2 2 2 2	13765	12	15	93	13	38	3.58	0.45	0.00	0.161	0.017	0.009	28.8	C
249	A533-1 A	33	2 2 2 2 2	12423	28	85	315	35	3	4.67	0.31	0.18	0.348	0.012	0.000	34.1	A
250	A533-1 B	33	1 1 1 1 1	11963	26	—	182	49	—	5.04	0.40	—	0.322	0.015	—	22.6	B
251	A533-1 C	33	2 2 2 2 2	12599	27	28	274	39	9	4.67	0.42	0.46	0.354	0.018	0.001	31.0	A
252	A533-1 D	33	2 2 2 2 2	10346	17	29	214	22	11	4.16	0.35	0.25	0.333	0.015	0.036	29.4	A
253	A533-1 E	33	1 1 1 1 1	11076	23	—	163	27	—	4.81	0.66	—	0.257	0.024	—	16.4	B
254	A533-1 F	34	1 1 1 1 1	14262	36	—	261	44	—	4.31	0.43	—	0.319	0.027	—	26.3	B
255	A533-1 G	34	1 1 1 1 1	14061	34	—	237	38	—	4.67	0.54	—	0.326	0.028	—	25.7	B
257	A533 A	34	2 2 2 2 1	14512	21	31	274	25	12	5.09	0.29	0.29	0.350	0.028	—	31.6	A
258	A533 B	34	1 1 1 1 1	14287	74	—	381	122	—	4.03	0.77	—	0.307	0.023	—	14.8	B
259	A533 C	34	2 2 2 1 1	15838	24	5	303	33	29	4.61	1.62	—	0.371	0.041	—	28.2	A
260	A533 D	34	2 2 2 2 1	14922	27	27	326	37	35	5.85	0.35	0.54	0.305	0.025	—	29.3	A
261	A533 E	34	2 2 1 1 0	14849	22	14	251	28	—	4.75	0.36	—	—	—	—	23.9	A
262	A533 F	34	1 1 0 0 0	13256	71	—	—	—	—	—	—	—	—	—	—	12.1	D
263	A533 G	33	2 1 1 1 0	12362	18	—	208	16	—	4.97	0.23	—	—	—	—	33.5	A
264	A533 H	33	1 1 0 0 0	10802	71	—	—	—	—	—	—	—	—	—	—	19.3	D*
266	A548-1 A	35	1 1 1 1 1	11874	19	—	207	17	—	5.96	0.31	—	0.384	0.012	—	25.8	B
267	A548-1 B	35	1 1 1 1 1	12000	22	—	295	24	—	5.85	0.22	—	0.345	0.011	—	34.0	A
268	A548-1 C	35	3 3 3 3 3	12806	10	9	148	9	34	4.29	0.23	0.42	0.241	0.009	0.045	45.6	A
269	A548-1 D	35	3 3 3 3 3	13096	12	11	270	10	2	5.19	0.16	0.26	0.313	0.007	0.006	72.0	A
270	A548-1 E	36	3 3 3 3 3	13615	14	7	323	14	33	5.22	0.15	0.16	0.331	0.006	0.009	73.4	A
271	A548-1 F	35	3 3 3 3 3	12775	9	13	115	8	28	4.34	0.22	0.32	0.233	0.009	0.012	46.3	A
272	A548-1 G	35	3 3 3 3 3	11287	11	21	225	9	8	4.45	0.14	0.19	0.269	0.007	0.016	72.2	A
273	A548-1 H	35	1 1 1 1 1	11845	18	—	244	14	—	4.82	0.16	—	0.281	0.010	—	39.9	A
274	A548-1 I	36	1 1 1 1 1	12769	19	—	246	17	—	5.26	0.22	—	0.300	0.010	—	36.2	A
275	A548-1 J	36	1 1 1 1 1	12653	19	—	243	18	—	5.41	0.25	—	0.319	0.012	—	33.4	A
276	A548-1 K	35	3 1 1 1 1	12651	19	—	178	16	—	4.18	0.40	—	0.268	0.015	—	27.4	B
277	A548-1 L	35	1 1 1 1 1	11386	17	—	238	13	—	4.28	0.20	—	0.237	0.010	—	43.3	A

Table 5 – *continued*

(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)	(11)	(12)	(13)	(14)	(15)	(16)	(17)	(18)
GIN	Galaxy Name	CAN	N $s z \sigma b 2$	cz (km s^{-1})	Δcz	δcz	σ (km s^{-1})	$\Delta \sigma$	$\delta \sigma$	Mgb	ΔMgb (\AA)	δMgb	Mg_2	ΔMg_2 (mag)	δMg_2	S/N	Q
278	A548-1 M	35	1 1 1 1 1	11558	17	—	208	13	—	4.73	0.22	—	0.258	0.010	—	35.2	A
279	A548-1 N	35	3 3 3 3 3	11541	12	25	250	10	11	5.53	0.16	0.21	0.301	0.006	0.013	68.5	A
280	A548-1 O	35	3 3 3 3 3	10730	10	7	177	8	12	4.32	0.14	0.34	0.247	0.006	0.014	70.9	A
281	A548-1 P	36	1 1 1 1 1	12709	21	—	250	22	—	4.67	0.28	—	0.280	0.012	—	28.3	B
282	A548-1 Q	36	4 4 4 4 3	13755	11	50	223	10	27	5.28	0.18	0.44	0.330	0.008	0.005	57.9	A
283	A548-1 R	35	4 2 2 2 2	12645	20	20	230	24	10	4.06	0.35	0.09	0.285	0.014	0.009	26.0	B
285	A548-2 A	35	5 5 5 5 3	11861	11	107	269	11	12	5.09	0.12	0.24	0.303	0.007	0.016	66.0	A
286	A548-2 B	35	3 3 3 3 2	11548	14	26	284	14	17	5.05	0.15	0.27	0.274	0.008	0.028	53.2	A
287	A548-2 C	35	3 3 0 0 0	11040	41	34	—	—	—	—	—	—	—	—	—	30.5	D*
288	A548-2 D	35	2 2 2 2 1	11831	15	11	272	14	10	5.16	0.16	0.06	0.295	0.010	—	50.8	A
289	A548-2 E	35	1 1 1 1 0	11141	18	—	80	51	—	4.11	0.96	—	—	—	—	7.9	C
290	A548-2 F	35	4 4 3 3 1	11345	12	32	215	12	17	4.25	0.19	0.15	0.287	0.013	—	47.4	A
291	A548-2 G	35	3 3 3 3 1	11919	10	11	173	8	7	4.32	0.14	0.31	0.240	0.010	—	52.5	A
292	A548-2 H	35	4 4 4 4 1	12765	11	22	220	11	12	4.90	0.17	0.13	0.250	0.014	—	49.0	A
293	A548-2 I	136	2 2 2 2 1	8998	14	27	226	11	4	4.75	0.15	0.04	0.296	0.010	—	45.7	A
294	A548-2 J	136	2 2 2 2 1	8529	11	20	62	12	9	3.56	0.32	0.19	0.202	0.019	—	24.9	B
295	A548-2 K	35	4 4 4 4 2	12662	11	17	197	10	12	5.61	0.19	0.59	0.304	0.010	0.008	48.6	A
296	A548-2 L	35	3 1 0 0 0	12569	71	—	—	—	—	—	—	—	—	—	—	12.3	D
297	A548-2 M	35	3 1 0 0 0	10600	71	—	—	—	—	—	—	—	—	—	—	7.9	D*
298	A548-2 N	36	2 2 2 2 1	11667	20	9	311	24	20	4.83	0.25	0.24	0.290	0.015	—	38.4	A
299	A548-2 O	35	3 3 3 3 1	11310	14	8	204	16	23	4.66	0.27	0.43	0.303	0.028	—	36.5	A
300	A548-2 1	35	3 3 3 3 0	11161	10	16	69	15	2	3.53	0.31	0.25	—	—	—	26.2	B
301	A548-2 2	35	3 3 3 3 0	11581	16	37	277	19	19	4.57	0.21	0.07	—	—	—	39.7	A
302	J11 A	37	7 7 7 7 3	8853	9	22	217	8	11	4.87	0.16	0.14	0.297	0.010	0.007	75.2	A
303	J11 B	37	5 4 4 4 3	8825	11	20	249	10	11	4.93	0.18	0.13	0.305	0.009	0.003	81.6	A
304	J11 C	37	2 2 2 2 2	8915	19	60	219	20	3	3.98	0.43	0.15	0.261	0.015	0.001	40.7	A
305	J11 D	37	1 1 1 1 1	8705	21	—	143	19	—	4.10	0.49	—	0.239	0.021	—	30.5	A
306	J11 E	137	2 1 1 1 1	11399	20	—	126	12	—	2.62	0.67	—	0.161	0.025	—	23.3	B
307	J11 F	37	3 3 3 3 1	8763	13	17	136	12	17	3.93	0.45	0.28	0.237	0.024	—	28.5	B
308	J11 1	37	1 1 1 1 1	9265	33	—	210	37	—	3.79	0.58	—	0.230	0.029	—	21.0	B
310	J12 A	38	1 1 1 1 1	12115	41	—	268	45	—	4.00	0.74	—	0.270	0.025	—	19.8	B
311	J12 B	338	1 1 1 1 1	10873	36	—	298	37	—	4.66	0.51	—	0.331	0.019	—	28.6	A
312	J12 C	38	1 1 1 1 1	13206	34	—	202	39	—	5.01	0.77	—	0.322	0.028	—	17.3	B
313	J12 D	238	1 1 1 1 1	17845	28	—	232	40	—	5.82	0.66	—	0.312	0.024	—	19.0	B
314	J12 E	38	2 2 2 2 2	13288	15	11	213	11	23	4.93	0.28	0.17	0.312	0.011	0.008	51.4	A
315	J12 F	38	2 2 2 2 2	12542	18	23	258	17	12	5.75	0.29	0.14	0.335	0.011	0.002	43.6	A
316	J12 1	138	2 2 2 2 1	15408	17	37	218	14	20	4.53	0.31	0.15	0.320	0.016	—	44.3	A
317	J12 2	38	1 1 1 1 1	34905	65	—	284	124	—	4.49	1.08	—	0.331	0.048	—	10.7	C
318	J13 A	39	1 1 1 1 1	9333	32	—	264	31	—	4.39	0.56	—	0.286	0.022	—	20.0	A
319	J13 B	39	1 1 1 1 1	9323	24	—	183	23	—	4.28	0.54	—	0.238	0.020	—	27.9	B
320	J13 C	39	1 1 1 1 1	8724	27	—	130	45	—	2.35	1.04	—	0.107	0.038	—	11.7	C
321	J13 D	39	1 1 1 1 1	8728	32	—	217	34	—	3.93	0.64	—	0.269	0.022	—	21.3	A
322	J13 E	239	1 1 1 1 1	17438	56	—	368	64	—	5.91	0.69	—	0.305	0.024	—	21.2	B
323	J13 F	39	1 1 1 1 1	8625	25	—	172	26	—	3.75	0.69	—	0.253	0.022	—	24.0	B
324	J13 G	39	1 1 1 1 1	9060	37	—	303	39	—	4.96	0.50	—	0.313	0.022	—	28.3	A
325	J13 H	39	1 1 1 1 1	8764	21	—	248	30	—	4.83	0.56	—	0.282	0.023	—	24.6	A
326	J13 I	39	1 1 1 1 1	8857	20	—	179	29	—	4.20	0.95	—	0.287	0.025	—	19.7	A
328	J13 K	39	1 1 1 1 0	8927	24	—	199	26	—	4.25	0.57	—	—	—	—	24.1	A
329	J13 L	339	1 1 1 1 0	15651	31	—	225	40	—	5.40	0.63	—	—	—	—	13.6	C
330	J13 M	39	2 2 2 2 2	8808	16	20	145	18	21	4.56	0.38	0.08	0.239	0.013	0.011	28.3	B
331	J13 N	139	1 1 0 0 0	18872	72	—	—	—	—	—	—	—	—	—	—	17.3	D
332	J13 O	39	1 1 0 0 0	8762	75	—	—	—	—	—	—	—	—	—	—	11.8	D
333	J13 P	39	1 1 1 1 1	8671	17	—	180	10	—	4.22	0.27	—	0.269	0.013	—	44.6	A
334	J13 1	139	1 1 1 1 0	18851	47	—	331	65	—	5.57	0.67	—	—	—	—	15.0	C
335	J13 2	139	1 1 1 1 0	19410	25	—	178	29	—	6.21	0.64	—	—	—	—	13.4	C
336	J36 A	40	2 2 2 2 2	11230	23	11	289	25	38	5.85	0.50	0.71	0.360	0.016	0.001	29.4	A
337	J36 B	40	2 2 2 2 2	11624	19	14	231	21	33	6.14	0.50	0.47	0.318	0.019	0.008	23.7	B
338	J36 C	140	2 2 2 2 0	9966	27	10	194	23	5	5.41	0.60	0.35	—	—	—	18.5	B
339	J36 D	40	3 3 3 3 3	11982	16	26	250	14	14	5.48	0.36	0.35	0.313	0.014	0.010	35.5	A
340	J36 1	340	2 2 1 1 1	28974	32	10	159	54	—	4.42	0.99	—	0.316	0.044	—	11.2	C
341	J36 2	240	1 1 1 1 0	17291	27	—	253	23	—	5.31	0.35	—	—	—	—	28.6	A
342	J36 3	140	1 1 1 1 1	9647	21	—	187	14	—	5.16	0.35	—	0.295	0.015	—	29.6	B

Table 5 – continued

(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)	(11)	(12)	(13)	(14)	(15)	(16)	(17)	(18)
GIN	Galaxy Name	CAN	N $s z \sigma b 2$	cz (km s^{-1})	Δcz	δcz	σ (km s^{-1})	$\Delta \sigma$	$\delta \sigma$	Mgb	ΔM_{gb} (\AA)	δM_{gb}	M_{g_2}	ΔM_{g_2} (mag)	δM_{g_2}	S/N	Q
343	J14 A	41	1 1 1 1 1	8722	30	—	255	30	—	4.45	0.52	—	0.303	0.018	—	29.4	A
344	J14 B	41	1 1 1 1 1	8955	28	—	205	28	—	4.30	0.62	—	0.277	0.023	—	25.5	A
345	J14 C	41	1 1 1 1 1	8632	31	—	158	61	—	4.16	1.78	—	0.193	0.048	—	7.6	C
346	J14 D	41	1 1 1 1 1	9247	34	—	213	78	—	4.00	1.12	—	0.208	0.042	—	9.4	C
347	J14 1	241	1 1 1 1 0	20534	81	—	529	142	—	5.49	0.58	—	—	—	—	14.6	C
348	J14 2	41	2 2 2 2 2	9307	16	11	157	14	19	4.22	0.43	0.37	0.250	0.016	0.023	23.6	C
349	J14 3	141	1 1 1 1 1	13280	25	—	196	30	—	5.85	0.60	—	0.299	0.024	—	14.5	C*
350	J14 4	141	1 1 1 1 0	14090	19	—	164	11	—	2.83	0.31	—	—	—	—	32.0	B
351	J14 5	241	1 1 1 1 1	20345	41	—	151	61	—	4.16	1.22	—	0.249	0.050	—	9.4	C
352	J14-1 A	42	1 1 1 1 1	8721	29	—	247	29	—	4.53	0.46	—	0.307	0.020	—	29.4	A
353	J14-1 B	42	1 1 1 1 1	8954	28	—	205	28	—	4.29	0.61	—	0.276	0.021	—	25.5	A
354	J14-1 C	42	1 1 1 1 1	8306	28	—	179	30	—	3.96	0.69	—	0.258	0.023	—	21.3	A
355	J14-1 D	42	1 1 1 1 1	8842	21	—	122	21	—	3.99	0.59	—	0.238	0.021	—	26.2	A
356	J14-1 E	42	2 2 2 2 0	9138	19	6	182	15	2	4.38	0.53	0.03	—	—	—	23.2	B
357	J14-1 F	42	2 2 2 2 1	8734	14	9	188	8	20	4.31	0.25	0.13	0.263	0.013	—	59.4	A
358	J14-1 1	142	1 1 1 1 1	16133	32	—	282	38	—	6.35	0.29	—	0.311	0.016	—	30.9	B
359	J14-1 2	142	1 1 1 1 0	15649	37	—	272	34	—	4.28	0.50	—	—	—	—	20.9	B
360	J14-1 3	142	1 1 0 0 0	16096	72	—	—	—	—	—	—	—	—	—	—	13.6	D
361	J14-1 4	142	1 1 0 0 0	15574	73	—	—	—	—	—	—	—	—	—	—	25.7	D*
362	A1983 A	43	1 1 1 1 1	13601	47	—	298	56	—	5.36	0.67	—	0.350	0.030	—	18.3	A
363	A1983 B	43	2 2 2 2 2	13205	19	18	289	17	7	5.42	0.26	0.04	0.327	0.011	0.009	55.2	A
364	A1983 C	43	1 1 1 1 1	14069	36	—	178	48	—	6.21	1.01	—	0.364	0.040	—	9.9	B
365	A1983 D	43	1 1 1 1 1	13749	43	—	254	53	—	5.55	0.77	—	0.332	0.030	—	15.1	B
366	A1983 E	43	1 1 1 1 1	13426	37	—	275	51	—	5.04	0.61	—	0.278	0.024	—	18.0	B
367	A1983 F	43	1 1 1 1 1	13819	18	—	179	13	—	4.13	0.28	—	0.278	0.015	—	35.2	B
368	A1983 G	43	1 1 1 1 1	13718	18	—	221	13	—	4.79	0.28	—	0.290	0.013	—	42.9	A
369	A1983 1	43	2 2 2 2 1	13631	22	8	210	19	20	4.80	0.37	0.14	0.314	0.036	—	27.5	B
370	A1983 2	43	1 1 1 1 0	13786	18	—	127	12	—	2.96	0.55	—	—	—	—	22.7	C*
371	A1983 3	143	1 1 1 1 1	21438	23	—	218	20	—	4.05	0.46	—	0.357	0.018	—	33.8	B
372	A1991 A	44	1 1 1 1 1	17752	46	—	365	67	—	6.99	0.81	—	0.312	0.029	—	14.5	B
373	A1991 B	44	1 1 1 1 1	16954	35	—	260	61	—	7.28	1.50	—	0.297	0.043	—	9.1	C
374	A1991 C	44	1 1 1 1 1	16696	30	—	229	56	—	4.65	1.65	—	0.227	0.048	—	9.5	C
375	A1991 D	144	1 1 1 1 1	13213	44	—	369	65	—	5.45	0.85	—	0.303	0.028	—	15.4	B
376	A1991 E	144	1 1 1 1 1	13167	51	—	446	84	—	5.75	0.93	—	0.329	0.027	—	16.9	B
377	A1991 F	244	1 1 1 1 0	21199	64	—	383	100	—	7.68	0.63	—	—	—	—	13.0	C
378	A1991 1	44	1 1 1 1 1	17578	20	—	83	17	—	4.01	0.64	—	0.277	0.025	—	22.3	C
379	A1991 2	44	2 2 2 2 0	16636	28	20	266	38	12	5.25	0.51	0.55	—	—	—	18.9	C
380	A1991 3	44	2 2 2 2 0	16214	23	28	257	20	17	5.87	0.27	0.30	—	—	—	33.5	B
381	A1991 4	44	1 1 0 0 0	16504	73	—	—	—	—	—	—	—	—	—	—	25.4	D
382	A1991 5	44	1 1 1 1 0	17261	23	—	189	19	—	5.91	0.38	—	—	—	—	22.8	B
383	J16 A	45	1 1 1 1 1	10948	44	—	304	49	—	5.27	0.73	—	0.337	0.021	—	22.3	A
384	J16 B	245	1 1 1 1 1	14024	40	—	288	44	—	4.80	0.67	—	0.321	0.021	—	23.1	A
385	J16 C	345	1 1 1 1 1	7789	37	—	216	44	—	4.24	0.94	—	0.219	0.033	—	14.9	B
386	J16 D	145	1 1 1 1 1	15684	51	—	309	65	—	4.77	0.82	—	0.338	0.023	—	16.9	B
387	J16 E	45	2 2 2 2 2	10810	15	14	255	11	11	4.89	0.22	0.12	0.305	0.011	0.009	62.4	A
388	J16 F	45	2 2 2 2 1	11405	14	53	241	12	7	4.65	0.21	0.21	0.301	0.014	—	77.1	A
390	J16 1	145	1 1 1 1 1	15972	23	—	161	14	—	3.85	0.38	—	0.277	0.016	—	30.0	B
392	J16 3	45	2 1 0 0 0	11533	72	—	—	—	—	—	—	—	—	—	—	13.6	D
393	J16 4	145	1 1 1 1 0	15616	24	—	252	19	—	4.64	0.36	—	—	—	—	33.9	A
394	J16 5	45	1 1 1 1 1	10930	20	—	187	13	—	5.58	0.32	—	0.326	0.016	—	33.2	A
395	J16-W A	46	1 1 1 1 0	10800	22	—	258	35	—	5.57	0.34	—	—	—	—	76.0	A
396	J16-W B	46	1 1 1 1 1	11325	24	—	127	27	—	4.33	0.71	—	0.282	0.027	—	16.5	B
397	J16-W C	46	1 1 1 1 0	10404	22	—	218	33	—	4.39	0.36	—	—	—	—	58.4	B
398	J16-W D	46	1 1 1 1 1	11137	27	—	325	24	—	5.40	0.42	—	0.356	0.021	—	29.5	A
399	J16-W E	46	1 1 1 1 1	10788	35	—	209	44	—	5.28	0.80	—	0.273	0.028	—	15.9	B
400	J16-W F	46	1 1 1 1 0	10717	21	—	142	70	—	3.51	0.36	—	—	—	—	67.9	A
401	J16-W G	46	1 1 1 1 1	11325	23	—	203	16	—	4.83	0.61	—	0.306	0.023	—	21.9	A
402	J16-W H	46	1 1 1 1 0	12462	23	—	234	36	—	5.23	0.38	—	—	—	—	48.7	B
403	J16-W I	146	1 1 1 1 1	23945	25	—	215	23	—	2.80	0.39	—	0.233	0.016	—	29.3	B
404	J16-W 1	46	1 1 1 1 1	11269	21	—	93	9	—	2.27	0.38	—	0.156	0.016	—	30.0	B
406	A2063-S A	50	2 2 2 2 1	11785	19	42	221	48	42	5.07	0.55	0.11	0.286	0.038	—	26.4	A
408	A2063-S C	47	1 1 1 1 0	13876	27	—	246	23	—	6.12	0.38	—	—	—	—	26.8	B

Table 5 – *continued*

(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)	(11)	(12)	(13)	(14)	(15)	(16)	(17)	(18)
GIN	Galaxy Name	CAN	N $s z \sigma b 2$	cz (km s^{-1})	Δcz	δcz	σ (km s^{-1})	$\Delta \sigma$	$\delta \sigma$	Mgb	ΔM_{gb} (\AA)	δM_{gb}	M_{g_2}	ΔM_{g_2} (mag)	δM_{g_2}	S/N	Q
409	A2040 A	48	1 1 1 1 0	13709	24	—	187	28	—	6.14	0.60	—	—	—	—	15.1	C
410	A2040 B	48	1 1 1 1 0	13722	24	—	223	20	—	5.11	0.36	—	—	—	—	27.9	A
411	A2040 C	48	2 2 2 2 2	13213	15	28	196	12	4	5.20	0.28	0.01	0.300	0.012	0.010	37.1	B
412	A2040 D	48	1 1 1 1 1	31895	23	—	239	22	—	5.34	0.34	—	0.311	0.017	—	25.5	B
413	A2040 E	48	1 1 1 1 1	12927	63	—	336	125	—	4.97	1.11	—	0.283	0.039	—	8.1	B
414	A2040 F	48	1 1 1 1 1	13449	42	—	353	61	—	4.41	0.53	—	0.314	0.027	—	29.8	A
415	A2040 G	48	1 1 0 0 0	13520	75	—	—	—	—	—	—	—	—	—	—	16.6	D
416	A2040 H	148	1 1 0 0 0	23099	73	—	—	—	—	—	—	—	—	—	—	6.6	D
418	A2040 2	48	1 1 1 1 1	13723	48	—	272	82	—	4.58	0.99	—	0.312	0.037	—	9.6	A
419	A2052 A	49	2 2 2 2 2	10308	30	9	300	45	19	6.04	0.71	0.32	0.323	0.023	0.020	18.2	B
420	A2052 B	49	2 2 2 2 2	9359	17	0	168	21	25	5.81	0.42	0.13	0.299	0.014	0.003	24.2	B
421	A2052 C	49	1 1 1 1 1	9983	21	—	218	15	—	5.63	0.30	—	0.302	0.014	—	36.5	A
422	A2052 D	49	1 1 1 1 1	9085	20	—	134	18	—	4.90	0.58	—	0.255	0.023	—	16.5	C
423	A2052 E	49	1 1 1 1 1	9571	22	—	160	20	—	4.67	0.61	—	0.255	0.024	—	16.3	C
424	A2052 F	49	1 1 1 1 1	10328	19	—	119	15	—	4.39	0.49	—	0.183	0.024	—	18.3	C
425	A2052 G	49	1 1 1 1 0	10974	22	—	230	15	—	4.54	0.41	—	—	—	—	51.0	A
426	A2063 A	50	1 1 1 1 1	10223	35	—	240	52	—	5.97	0.69	—	0.344	0.023	—	17.3	B
427	A2063 B	50	2 2 2 2 2	10354	16	90	175	11	5	4.73	0.29	0.13	0.268	0.012	0.001	36.3	A
428	A2063 C	50	1 1 1 1 1	10351	21	—	198	14	—	4.36	0.33	—	0.307	0.014	—	31.8	A
429	A2063 D	50	1 1 1 1 1	10375	20	—	147	12	—	5.11	0.39	—	0.304	0.015	—	25.5	B
430	A2063 E	50	1 1 1 1 1	9835	23	—	204	18	—	5.44	0.38	—	0.286	0.017	—	26.3	B
431	A2063 F	50	1 1 1 1 1	10732	19	—	137	10	—	4.82	0.38	—	0.277	0.016	—	26.2	B
432	A2063 G	50	1 1 1 1 1	9868	22	—	186	16	—	6.16	0.39	—	0.367	0.017	—	26.8	B
433	A2063 H	50	1 1 1 1 0	10318	21	—	217	14	—	4.73	0.42	—	—	—	—	53.6	A
436	A2107 A	51	2 2 2 2 2	12525	25	39	324	31	46	4.41	0.46	0.60	0.347	0.014	0.029	25.7	B
437	A2107 B	51	1 1 1 1 1	12579	21	—	258	33	—	4.61	0.70	—	0.256	0.021	—	19.9	B
438	A2107 C	51	1 1 1 1 1	12165	26	—	172	29	—	4.16	0.70	—	0.218	0.023	—	17.8	B
439	A2107 D	51	1 1 1 1 1	12955	23	—	265	22	—	4.90	0.31	—	0.316	0.015	—	40.6	A
440	A2107 E	51	1 1 1 1 1	12388	22	—	136	13	—	4.51	0.36	—	0.225	0.016	—	26.5	B
441	A2107 I	51	2 1 1 1 1	12949	20	—	75	11	—	4.65	0.44	—	0.292	0.018	—	30.5	B
443	J17 A	52	2 2 2 2 2	10102	26	11	264	27	6	5.23	0.44	0.20	0.329	0.017	0.000	33.8	A
444	J17 B	52	2 2 2 2 2	9980	22	59	175	28	10	3.98	0.70	0.22	0.261	0.020	0.004	21.3	B
445	J17 C	52	1 1 1 1 1	8956	32	—	195	36	—	3.71	0.80	—	0.239	0.025	—	17.9	A
446	J17 D	52	1 1 1 1 1	9925	43	—	308	47	—	5.13	0.66	—	0.330	0.022	—	23.9	A
447	J17 E	52	1 1 1 1 1	9955	36	—	224	57	—	6.22	0.83	—	0.251	0.029	—	13.6	B
449	J17 G	52	1 1 1 1 1	10978	27	—	147	49	—	4.19	1.43	—	0.219	0.037	—	9.4	C
453	A2147 A	53	1 1 1 1 1	10588	38	—	264	59	—	4.82	0.72	—	0.289	0.022	—	15.4	B
454	A2147 B	53	1 1 1 1 1	10700	35	—	246	51	—	5.97	0.80	—	0.352	0.029	—	15.9	B
455	A2147 C	53	1 1 1 1 1	11654	39	—	221	69	—	5.83	1.13	—	0.348	0.039	—	10.5	B
456	A2147 D	53	1 1 1 1 1	11877	57	—	337	109	—	5.32	1.08	—	0.330	0.033	—	11.4	B
457	A2147 E	53	1 1 1 1 0	11032	22	—	253	16	—	4.42	0.41	—	—	—	—	59.1	A
459	A2147 G	53	1 1 1 1 1	11056	30	—	222	40	—	5.11	0.71	—	0.248	0.022	—	18.2	A
460	A2147 H	53	1 1 1 1 1	10602	48	—	259	100	—	4.57	1.13	—	0.334	0.036	—	10.3	C
461	A2147 I	53	1 1 1 1 1	10418	25	—	249	21	—	5.35	0.35	—	0.307	0.014	—	30.3	B
462	A2147 J	53	1 1 1 1 1	11538	21	—	155	14	—	4.94	0.40	—	0.302	0.015	—	22.1	B
466	A2147 4	53	1 1 1 1 0	10673	21	—	155	12	—	4.39	0.43	—	—	—	—	40.9	B
467	A2147 5	53	1 1 1 1 0	9498	21	—	216	14	—	3.91	0.42	—	—	—	—	51.4	A
468	P386-1 A	54	1 1 1 1 1	4410	21	—	249	27	—	4.46	0.52	—	0.309	0.019	—	30.3	A
469	P386-1 B	154	1 1 1 1 1	13830	22	—	191	34	—	4.56	0.94	—	0.310	0.028	—	16.0	B
471	P386-1 I	254	1 1 1 1 1	26612	74	—	382	160	—	4.54	1.45	—	0.308	0.049	—	6.6	B
472	P386-2 A	55	2 2 2 2 1	9607	15	29	259	11	2	4.90	0.21	0.10	0.296	0.013	—	68.6	A
473	P386-2 B	55	2 2 1 1 1	9715	17	21	170	10	—	4.51	0.28	—	0.287	0.013	—	44.1	A
474	P386-2 C	55	2 2 2 2 1	9589	17	11	210	11	1	4.90	0.24	0.09	0.287	0.013	—	49.0	A
476	P386-2 I	155	1 1 1 1 1	20285	32	—	134	48	—	4.59	1.86	—	0.229	0.056	—	4.9	B
478	A2148 A	56	1 1 1 1 1	26874	34	—	284	40	—	4.88	0.35	—	0.283	0.017	—	30.2	A
479	A2148 B	56	1 1 1 1 1	25701	29	—	255	30	—	4.83	0.33	—	0.274	0.016	—	30.6	A
480	A2148 C	256	1 1 1 1 1	12730	23	—	123	27	—	3.83	0.60	—	0.096	0.029	—	12.9	C*
481	A2148 D	156	1 1 1 1 1	7955	21	—	65	12	—	1.45	0.58	—	0.018	0.023	—	21.3	C*
482	A2148 E	356	1 1 0 0 0	21558	73	—	—	—	—	—	—	—	—	—	—	18.7	D
485	A2148 2	56	1 1 1 1 1	26265	55	—	278	96	—	4.54	1.37	—	0.288	0.045	—	6.3	B
487	J18 A	57	1 1 1 1 1	9630	54	—	384	60	—	5.65	0.63	—	0.371	0.019	—	24.9	A
488	J18 B	57	1 1 1 1 1	9018	24	—	128	27	—	3.97	0.66	—	0.234	0.022	—	20.2	A

Table 5 – continued

(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)	(11)	(12)	(13)	(14)	(15)	(16)	(17)	(18)
GIN	Galaxy Name	CAN	N $s z \sigma b 2$	cz (km s^{-1})	Δcz	δcz	σ (km s^{-1})	$\Delta \sigma$	$\delta \sigma$	Mgb	ΔM_{gb} (\AA)	δM_{gb}	M_{g_2}	ΔM_{g_2} (mag)	δM_{g_2}	S/N	Q
489	J18 C	57	1 1 1 1 1	9517	60	—	358	80	—	5.43	0.80	—	0.299	0.028	—	16.8	A
490	J18 D	57	2 2 2 2 1	9729	17	66	106	10	14	3.82	0.43	0.07	0.253	0.046	—	41.2	A
491	J18 E	57	1 1 1 1 1	9507	23	—	179	15	—	4.12	0.33	—	0.274	0.016	—	30.6	A
492	J18 F	57	1 1 1 1 1	9725	21	—	114	13	—	4.68	0.32	—	0.275	0.018	—	23.6	B
493	J18 I	157	1 1 1 1 1	16075	47	—	221	74	—	5.09	1.86	—	0.254	0.055	—	4.8	B
495	A2151 A	58	1 1 1 1 1	10526	68	—	409	128	—	5.91	0.91	—	0.349	0.028	—	14.8	B
496	A2151 B	58	1 1 1 1 1	10705	22	—	209	24	—	5.06	0.43	—	0.310	0.023	—	24.9	A
497	A2151 C	58	1 1 1 1 1	9917	46	—	224	73	—	4.80	1.63	—	0.272	0.055	—	5.7	B
498	A2151 D	58	1 1 1 1 1	9360	22	—	213	16	—	4.74	0.30	—	0.327	0.015	—	38.8	A
499	A2151 E	58	1 1 1 1 1	9936	21	—	162	10	—	3.49	0.31	—	0.250	0.015	—	40.4	A
500	A2151 F	58	1 1 1 1 1	10355	37	—	192	53	—	3.48	1.19	—	0.221	0.037	—	8.9	A
501	A2151 G	58	1 1 1 1 1	12076	33	—	366	51	—	5.12	0.33	—	0.325	0.016	—	36.3	A
502	A2151 H	58	1 1 1 1 1	11689	23	—	225	18	—	4.11	0.33	—	0.280	0.015	—	38.2	A
503	A2151 I	58	1 1 1 1 1	12020	21	—	108	10	—	4.22	0.36	—	0.203	0.016	—	27.7	B
504	A2151 J	58	1 1 1 1 1	11324	25	—	270	27	—	4.53	0.31	—	0.272	0.015	—	37.1	A
505	A2151 K	58	2 2 2 2 2	10595	16	9	195	11	19	4.72	0.24	0.50	0.298	0.011	0.011	45.0	A
506	A2151 L	58	1 1 1 1 1	11373	21	—	110	11	—	3.36	0.34	—	0.188	0.017	—	27.0	B
507	A2151 M	58	1 1 1 1 1	11044	27	—	323	29	—	5.03	0.26	—	0.304	0.015	—	46.7	A
508	A2151 N	58	1 1 1 1 1	11404	23	—	251	20	—	4.67	0.30	—	0.304	0.016	—	41.3	A
509	A2151 O	58	1 1 1 1 1	11631	24	—	213	19	—	4.97	0.31	—	0.282	0.016	—	33.9	B
510	A2151 I	58	1 1 1 1 1	11136	58	—	285	149	—	2.11	2.74	—	0.128	0.079	—	3.7	C
511	A2152 A	59	2 2 2 2 2	12066	20	13	280	19	0	5.96	0.33	0.44	0.359	0.013	0.003	32.3	A
512	A2152 B	59	1 1 1 1 1	13246	26	—	237	19	—	5.32	0.65	—	0.359	0.025	—	21.3	A
513	A2152 C	59	2 2 2 2 2	12311	17	7	240	15	13	4.56	0.33	0.30	0.313	0.012	0.016	50.1	A
514	A2152 D	53	2 2 2 2 2	10896	16	18	162	13	13	4.95	0.40	0.35	0.274	0.014	0.001	28.4	B
515	A2152 E	59	1 1 1 1 1	11893	24	—	134	33	—	5.58	0.46	—	0.275	0.017	—	17.1	B
517	A2152 G	53	2 2 1 1 0	10142	19	9	153	11	—	4.37	0.43	—	—	—	—	46.4	B
518	A2152 H	59	1 1 1 1 1	13283	30	—	280	26	—	5.19	0.61	—	0.321	0.024	—	20.6	A
519	A2152 I	59	1 1 1 1 1	13624	40	—	313	34	—	6.03	0.74	—	0.383	0.022	—	19.1	A
524	A2152 N	59	2 2 2 2 2	12930	19	21	150	13	16	3.98	0.46	0.02	0.267	0.016	0.003	32.4	B
525	A2152 I	59	1 1 1 1 1	13559	26	—	263	26	—	4.18	0.46	—	0.294	0.016	—	33.6	A
527	A2152 3	53	1 1 1 1 0	10412	20	—	112	12	—	3.59	0.45	—	—	—	—	33.7	B
528	P445-1 A	60	1 1 1 1 1	13618	25	—	252	25	—	5.02	0.28	—	0.313	0.016	—	35.5	A
529	P445-1 B	60	1 1 1 1 1	14323	24	—	181	19	—	4.19	0.31	—	0.233	0.018	—	26.2	A
530	P445-1 C	60	1 1 1 1 1	13880	30	—	283	35	—	4.78	0.30	—	0.304	0.017	—	33.3	A
531	P445-1 D	60	1 1 1 1 1	14329	23	—	161	15	—	3.53	0.38	—	0.217	0.016	—	27.5	A
532	P445-1 E	260	1 1 1 1 1	18545	30	—	247	33	—	4.23	0.38	—	0.365	0.017	—	27.5	B
533	P445-1 F	60	1 1 1 1 1	13817	24	—	211	18	—	5.16	0.31	—	0.315	0.016	—	34.6	A
534	P445-1 G	160	1 1 1 1 1	9823	22	—	162	13	—	4.10	0.36	—	0.297	0.016	—	31.9	B
535	P445-1 I	160	1 1 1 1 1	11125	30	—	132	44	—	2.44	1.62	—	0.232	0.047	—	6.6	B
536	P445-2 A	53	1 1 1 1 1	11106	25	—	343	35	—	5.24	0.31	—	0.334	0.016	—	40.0	A
537	P445-2 B	53	1 1 1 1 1	10388	24	—	225	20	—	5.18	0.34	—	0.320	0.016	—	36.2	A
538	P445-2 C	61	1 1 1 1 1	4742	21	—	185	10	—	4.39	0.27	—	0.251	0.014	—	48.6	A
539	P445-2 D	61	1 1 1 1 1	4742	21	—	137	7	—	4.13	0.26	—	0.227	0.014	—	50.4	A
540	A2162-N A	62	1 1 1 1 1	14847	22	—	172	16	—	6.16	0.65	—	0.345	0.024	—	19.8	A
541	A2162-N B	62	1 1 1 1 0	14395	28	—	222	25	—	5.00	0.43	—	—	—	—	21.2	B
542	A2162-N C	62	2 2 2 2 1	15172	20	14	201	18	19	5.71	0.43	0.36	0.306	0.040	—	26.4	B
543	A2162-N D	62	1 1 1 1 1	15367	31	—	234	24	—	5.01	0.76	—	0.326	0.025	—	18.6	B
545	A2162-N I	62	1 1 1 1 1	14811	53	—	301	100	—	4.89	1.12	—	0.298	0.036	—	9.4	A
546	A2162-S A	63	2 2 2 2 2	9602	28	16	288	31	31	5.53	0.41	0.24	0.360	0.015	0.006	34.6	A
547	A2162-S B	163	1 1 1 1 1	14994	85	—	487	142	—	4.45	1.11	—	0.354	0.032	—	15.3	B
548	A2162-S C	163	2 2 2 2 2	14905	18	5	256	13	2	5.17	0.24	0.05	0.297	0.013	0.005	49.2	A
549	A2162-S D	163	1 1 1 1 0	16237	18	—	71	14	—	4.95	0.72	—	—	—	—	14.3	C*
550	A2162-S E	163	2 2 2 2 1	16086	14	7	183	10	11	4.48	0.24	0.04	0.286	0.014	—	43.0	B
552	A2162-S G	63	1 1 1 1 1	9613	84	—	448	196	—	4.77	1.30	—	0.281	0.045	—	7.4	A
554	A2162-S I	163	1 1 1 1 1	15593	40	—	223	61	—	4.27	1.14	—	0.273	0.036	—	9.6	A
555	A2162-S 1	63	1 1 1 1 1	9883	32	—	147	45	—	3.45	1.13	—	0.190	0.039	—	7.5	B
556	A2162-S 2	263	1 1 0 0 0	27486	73	—	—	—	—	—	—	—	—	—	—	3.4	D
557	J20 A	64	1 1 1 1 0	9429	31	—	266	43	—	5.14	0.59	—	—	—	—	21.3	A
558	J20 B	64	2 2 2 2 0	8921	16	17	204	21	12	4.56	0.33	0.26	—	—	—	70.2	A
559	J20 C	64	1 1 1 1 0	9431	23	—	283	37	—	5.53	0.33	—	—	—	—	77.4	A
560	J20 D	64	1 1 1 1 0	8695	22	—	192	39	—	5.11	0.33	—	—	—	—	67.1	B

Table 5 – *continued*

(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)	(11)	(12)	(13)	(14)	(15)	(16)	(17)	(18)
GIN	Galaxy Name	CAN	N $s z \sigma b 2$	cz (km s^{-1})	Δcz	δcz	σ (km s^{-1})	$\Delta \sigma$	$\delta \sigma$	Mgb	ΔMgb (\AA)	δMgb	Mg_2	ΔMg_2 (mag)	δMg_2	S/N	Q
561	J20 E	64	1 1 1 1 0	8836	21	—	206	13	—	5.31	0.40	—	—	—	—	49.7	A
562	A2197 A	65	4 3 3 3 3	8784	15	12	257	18	5	4.87	0.32	0.05	0.313	0.011	0.002	49.7	A
563	A2197 B	65	2 2 2 2 2	9523	21	0	210	21	22	5.27	0.47	0.28	0.331	0.015	0.000	33.9	A
564	A2197 C	65	3 3 3 3 2	8820	16	7	276	18	8	4.48	0.28	0.04	0.284	0.013	0.000	59.5	A
565	A2197 D	65	1 1 1 1 1	8970	19	—	110	17	—	3.11	0.64	—	0.119	0.025	—	14.5	C*
566	A2197 E	65	4 3 3 3 3	8922	13	60	261	13	13	5.46	0.24	0.36	0.310	0.010	0.010	52.2	A
568	A2197 G	65	1 1 1 1 1	8248	24	—	156	29	—	4.69	0.78	—	0.264	0.029	—	12.2	C
571	A2197 J	65	3 3 3 3 3	9092	13	15	154	8	13	3.81	0.25	0.09	0.247	0.011	0.008	38.2	A
572	A2197 K	65	1 1 1 1 1	9553	33	—	168	47	—	3.82	1.24	—	0.243	0.036	—	8.7	B
573	A2197 L	65	1 1 1 1 1	9673	20	—	192	12	—	4.54	0.30	—	0.278	0.014	—	37.7	A
574	A2197 M	65	1 1 1 1 1	9524	34	—	177	47	—	3.72	0.98	—	0.244	0.035	—	9.7	B
578	A2197 3	65	1 1 1 1 1	10344	34	—	168	47	—	4.37	1.11	—	0.266	0.037	—	8.6	A
579	A2199 A	66	1 1 1 1 1	9303	30	—	281	39	—	5.67	0.83	—	0.365	0.029	—	17.3	B
580	A2199 B	66	1 1 1 1 1	8800	21	—	158	30	—	4.98	0.74	—	0.301	0.028	—	16.3	B
581	A2199 C	66	1 1 1 1 1	7834	21	—	160	11	—	4.59	0.29	—	0.272	0.015	—	35.9	A
582	A2199 D	66	1 1 1 1 1	8716	21	—	194	12	—	4.80	0.27	—	0.309	0.015	—	41.6	A
583	A2199 E	66	1 1 1 1 1	9517	21	—	117	8	—	3.48	0.29	—	0.188	0.015	—	39.9	B*
584	A2199 F	66	2 2 2 2 2	8652	14	4	192	9	0	4.39	0.22	0.09	0.288	0.010	0.034	55.8	A
585	A2199 G	66	1 1 1 1 1	10701	21	—	109	8	—	2.92	0.32	—	0.183	0.017	—	32.6	B*
586	A2199 H	66	1 1 1 1 1	8856	23	—	189	15	—	4.78	0.32	—	0.331	0.016	—	34.0	A
588	A2199 J	66	1 1 1 1 1	9092	24	—	171	16	—	4.57	0.32	—	0.298	0.016	—	27.4	A
589	A2199 K	66	1 1 1 1 1	9127	21	—	168	12	—	5.09	0.31	—	0.292	0.016	—	36.1	A
590	A2199 L	66	1 1 1 1 1	8897	21	—	166	12	—	4.47	0.32	—	0.308	0.015	—	36.4	A
591	A2199 M	66	1 1 1 1 1	10108	21	—	181	12	—	3.89	0.33	—	0.247	0.016	—	38.5	A
592	A2199 N	66	1 1 1 1 1	9918	21	—	144	11	—	2.70	0.36	—	0.178	0.016	—	33.4	B*
593	A2199 O	66	1 1 1 1 1	9006	22	—	192	14	—	5.11	0.31	—	0.355	0.015	—	37.2	A
595	A2199 1	66	2 2 2 2 1	8814	19	1	193	14	0	4.10	0.43	0.24	0.229	0.049	—	39.6	A
596	A2199 2	166	1 1 1 1 1	17795	46	—	256	77	—	4.18	1.31	—	0.288	0.035	—	9.0	B
598	J21 A	167	2 2 2 2 0	6085	16	5	200	21	8	4.61	0.30	0.53	—	—	—	49.0	A
600	J21 C	67	1 1 1 1 1	13119	33	—	287	47	—	5.12	0.51	—	0.314	0.020	—	22.9	A
601	J21 D	67	1 1 1 1 0	13778	22	—	198	17	—	6.14	0.37	—	—	—	—	26.7	A
602	J21 E	67	2 2 2 2 1	14348	21	20	327	22	22	5.13	0.23	0.04	0.333	0.023	—	52.8	A
604	J21 G	67	1 1 1 1 1	14201	21	—	178	96	—	3.38	0.57	—	0.179	0.040	—	30.3	C
606	J21 I	67	1 1 1 1 1	13259	24	—	222	20	—	5.65	0.39	—	0.294	0.015	—	27.3	B
607	J21 1	67	1 1 1 1 0	13750	23	—	197	18	—	5.70	0.41	—	—	—	—	25.9	B
608	J21 2	67	1 1 1 1 0	14675	23	—	172	17	—	4.61	0.46	—	—	—	—	20.8	B
609	A2247 A	68	1 1 1 1 0	11111	22	—	209	33	—	2.81	0.38	—	—	—	—	63.6	A
610	A2247 B	68	1 1 1 1 0	11844	22	—	229	34	—	4.40	0.37	—	—	—	—	57.0	A
611	A2247 C	68	1 1 1 1 0	11796	22	—	226	37	—	5.67	0.35	—	—	—	—	49.8	A
612	A2247 D	68	1 1 1 1 0	11251	22	—	214	36	—	4.05	0.38	—	—	—	—	44.4	A
613	A2247 E	68	1 1 1 1 0	11366	22	—	223	36	—	4.52	0.40	—	—	—	—	47.3	A
614	A2247 F	68	1 1 1 1 1	11679	24	—	270	20	—	5.00	0.28	—	0.312	0.014	—	37.0	A
615	A2247 1	68	1 1 1 1 1	11696	32	—	136	69	—	2.97	2.98	—	0.239	0.077	—	3.4	C
616	A2247 2	68	1 1 1 1 1	11371	17	—	97	25	—	4.39	0.45	—	0.246	0.017	—	24.9	B
617	P332-1 A	70	2 2 2 2 1	10675	14	15	204	11	0	4.79	0.21	0.10	0.281	0.013	—	89.1	A
618	P332-1 B	70	1 1 1 1 0	10785	23	—	260	38	—	2.22	0.41	—	—	—	—	51.5	B
619	P332-1 C	70	2 2 2 2 1	10066	14	34	164	17	5	3.95	0.27	0.06	0.221	0.019	—	62.1	A
620	P332-1 D	69	1 1 1 1 1	16941	24	—	204	33	—	1.19	0.79	—	0.079	0.029	—	14.4	C
621	P332-1 E	69	1 1 1 1 1	17747	24	—	234	25	—	4.21	0.41	—	0.265	0.017	—	25.8	B
622	P332-1 F	70	1 1 1 1 1	10522	22	—	177	20	—	3.82	0.54	—	0.247	0.022	—	21.6	B
623	P332-1 G	70	1 1 1 1 1	10042	47	—	284	84	—	4.84	0.89	—	0.288	0.037	—	11.0	A
626	J22 A	70	3 3 3 3 3	10448	24	37	330	32	17	5.38	0.28	0.14	0.339	0.011	0.019	40.9	A
627	J22 B	70	2 2 2 2 2	10358	30	20	296	33	29	5.41	0.34	0.68	0.309	0.013	0.021	34.5	A
628	J22 C	70	2 2 2 2 2	9729	21	1	222	23	15	5.37	0.30	0.60	0.315	0.013	0.019	37.2	A
629	J22 D	70	1 1 1 1 1	9723	34	—	243	38	—	5.95	0.33	—	0.296	0.016	—	26.3	B
630	J22 E	70	2 2 2 2 2	10542	20	10	233	24	9	4.40	0.34	0.18	0.281	0.012	0.003	34.8	A
631	J22 F	70	1 1 1 1 1	10952	28	—	141	32	—	3.79	0.76	—	0.220	0.025	—	15.0	B
632	J22 G	70	2 2 2 2 2	9949	21	20	273	20	18	5.58	0.30	0.27	0.318	0.011	0.009	36.6	A
633	J22 H	70	2 2 2 2 2	10376	14	34	172	9	10	4.49	0.24	0.26	0.253	0.011	0.009	49.3	A
634	J23 A	71	2 2 2 2 1	8220	26	3	347	40	64	5.44	0.44	0.18	0.377	0.019	—	30.1	A
635	J23 B	71	2 2 2 2 2	8867	19	18	245	25	3	4.01	0.49	0.05	0.292	0.015	0.003	26.1	B
636	J23 C	71	2 2 2 2 2	8496	20	69	185	29	29	4.63	0.63	0.08	0.258	0.021	0.003	18.7	B

Table 5 – continued

(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)	(11)	(12)	(13)	(14)	(15)	(16)	(17)	(18)
GIN	Galaxy Name	CAN	N $s z \sigma b 2$	cz (km s^{-1})	Δcz	δcz	σ (km s^{-1})	$\Delta \sigma$	$\delta \sigma$	Mgb	ΔMgb (\AA)	δMgb	M_{g_2}	ΔM_{g_2} (mag)	δM_{g_2}	S/N	Q
638	J23 E	71	2 2 2 2 1	8275	20	15	147	24	4	3.42	0.64	0.31	0.206	0.040	—	18.7	B
639	J24 A	72	2 2 2 2 2	9858	27	46	253	37	30	5.58	0.60	0.91	0.324	0.018	0.004	19.7	B
640	J24 B	72	2 2 2 2 2	10228	16	11	103	22	20	3.69	0.63	0.25	0.248	0.021	0.002	17.8	B
641	J24 C	72	3 3 2 2 2	10346	22	45	208	40	3	4.72	0.70	0.36	0.290	0.025	0.007	14.6	A
642	J24 D	72	1 1 0 0 0	10695	72	—	—	—	—	—	—	—	—	—	—	10.6	D
643	J24 E	72	1 1 1 1 1	10427	25	—	195	37	—	4.09	1.02	—	0.243	0.030	—	13.9	B
645	J25 A	73	2 2 2 2 1	8233	23	12	204	27	12	5.14	0.54	0.73	0.381	0.027	—	25.0	A
646	J25 B	73	2 2 2 2 2	8354	21	11	206	24	4	5.35	0.42	0.21	0.329	0.017	0.000	27.8	A
647	J25 C	73	3 3 2 2 2	7775	18	46	140	24	33	4.30	0.76	0.00	0.306	0.025	0.019	16.4	A
648	J26 A	74	3 2 2 2 2	15045	20	43	269	30	20	6.31	0.44	1.06	0.366	0.015	0.006	26.8	B
649	J26 B	174	1 1 1 1 1	18187	37	—	284	65	—	6.40	1.38	—	0.286	0.043	—	9.8	C
650	J26 C	74	2 2 2 2 2	15292	20	15	281	41	2	4.30	0.43	0.11	0.315	0.019	0.001	35.5	B
651	J26 1	74	1 1 0 0 0	15341	71	—	—	—	—	—	—	—	—	—	—	15.5	D
652	J26 2	74	1 1 1 1 1	14188	19	—	134	35	—	5.11	0.72	—	0.269	0.027	—	14.8	C
653	J27 A	75	3 2 2 2 2	10646	30	13	289	46	31	5.25	0.55	0.23	0.340	0.020	0.001	22.7	A
654	J27 B	75	3 2 2 2 1	9325	29	0	235	40	5	5.68	0.64	0.62	0.328	0.032	—	20.2	B
655	J27 C	75	2 1 1 1 1	9784	32	—	221	43	—	4.72	0.60	—	0.248	0.024	—	17.6	B
656	J27 D	75	2 1 1 1 1	9424	20	—	95	21	—	4.05	0.81	—	0.279	0.028	—	16.5	B
657	J27 E	75	3 2 2 2 2	9740	15	17	192	22	32	4.91	0.48	0.16	0.278	0.016	0.016	27.2	B
658	J38 A	76	3 3 3 3 3	15519	31	45	385	42	7	4.48	0.28	0.06	0.335	0.012	0.026	34.6	A
659	J38 B	76	3 3 3 3 3	14818	16	16	190	21	23	4.09	0.44	1.46	0.248	0.015	0.014	25.2	B
660	J38 C	76	5 5 2 2 2	16207	19	19	125	21	25	4.29	0.56	0.20	0.206	0.022	0.008	17.6	C
661	J38 1	76	1 1 1 1 1	14918	26	—	108	43	—	0.56	1.14	—	0.088	0.041	—	9.7	C
662	P522-1 A	77	7 6 6 6 4	7318	10	18	288	6	11	5.22	0.09	0.33	0.331	0.005	0.004	147.2	A
663	P522-1 B	77	6 5 5 5 3	7847	10	11	249	7	14	5.35	0.10	0.18	0.328	0.006	0.014	119.3	A
664	P522-1 C	77	4 4 4 4 2	7697	10	5	101	6	7	4.11	0.16	0.26	0.246	0.008	0.001	74.1	A
665	P522-1 D	77	5 5 5 5 3	7933	8	7	184	6	16	4.43	0.11	0.35	0.266	0.006	0.002	107.6	A
666	P522-1 E	77	5 5 5 5 3	7315	10	9	166	8	12	4.54	0.18	0.27	0.274	0.009	0.006	61.9	A
667	P522-1 F	77	4 4 4 4 3	8340	10	39	134	6	10	4.38	0.16	0.19	0.255	0.008	0.004	74.0	A
668	P522-1 G	77	4 4 3 3 2	7118	11	22	191	6	8	4.45	0.12	0.15	0.267	0.007	0.007	111.9	A
669	P522-1 1	77	3 3 3 3 2	7572	11	2	186	6	8	4.60	0.13	0.21	0.292	0.007	0.000	88.5	A
670	P522-1 2	177	2 2 2 2 2	23926	13	12	216	11	5	3.97	0.23	0.12	0.266	0.009	0.004	51.8	A
672	P522-1 4	77	3 3 3 3 3	7607	10	8	85	6	17	3.83	0.19	0.10	0.200	0.009	0.010	56.4	A
673	P522-1 5	277	2 2 2 2 2	27282	13	31	119	12	42	5.17	0.33	0.25	0.234	0.012	0.001	35.8	B
674	A2572 A	78	2 2 2 2 2	11793	20	12	305	21	9	5.11	0.21	0.68	0.300	0.009	0.015	45.0	A
675	A2572 B	78	3 3 3 3 3	11384	15	16	263	17	20	5.10	0.24	0.07	0.320	0.009	0.013	45.9	A
676	A2572 C	78	2 2 2 2 2	11263	23	34	296	24	15	5.38	0.32	0.26	0.328	0.012	0.003	31.9	A
677	A2572 D	78	2 2 2 2 2	11260	16	12	182	12	16	5.63	0.23	0.65	0.315	0.011	0.016	37.6	B
678	A2572 E	78	2 2 0 0 0	11297	50	54	—	—	—	—	—	—	—	—	—	20.5	D*
680	A2572 2	78	2 2 2 2 2	12027	14	7	149	16	8	3.22	0.45	0.96	0.290	0.017	0.042	23.7	B
681	A2589 A	79	6 5 5 5 5	12327	10	7	254	9	7	5.45	0.13	0.16	0.345	0.006	0.014	85.7	A
682	A2589 B	79	5 5 5 5 5	12773	11	35	261	10	16	4.80	0.14	0.25	0.301	0.006	0.007	81.1	A
683	A2589 C	79	3 3 3 3 3	14145	10	12	165	8	31	4.42	0.17	0.13	0.272	0.007	0.004	65.4	A
684	A2589 D	79	3 3 3 3 3	13549	11	15	225	9	11	5.08	0.18	0.30	0.306	0.007	0.012	63.7	A
685	A2589 E	79	3 3 3 3 3	10403	9	24	70	8	12	2.53	0.27	0.34	0.103	0.010	0.002	40.8	B
686	A2589 F	79	3 3 3 3 3	12124	12	2	254	10	11	5.17	0.13	0.07	0.284	0.006	0.007	77.0	A
687	A2589 G	79	3 3 3 3 3	12535	12	30	183	10	20	5.33	0.21	0.20	0.277	0.009	0.014	45.0	A
688	A2589 H	79	3 3 3 3 3	13420	13	12	244	11	21	4.63	0.18	0.31	0.269	0.008	0.003	56.2	A
689	A2589 1	79	3 3 3 3 3	13329	11	26	121	9	3	4.75	0.24	0.13	0.248	0.010	0.003	45.9	B
690	A2589 2	79	3 3 2 2 2	12312	13	32	113	17	38	4.85	0.53	0.57	0.233	0.018	0.006	21.4	C
691	A2593-N A	80	2 2 2 2 2	12489	17	3	246	13	4	5.62	0.15	0.10	0.345	0.009	0.015	49.9	A
692	A2593-N B	80	6 6 6 6 4	11601	8	18	208	8	18	4.83	0.14	0.29	0.321	0.007	0.029	74.1	A
693	A2593-N C	80	3 3 3 3 3	13381	10	4	204	8	18	4.61	0.14	0.24	0.293	0.007	0.012	74.0	A
694	A2593-N D	80	4 4 4 4 4	12637	9	13	231	6	15	4.83	0.11	0.23	0.308	0.005	0.014	91.8	A
695	A2593-N E	80	4 4 4 4 4	11591	11	8	311	10	9	5.18	0.10	0.13	0.330	0.006	0.014	103.3	A
696	A2593-N F	80	4 4 4 4 4	12424	9	6	165	6	10	4.99	0.11	0.07	0.294	0.006	0.003	82.4	A
697	A2593-N G	80	5 5 4 4 3	11747	9	22	142	7	15	4.46	0.18	0.27	0.240	0.008	0.023	55.6	A
698	A2593-N H	80	4 4 4 4 4	12645	10	12	239	7	9	5.37	0.11	0.23	0.325	0.006	0.024	88.1	A
699	A2593-N I	80	5 5 5 5 3	12263	10	14	256	8	10	5.06	0.13	0.23	0.318	0.006	0.012	84.2	A
700	A2593-N J	80	7 7 7 7 5	12617	8	12	235	6	18	5.46	0.11	0.39	0.319	0.005	0.019	101.7	A
701	A2593-N K	80	7 7 7 7 5	12634	8	17	238	6	6	5.03	0.10	0.22	0.337	0.005	0.009	109.8	A
702	A2593-N L	80	5 5 5 5 3	11823	8	16	202	6	11	5.03	0.13	0.28	0.309	0.006	0.014	84.0	A

Table 5 – continued

(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)	(11)	(12)	(13)	(14)	(15)	(16)	(17)	(18)
GIN	Galaxy Name	CAN	N $s z \sigma b 2$	cz (km s^{-1})	Δcz	δcz	σ (km s^{-1})	$\Delta \sigma$	$\delta \sigma$	Mgb	ΔMgb (\AA)	δMgb	M_{g_2}	ΔM_{g_2} (mag)	δM_{g_2}	S/N	Q
703	A2593-N M	80	5 5 5 5 3	13189	8	24	182	7	12	4.62	0.14	0.27	0.288	0.007	0.018	74.3	A
704	A2593-N 1	80	3 3 3 3 3	11704	11	19	240	8	11	4.73	0.15	0.24	0.290	0.007	0.011	76.4	A
705	A2593-N 2	80	3 3 3 3 3	13110	10	18	170	7	10	4.52	0.17	0.26	0.298	0.008	0.014	63.5	A
706	A2593-N 3	180	7 7 6 4 4	27822	8	28	229	7	29	4.90	0.15	0.21	0.325	0.007	0.009	84.0	A
707	A2593-N 4	80	5 5 3 3 3	11781	10	20	125	6	10	4.29	0.19	0.35	0.256	0.008	0.013	57.4	A
708	A2593-N 5	80	5 5 2 2 2	11096	11	28	57	9	3	3.65	0.36	0.24	0.175	0.012	0.002	36.3	B
709	A2593-S A	80	5 5 5 5 3	12713	11	19	330	11	39	5.23	0.10	0.23	0.326	0.006	0.008	100.9	A
710	A2593-S B	80	5 5 5 5 3	12150	8	15	178	7	8	4.86	0.12	0.20	0.279	0.007	0.031	78.4	A
711	A2593-S C	80	3 3 3 3 3	13066	12	41	293	12	8	5.18	0.14	0.11	0.317	0.007	0.010	77.6	A
712	A2593-S D	80	5 5 5 5 3	11046	9	21	214	8	6	4.75	0.12	0.18	0.280	0.006	0.020	85.0	A
713	A2593-S E	80	7 5 3 3 3	12554	10	9	143	7	7	4.19	0.19	0.60	0.225	0.009	0.024	57.4	A
714	A2593-S F	80	3 3 3 3 3	12805	10	10	124	8	16	4.59	0.16	0.18	0.253	0.008	0.011	56.8	A
715	A2593-S G	80	3 3 3 3 3	12880	10	7	107	9	8	3.68	0.25	0.29	0.201	0.010	0.011	42.6	B
716	A2634 A	82	3 3 3 3 2	9060	14	11	312	11	10	5.40	0.08	0.06	0.329	0.006	0.011	105.6	A
717	A2634 B	82	3 3 3 3 3	9274	12	10	236	8	9	4.98	0.08	0.11	0.295	0.006	0.003	90.1	A
718	A2634 C	82	3 3 3 3 3	9502	12	23	190	7	20	3.65	0.12	1.03	0.297	0.007	0.022	69.4	A
719	A2634 D	82	3 2 2 2 2	10774	15	1	231	12	15	4.84	0.16	0.02	0.297	0.009	0.005	45.4	A
720	A2634 E	82	3 3 3 3 3	9272	14	11	269	13	7	4.85	0.16	0.11	0.316	0.007	0.006	54.7	A
721	A2634 F	82	2 2 2 2 2	10211	13	16	135	10	4	4.75	0.16	0.03	0.288	0.009	0.014	41.3	A
722	A2634 G	82	4 4 4 4 4	9398	15	16	333	17	7	5.06	0.13	0.31	0.327	0.007	0.018	62.8	A
723	A2634 H	82	3 3 3 3 3	11694	15	21	271	19	24	4.60	0.18	0.24	0.266	0.008	0.006	48.5	A
724	A2634 I	82	3 3 3 3 3	9958	13	25	211	12	8	4.84	0.18	0.10	0.297	0.009	0.013	43.2	A
725	A2634 J	82	2 2 2 2 2	8701	17	25	219	16	15	5.30	0.21	0.55	0.309	0.011	0.016	35.3	A
726	A2634 K	82	3 3 3 3 3	10011	13	55	214	12	8	5.28	0.15	0.22	0.333	0.008	0.008	46.5	A
728	A2634 2	82	1 1 1 1 1	9406	31	—	296	35	—	5.76	0.41	—	0.295	0.018	—	24.1	B
729	A2657 A	83	2 2 2 2 1	12063	14	15	179	18	19	4.84	0.30	0.24	0.312	0.020	—	24.7	A
730	A2657 B	83	2 2 2 2 2	12354	18	9	281	20	14	4.86	0.20	0.08	0.271	0.010	0.032	41.9	B
731	A2657 C	83	3 3 3 3 2	12298	12	5	196	13	9	4.25	0.20	0.14	0.254	0.011	0.006	37.8	A
732	A2657 D	83	1 1 1 1 1	12542	19	—	221	19	—	4.14	0.29	—	0.263	0.012	—	27.0	B
733	A2657 E	83	2 2 2 2 1	10725	15	17	215	14	9	5.12	0.20	0.29	0.295	0.012	—	35.8	A
734	A2657 F	83	2 2 2 2 1	11876	15	9	239	15	24	5.72	0.20	0.08	0.351	0.011	—	41.8	A
735	A2657 G	83	2 2 2 2 1	13594	12	17	150	11	7	4.18	0.22	0.13	0.224	0.014	—	33.5	A
736	A2657 H	83	2 2 2 2 1	12249	14	0	219	13	3	5.05	0.19	0.23	0.321	0.012	—	40.7	A
737	A2657 1	83	1 1 1 1 0	11719	20	—	152	25	—	4.73	0.46	—	—	—	—	18.4	B
738	A2666 A	84	2 2 2 2 1	8153	17	3	257	19	15	5.24	0.28	0.19	0.311	0.013	—	35.5	A
739	A2666 B	84	2 2 2 2 1	7649	14	64	84	20	1	3.43	0.66	0.30	0.185	0.030	—	12.8	C
740	A2666 C	84	2 2 1 1 1	8511	17	8	93	11	—	4.53	0.39	—	0.259	0.015	—	22.8	B
741	A2666 D	84	1 1 0 0 0	9018	71	—	—	—	—	—	—	—	—	—	—	19.4	D
742	A2666 E	84	1 1 1 1 1	8391	19	—	181	14	—	4.31	0.35	—	0.277	0.014	—	27.4	B
750	COMA 70	90	1 1 1 1 1	6368	21	—	139	13	—	5.02	0.44	—	0.295	0.018	—	21.3	B
751	COMA 69	90	1 1 1 1 1	7065	29	—	206	25	—	5.37	0.43	—	0.333	0.018	—	18.9	B
752	COMA 240	90	1 1 1 1 1	6761	24	—	286	24	—	5.29	0.28	—	0.287	0.015	—	44.1	A
753	COMA 239	90	1 1 1 1 1	6266	22	—	196	14	—	4.76	0.29	—	0.257	0.015	—	37.4	A
754	COMA 194	90	1 1 1 1 0	7963	36	—	276	45	—	4.54	0.44	—	—	—	—	28.5	A
755	COMA 159	90	2 2 2 2 2	6811	15	1	179	9	18	4.47	0.25	0.11	0.261	0.011	0.009	48.7	A
756	COMA 133	90	2 2 2 2 2	4829	15	8	208	11	12	4.82	0.24	0.41	0.262	0.011	0.000	47.5	A
758	COMA 31	90	1 1 1 1 0	7376	33	—	204	36	—	4.30	0.55	—	—	—	—	20.0	A
759	COMA 49	90	1 1 1 1 1	7853	24	—	293	18	—	5.52	0.24	—	0.308	0.014	—	57.9	A
761	COMA 103	90	1 1 1 1 1	4752	16	—	205	17	—	3.94	0.26	—	0.274	0.011	—	46.1	A
762	COMA 106	90	1 1 1 1 1	5132	21	—	151	14	—	4.85	0.40	—	0.252	0.015	—	21.9	B
763	COMA 107	90	1 1 1 1 1	6499	21	—	133	20	—	4.89	0.74	—	0.229	0.025	—	13.9	C
764	COMA 118	90	1 1 1 1 1	7497	21	—	170	11	—	4.83	0.26	—	0.257	0.015	—	41.4	A
765	COMA 129	90	3 3 3 3 2	7199	21	13	250	20	21	4.64	0.24	0.20	0.285	0.011	0.011	40.6	B
766	COMA 134	90	1 1 1 1 1	7021	21	—	141	13	—	5.23	0.41	—	0.260	0.016	—	21.7	B
767	COMA 135	90	2 2 2 2 2	8312	14	15	117	15	2	4.38	0.54	0.29	0.215	0.019	0.025	17.7	C
768	COMA 136	90	1 1 1 1 1	5706	24	—	181	18	—	4.86	0.51	—	0.253	0.018	—	20.5	B
769	COMA 148	90	3 3 3 3 3	6451	20	31	356	20	11	5.45	0.18	0.14	0.333	0.008	0.009	62.0	A
770	COMA 160	90	1 1 1 1 0	7673	20	—	148	10	—	3.88	0.43	—	—	—	—	50.3	A
771	COMA 172	90	1 1 1 1 1	5734	34	—	171	48	—	3.96	1.29	—	0.212	0.039	—	8.4	A
772	COMA 217	90	1 1 1 1 1	6713	30	—	200	41	—	4.23	0.84	—	0.271	0.031	—	14.5	A
773	N3379	95	2 2 2 2 0	908	12	10	185	7	4	4.76	0.21	0.11	—	—	—	78.1	A
774	N 936	95	3 3 3 3 3	1425	9	17	150	6	2	4.67	0.11	0.03	0.330	0.006	0.019	102.7	A

Table 5 – continued

(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)	(11)	(12)	(13)	(14)	(15)	(16)	(17)	(18)
GIN	Galaxy Name	CAN	N $s z \sigma b 2$	cz (km s^{-1})	Δcz	δcz	σ (km s^{-1})	$\Delta \sigma$	$\delta \sigma$	Mgb	ΔMgb (\AA)	δMgb	Mg_2	ΔMg_2 (mag)	δMg_2	S/N	Q
775	N3115	95	1 1 1 1 0	732	20	—	239	14	—	4.44	0.34	—	—	—	—	50.8	A
776	N3377	95	1 1 1 1 1	643	18	—	116	15	—	3.72	0.40	—	0.289	0.017	—	52.5	A
777	N4365	95	1 1 1 1 0	1252	19	—	236	16	—	4.95	0.38	—	—	—	—	43.9	A
778	N4374	91	1 1 1 1 0	1028	18	—	236	13	—	4.45	0.35	—	—	—	—	53.9	A
779	N4382	91	1 1 1 1 0	719	18	—	134	11	—	2.88	0.26	—	—	—	—	59.1	A
780	N4406	91	1 1 1 1 0	−209	20	—	246	17	—	5.45	0.34	—	—	—	—	43.4	A
781	N4472	91	1 1 1 1 0	984	18	—	260	11	—	4.69	0.34	—	—	—	—	57.7	A
782	N4473	91	2 2 2 2 1	2250	13	5	164	6	10	5.03	0.19	0.06	0.311	0.013	—	103.5	A
783	N4486B	91	1 1 1 1 0	1494	21	—	158	17	—	3.40	0.44	—	—	—	—	29.2	A
784	N4494	95	1 1 1 1 0	1351	17	—	145	10	—	4.20	0.37	—	—	—	—	44.3	A
785	COMA 130	90	1 1 1 1 1	7247	29	—	236	26	—	4.70	0.49	—	0.258	0.016	—	23.0	B
786	COMA 153	90	2 1 1 1 0	6693	21	—	149	12	—	3.92	0.44	—	—	—	—	38.6	B
787	COMA 124	90	1 1 1 1 1	6705	22	—	204	16	—	4.58	0.36	—	0.260	0.016	—	29.0	B
788	COMA 125	90	1 1 1 1 1	6926	23	—	186	18	—	4.16	0.43	—	0.207	0.015	—	23.7	B
789	COMA 143	90	1 1 1 1 1	5000	38	—	219	57	—	4.68	0.95	—	0.294	0.033	—	10.9	A
790	COMA 168	90	1 1 1 1 1	7044	41	—	244	66	—	4.51	1.07	—	0.269	0.033	—	10.7	A
791	COMA 151	90	4 4 4 4 4	6379	10	16	152	7	7	4.85	0.21	0.66	0.256	0.009	0.033	47.0	A
792	COMA 150	90	2 2 2 2 2	7258	14	12	112	12	0	5.30	0.37	0.37	0.268	0.015	0.027	22.6	B
794	N 584	95	2 2 2 2 1	1854	11	8	179	6	3	4.28	0.11	0.25	0.337	0.008	—	113.2	A
795	N 596	95	2 2 2 2 1	1876	11	8	146	6	9	3.96	0.13	0.04	0.293	0.009	—	96.0	A
796	N4697	95	1 1 1 1 0	1237	19	—	164	7	—	3.96	0.40	—	—	—	—	90.8	A
797	N5846	95	1 1 1 1 1	1724	32	—	221	47	—	4.43	0.86	—	0.307	0.030	—	14.4	A
799	N7626	95	4 4 4 4 3	3403	10	35	244	8	8	5.07	0.13	0.29	0.338	0.007	0.018	80.9	A
800	N4486	91	1 1 1 1 0	1290	22	—	321	14	—	5.04	0.39	—	—	—	—	97.9	A

The columns of this table give: (1) galaxy identification number (GIN); (2) galaxy name; (3) the cluster assignment number (CAN); (4) the number of spectra N_s , redshifts N_z , dispersions N_σ , Mgb linestrengths N_b and Mg_2 linestrengths N_2 obtained for this object; then the combined estimate, its estimated total error (Δ) and the weighted rms error from any repeat observations (δ) for each of (5)–(7) redshift, (8)–(10) dispersion, (11)–(13) Mgb linestrength and (14)–(16) Mg_2 linestrength; (17) the combined S/N estimate; and (18) the overall quality parameter (with an asterisk if the galaxy possesses emission lines). Only objects with useful measurements are included; hence the lowest quality class present in this table is $Q = D$, and the seven galaxies with only $Q = E$ spectra (GINs 123, 284, 389, 448, 599, 637 and 679) are omitted. The data in this table are available from NASA’s Astrophysical Data Center (ADC, <http://adc.gsfc.nasa.gov>, J/MNRAS/305/259) and from the Centre de Données astronomiques de Strasbourg (CDS, <http://cdsweb.u-strasbg.fr>, J/MNRAS/305/259).

We determine the zero-point calibration of our linestrength measurements with respect to the Lick system (see Section 3.2) by comparing our Mgb' and Mg_2 linestrengths with measurements for the same galaxies given by Trager et al. (1998). We find that slightly different calibrations are needed for objects with different redshifts, the result of slight variations in the non-linear continuum shape as the spectra are redshifted with respect to the instrumental response and the sky background (see Section 3.2). Good agreement with Trager et al. is obtained if we use different zero-points for galaxies with redshifts above and below $cz = 3000 \text{ km s}^{-1}$ (although there are no objects in the comparison at $cz > 10000 \text{ km s}^{-1}$). Excluding a few outliers, we find weighted mean differences between the EFAR and Trager et al. linestrengths of $\langle \Delta \text{Mgb}' \rangle = -0.022 \text{ mag}$ and $\langle \Delta \text{Mg}_2 \rangle = -0.083 \text{ mag}$ for $cz < 3000 \text{ km s}^{-1}$, and $\langle \Delta \text{Mgb}' \rangle = -0.008 \text{ mag}$ and $\langle \Delta \text{Mg}_2 \rangle = -0.028 \text{ mag}$ for $cz \geq 3000 \text{ km s}^{-1}$. Subtracting these zero-point corrections gives the final, fully corrected linestrength measurements as listed in Tables 3 and 5. Figs 14 and 15 show the residual differences between the EFAR and Trager et al. linestrength measurements after applying these zero-point corrections. The rms scatter is 0.019 mag in Mgb' for the 41 objects in common, and 0.023 mag in Mg_2 for the 24 objects in common. There is no statistically significant trend with linestrength, velocity dispersion or redshift remaining in the residuals after these zero-point corrections are applied.

Fig. 16 compares our calibrated Mg_2 linestrengths with those obtained in A2199, A2634 and Coma by Lucey et al. (1997). The overall agreement for the 36 objects in common is very good, with a statistically non-significant zero-point offset and an rms scatter of 0.029 mag, similar to that found in the comparison with Trager et al.

The relation between the measured Mgb' and Mg_2 linestrengths for all the galaxies in the EFAR sample is shown in Fig. 17. We fit this relation using a maximum likelihood technique which accounts for both measurement errors and selection effects (Saglia et al., in preparation; Paper VI). We find

$$\text{Mg}_2 = 1.94 \text{Mgb}' - 0.05, \quad (13)$$

with a perpendicular rms residual of 0.019 mag (corresponding to an rms of 0.041 mag in Mg_2 , or 0.021 mag in Mgb'). The relation is the same if we fit ellipticals, E/S0s, cDs or spirals separately. This relation is similar to those derived by Burstein et al. (1984) and Jørgensen (1997). We can therefore use Mgb' as a predictor of Mg_2 (albeit with larger uncertainties) for those cases in which Mg_2 cannot be measured directly.

Also shown in Fig. 17 is the predicted relation between Mgb' and Mg_2 as a function of age and metallicity given by Worthey (1994). His models correctly predict the slope of the relation, but are offset by -0.025 mag in Mgb' (or by $+0.05 \text{ mag}$ in Mg_2), indicating a difference in the zero-point calibration of the model for one or both indices.

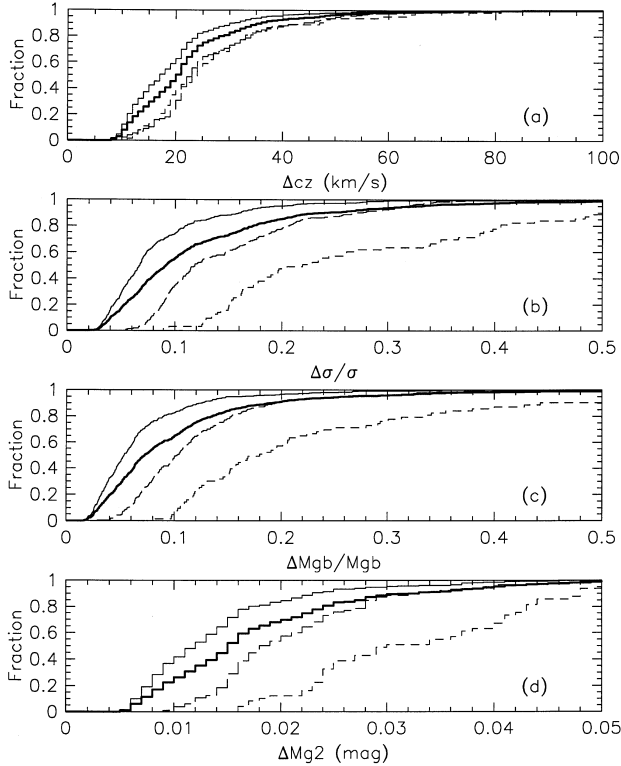


Figure 8. The cumulative distributions of the total estimated errors in the combined measurements of redshift, velocity dispersion, Mgb and Mg₂ for each galaxy. The distributions for quality classes A, B and C are shown as full, long-dashed and short-dashed lines respectively; the overall distribution is shown as the thick full line. (a) The distribution of combined errors in cz ; (b) combined relative errors in σ ; (c) combined relative errors in Mgb; (d) combined errors in Mg₂.

5 CLUSTER ASSIGNMENTS

The correct assignment of galaxies to clusters (or groups) is crucial to obtaining reliable redshifts and distances for the EFAR cluster sample. We also need to increase the precision of the cluster redshifts in order to minimize uncertainties in the peculiar velocities of the clusters. To achieve these goals we merged the EFAR redshifts with redshifts for all galaxies in ZCAT (Huchra et al. 1992, version of 1997 May 29) that lie within $3 h^{-1}$ Mpc (2 Abell radii) of each nominal EFAR cluster centre (see table 1 of Paper I). We then examined the redshift distributions of the combined sample in order to distinguish groups, clusters and field galaxies along the line of sight to a nominal EFAR ‘cluster’. We also considered the distribution of galaxies on the sky before assigning the EFAR galaxies to specific groupings.

The results of this process are shown in Fig. 18, which shows the redshift distributions of galaxies within $3 h^{-1}$ Mpc around each of the nominal EFAR clusters (labelled by their cluster ID number, CID; see Paper I) and the adopted groupings in redshift space. Note that CID = 81 (A2593-S) does not appear since it was merged with CID = 80 (A2593-N) – see below. Each EFAR galaxy was assigned to one of these groupings and given a cluster assignment number

Table 6. The distribution of errors per galaxy.

Q	Δcz (km s^{-1})		$\Delta\sigma/\sigma$		$\Delta Mgb/Mgb$		ΔMg_2 (mag)	
	50%	90%	50%	90%	50%	90%	50%	90%
All	20	36	0.091	0.240	0.072	0.188	0.015	0.032
A	17	30	0.067	0.161	0.053	0.120	0.012	0.023
B	23	44	0.118	0.270	0.103	0.194	0.018	0.033
C	24	47	0.220	0.507	0.180	0.425	0.030	0.048

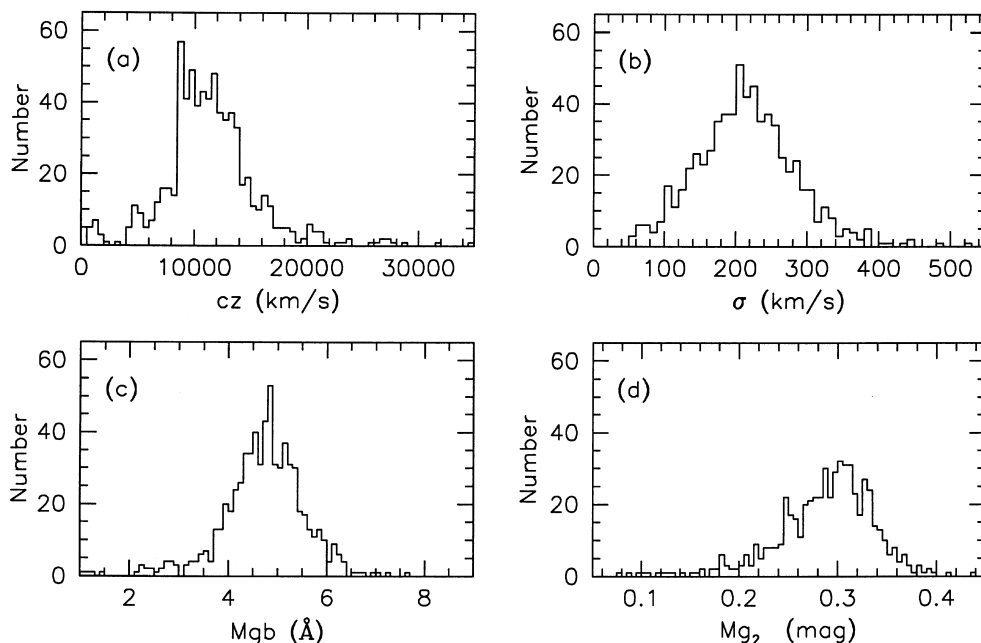


Figure 9. The distributions of (a) redshift, (b) velocity dispersion, (c) Mgb line strength and (d) Mg₂ line strength for the galaxies in the EFAR sample.

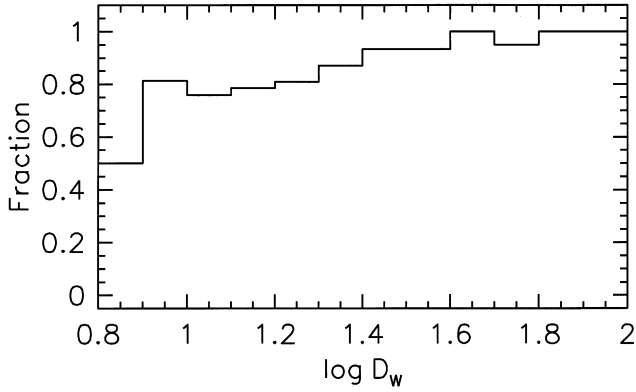


Figure 10. The fraction of programme objects for which we measured a velocity dispersion as a function of the logarithm of the selection diameter D_W (in arcsec).

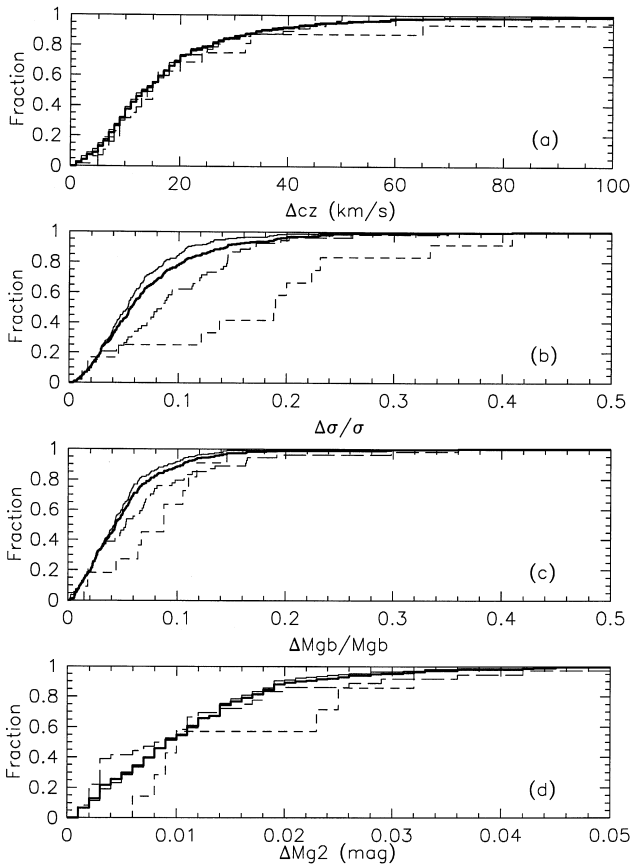


Figure 11. The cumulative distributions of the rms errors from repeat measurements of redshift, velocity dispersion, Mgb and Mg_2 . The distributions for spectral types A, B and C are shown as full, long-dashed and short-dashed lines respectively; the overall distribution is shown as the thick full line. (a) The distribution of rms errors in cz in km s^{-1} ; (b) relative rms errors in σ ; (c) relative rms errors in Mgb ; (d) rms errors in Mg_2 .

(CAN), listed in Table 5. The main grouping along the line of sight has a CAN which is simply the original two-digit CID; other groupings have CANs with a distinguishing third leading digit. The groupings (which we will hereafter call clusters regardless of their size) are labelled by their CANs in Fig. 18, which also shows the boundaries of each cluster in redshift space. The last two digits of each galaxy's CAN is its CID, apart from 41 galaxies which were

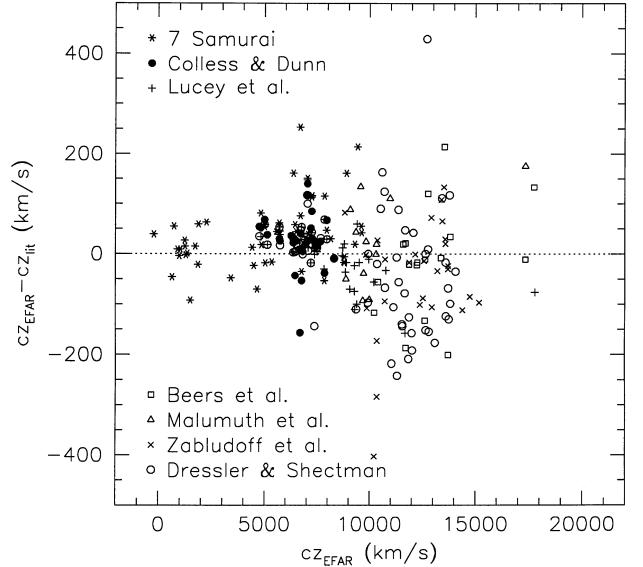


Figure 12. Differences between EFAR redshifts and those from various sources in the literature.

reassigned to other neighbouring clusters: two galaxies in CID = 33 were reassigned to CAN = 34 (GINs 254 and 255); two galaxies in CID = 34 were reassigned to CAN = 33 (GINs 263 and 264); five galaxies in CID = 35 were reassigned to CAN = 36 (GINs 270, 274, 275, 281 and 282); 14 galaxies in CID = 36 were reassigned to CAN = 35 (GINs 285–292, 295–297 and 299–301); one galaxy in CID = 47 was reassigned to CAN = 50 (GIN 406); three galaxies in CID = 59 and two in CID = 61 were reassigned to CAN = 53 (GINs 514, 517, 527, 536 and 537); five galaxies with CID = 69 were reassigned to CAN = 70 (GINs 617, 618, 619, 622 and 623); and all seven galaxies with CID = 81 were reassigned to CAN = 80 (GINs 709–715).

Table 7 lists, for each CAN, the number of EFAR galaxies, the number of EFAR+ZCAT galaxies, and the mean redshift, its standard error (taken to the error in the redshift for clusters with only one member) and the velocity dispersion. These quantities are computed both from the EFAR sample and from the EFAR+ZCAT sample. In many of the clusters the EFAR sample is greatly supplemented by the ZCAT galaxies, leading to much improved estimates of the mean cluster redshift: using EFAR galaxies only, the median uncertainty in the mean cluster redshift (for clusters with more than one member) is 177 km s^{-1} ; with EFAR+ZCAT galaxies, the median uncertainty is reduced to 133 km s^{-1} .

6 CONCLUSIONS

We have described the observations, reductions and analysis of 1319 spectra of 714 early-type galaxies studied as part of the EFAR project. We have obtained redshifts for 706 galaxies, velocity dispersions and Mgb linestrengths for 676 galaxies, and Mg_2 linestrengths for 582 galaxies. Although obtained in 33 observing runs spanning seven years and 10 different telescopes, we have applied uniform procedures to derive the spectroscopic parameters and brought all the measurements of each parameter on to a standard system, which we ensure is internally consistent through comparisons of the large numbers of repeat measurements, and externally consistent through comparisons with published data. We have performed detailed simulations to estimate measurement errors, and calibrated these error estimates using the repeat observations.

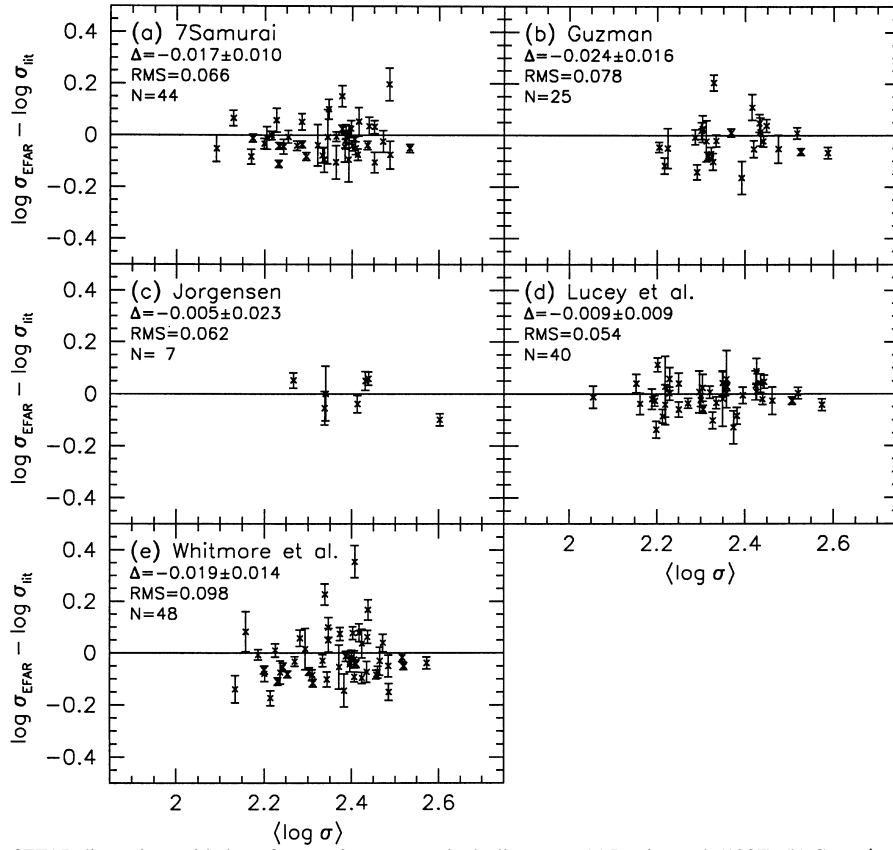


Figure 13. Comparisons of EFAR dispersions with those from various sources in the literature: (a) Davies et al. (1987), (b) Guzmán (1993), (c) Jørgensen (1997), (d) Lucey et al. (1997) and (e) Whitmore et al. (1985). In each case the mean difference, $\Delta = \langle \log \sigma_{\text{EFAR}} - \log \sigma_{\text{fit}} \rangle$, and its standard error are indicated, along with the rms scatter and the number of galaxies in the comparison.

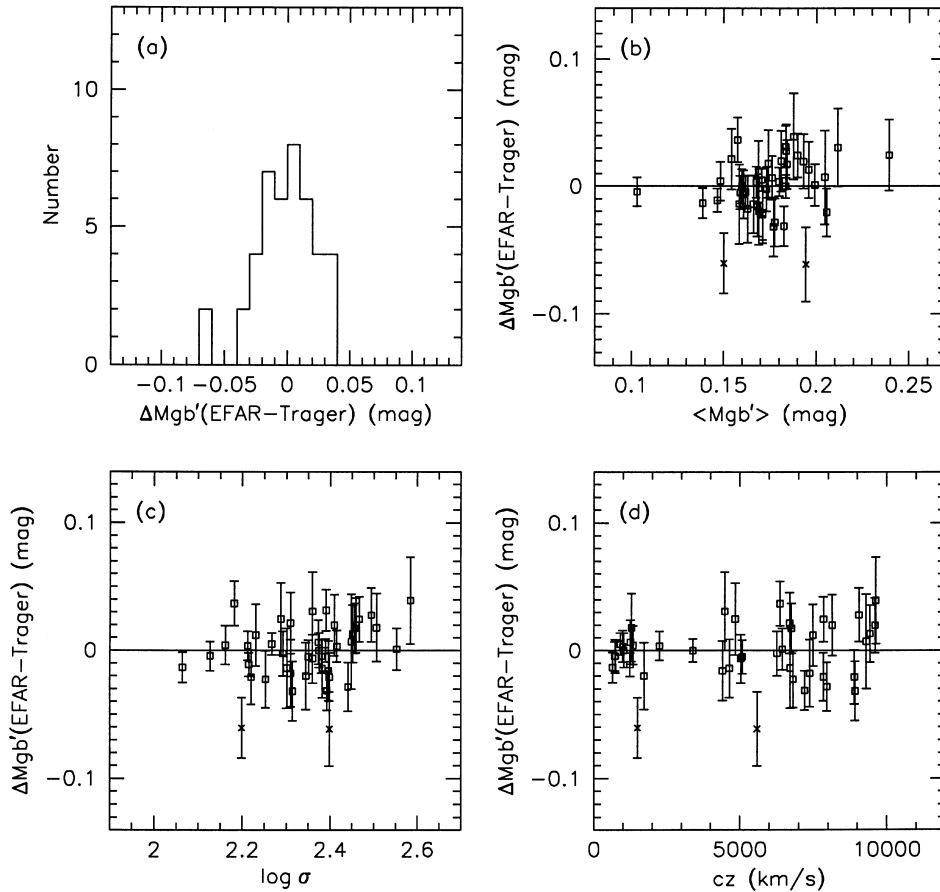


Figure 14. The residual differences in Mgb' linestrengths from EFAR and Trager et al. (1998) after applying the zero-point corrections discussed in the text: (a) the distribution of residuals; (b) the residuals as a function of Mgb' ; (c) the residuals as a function of $\log \sigma$; (d) the residuals as a function of redshift. Outliers excluded from the determination of the zero-point corrections are shown as crosses.

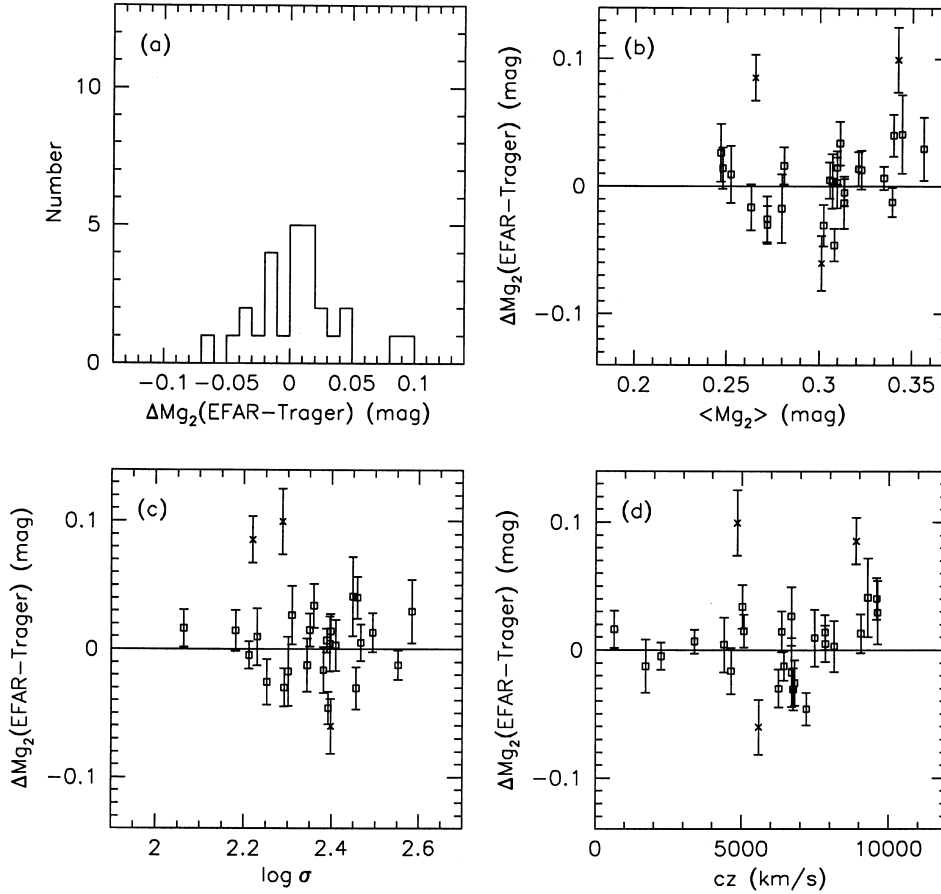


Figure 15. The residual differences in Mg_2 linestrengths from EFAR and Trager et al. (1998) *after* applying the zero-point corrections discussed in the text: (a) the distribution of residuals; (b) the residuals as a function of Mg_2 ; (c) the residuals as a function of $\log \sigma$; (d) the residuals as a function of redshift. Outliers excluded from the determination of the zero-point corrections are shown as crosses.

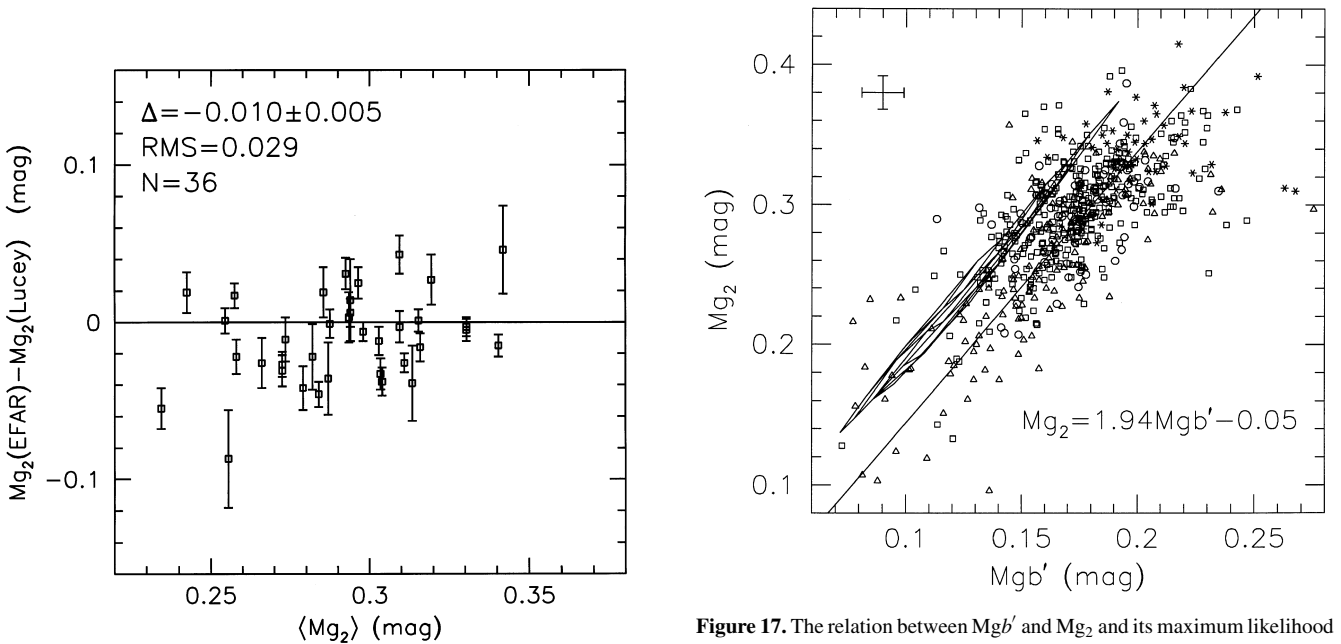


Figure 16. Comparisons of Mg_2 linestrengths obtained by EFAR and Lucey et al. (1997). The mean difference, $\Delta = Mg_2(\text{EFAR}) - Mg_2(\text{Lucey})$, and its standard error are indicated, along with the rms scatter and the number of galaxies in the comparison.

Figure 17. The relation between Mgb' and Mg_2 and its maximum likelihood fit. Ellipticals are marked by circles, E/SOs by squares, cDs by asterisks and spirals by triangles. Typical estimated errors are shown in the top left corner. The relation between Mgb' and Mg_2 as a function of age and metallicity, as predicted by Worthey (1994), is shown as the grid lying parallel to, but offset from, the data.

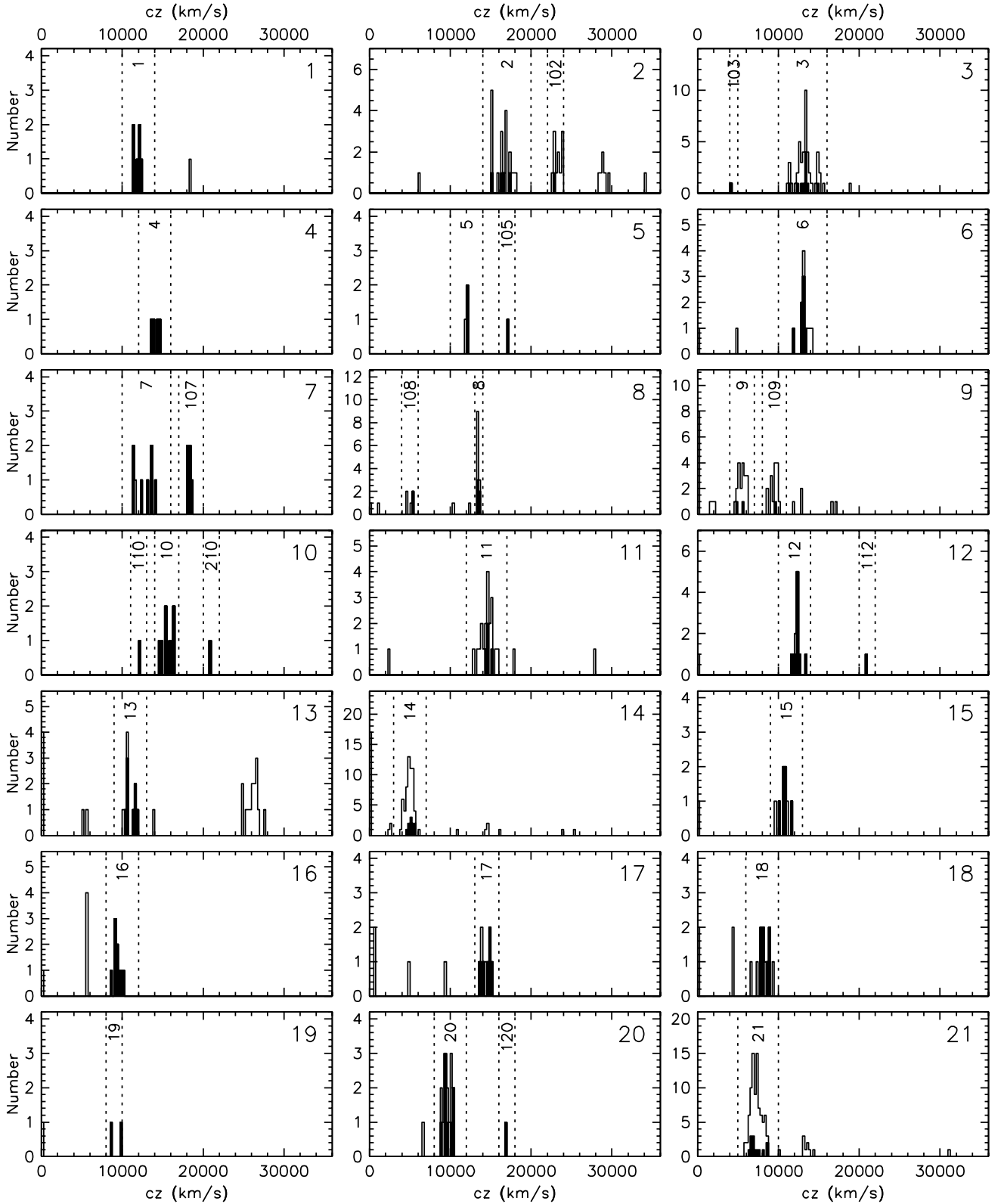


Figure 18. The redshift distributions of galaxies within $3 h^{-1}$ Mpc of each nominal EFAR cluster using the EFAR and ZCAT data. Each distribution is labelled at top right by the nominal cluster ID number (CID). The solid histogram shows the distribution of EFAR galaxies; the open histogram shows the extra ZCAT galaxies. The groupings adopted have boundaries in redshift marked by dotted lines and are labelled by their cluster assignment number (CAN). Clusters without numbers and boundaries contain no EFAR galaxies.

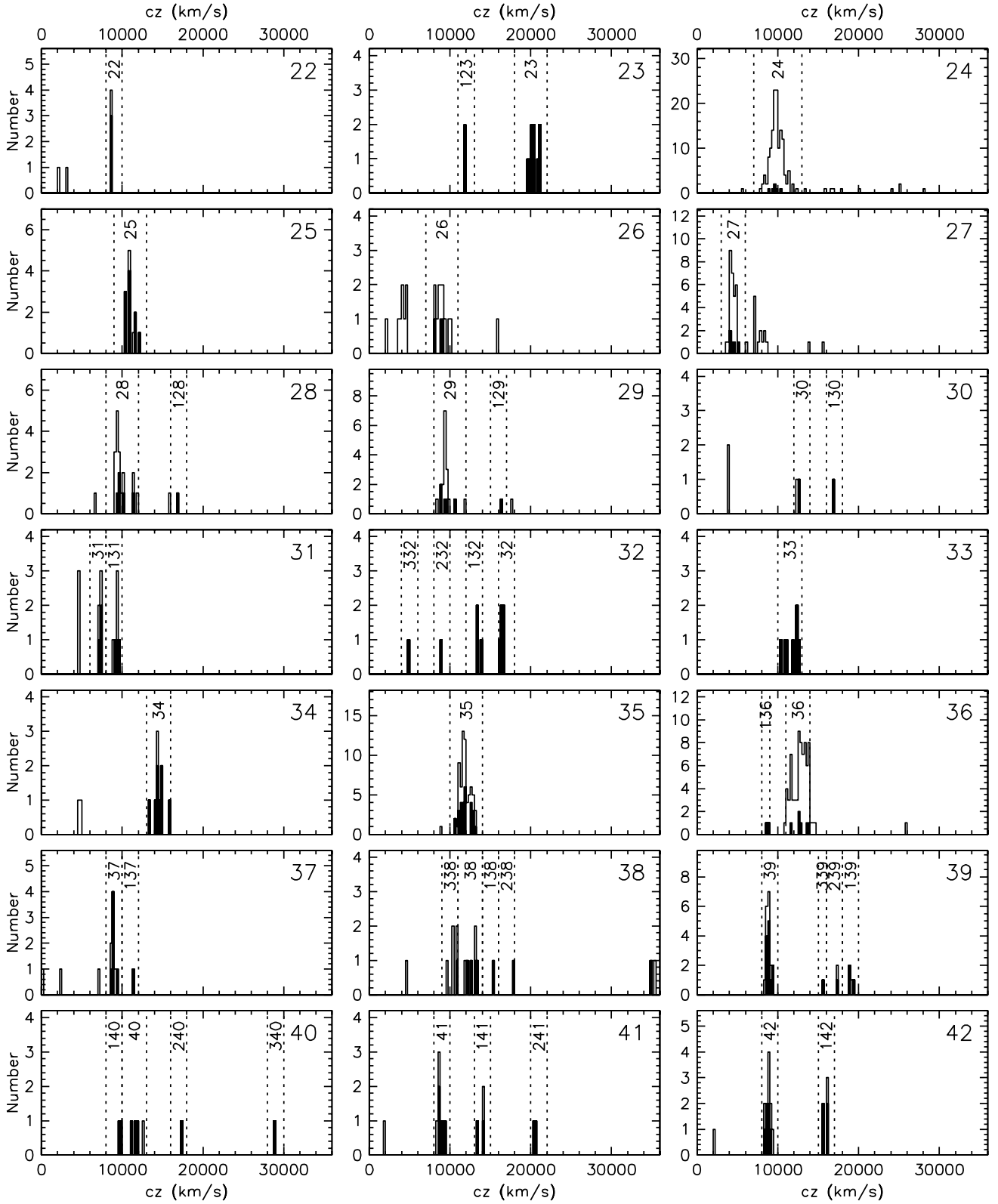


Figure 18 – continued

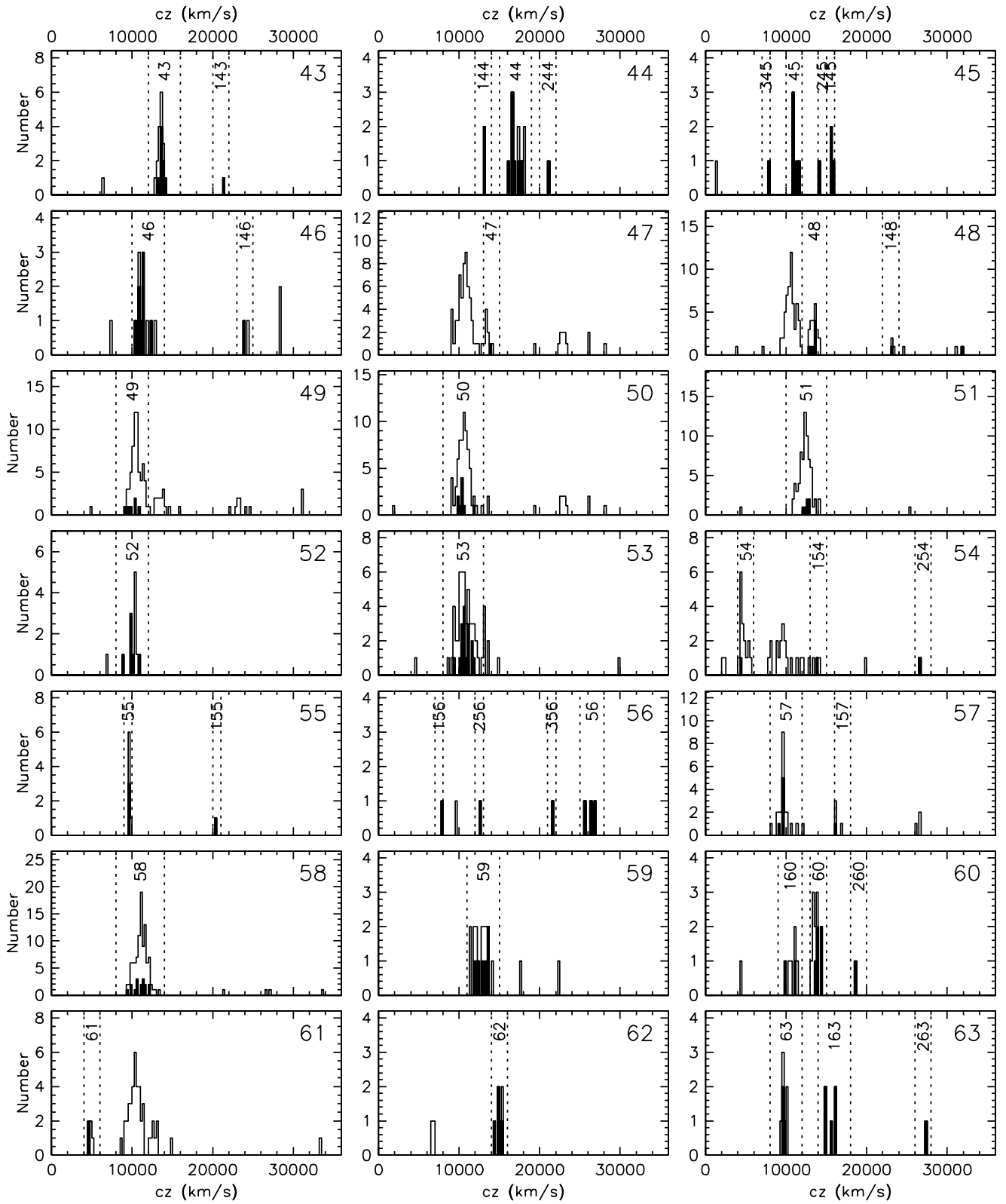


Figure 18 – *continued*

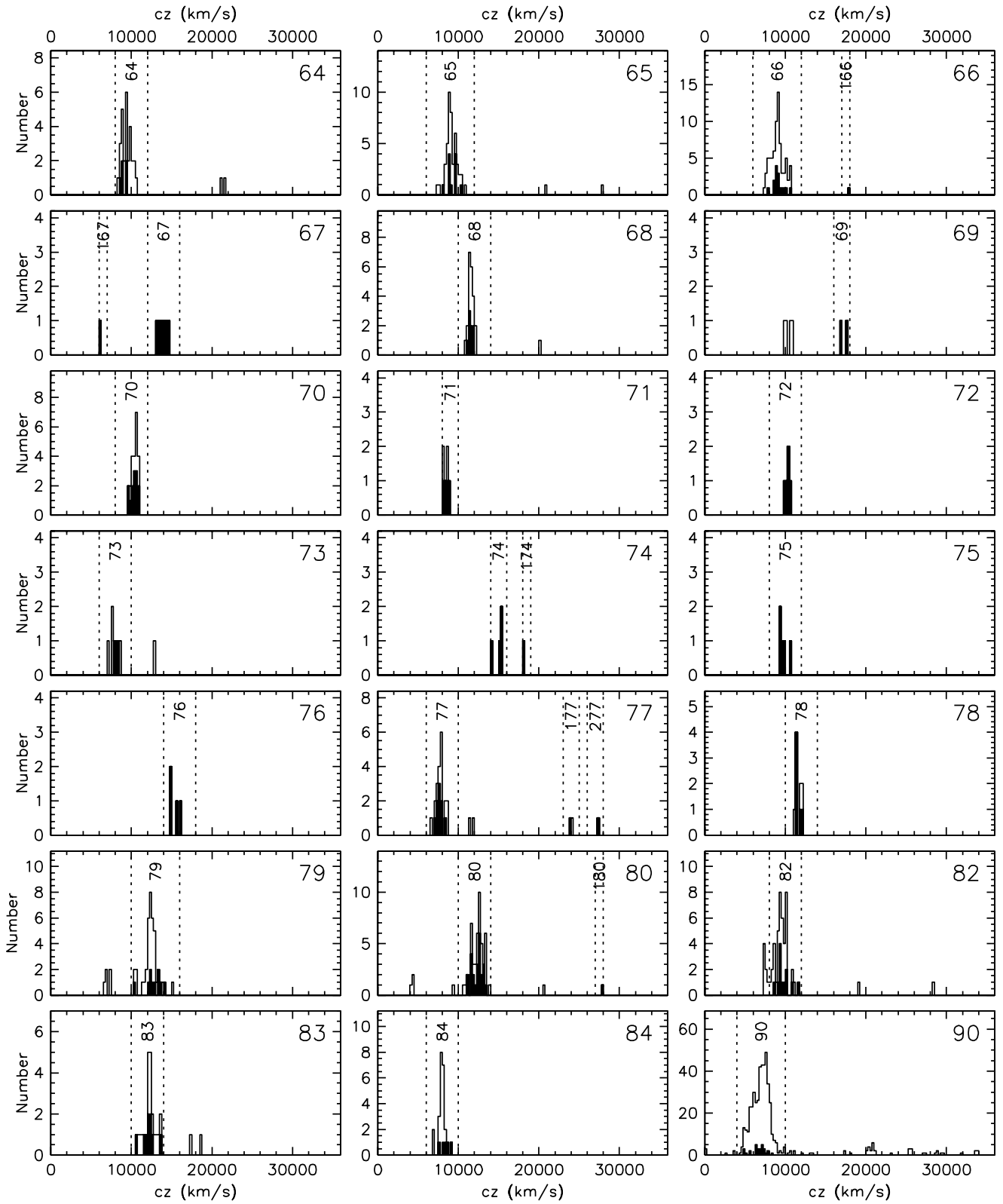


Figure 18 – *continued*

Table 7. Cluster mean redshifts and velocity dispersions.

CAN	<i>N</i>	EFAR (<i>cz</i>)	σ	<i>N</i>	EFAR+ZCAT (<i>cz</i>)	σ	CAN	<i>N</i>	EFAR (<i>cz</i>)	σ	<i>N</i>	EFAR+ZCAT (<i>cz</i>)	σ
1	6	11888 ± 191	468	6	11888 ± 191	468	241	2	20440 ± 94	134	2	20440 ± 94	134
2	4	16332 ± 466	931	19	16454 ± 227	991	42	6	8783 ± 114	280	11	8856 ± 93	309
102	1	22837 ± 71	—	11	23340 ± 134	444	142	4	15863 ± 146	292	5	15917 ± 125	280
3	10	13168 ± 287	907	45	13280 ± 158	1060	43	9	13667 ± 82	246	17	13536 ± 81	332
103	1	4127 ± 18	—	1	4127 ± 18	—	143	1	21438 ± 20	—	1	21438 ± 20	—
4	4	14090 ± 263	527	4	14090 ± 263	527	44	8	16950 ± 191	539	11	17184 ± 191	633
5	2	12082 ± 8	11	3	11994 ± 88	153	144	2	13190 ± 23	33	2	13190 ± 23	33
105	1	17154 ± 30	—	1	17154 ± 30	—	244	1	21199 ± 64	—	1	21199 ± 64	—
6	7	12899 ± 183	484	11	13193 ± 177	588	45	5	11123 ± 143	320	5	11123 ± 143	320
7	7	12782 ± 425	1123	8	12625 ± 400	1131	145	3	15757 ± 109	188	3	15757 ± 109	188
107	5	18296 ± 91	204	5	18296 ± 91	204	245	1	14024 ± 39	—	1	14024 ± 39	—
8	5	13466 ± 49	109	12	13414 ± 33	116	345	1	7789 ± 36	—	1	7789 ± 36	—
108	2	5276 ± 23	33	5	5000 ± 136	303	46	9	11136 ± 197	591	12	11321 ± 208	719
9	2	5212 ± 376	531	19	5482 ± 103	447	146	1	23945 ± 21	—	2	24166 ± 221	313
109	1	9599 ± 22	—	15	9473 ± 126	486	47	1	13876 ± 26	—	9	13576 ± 113	338
10	7	15546 ± 270	715	7	15546 ± 270	715	48	7	13466 ± 114	302	22	13455 ± 78	364
110	1	12206 ± 50	—	1	12206 ± 50	—	148	1	23099 ± 71	—	3	23233 ± 134	231
210	1	20957 ± 12	—	1	20957 ± 12	—	49	7	9944 ± 247	653	64	10528 ± 80	640
11	4	14747 ± 163	325	19	14499 ± 179	782	50	9	10427 ± 193	580	67	10548 ± 89	730
12	10	12342 ± 142	448	11	12315 ± 131	435	51	6	12593 ± 127	312	69	12353 ± 84	699
112	1	20993 ± 41	—	1	20993 ± 41	—	52	6	9982 ± 262	642	12	10215 ± 146	504
13	8	11074 ± 213	603	10	10944 ± 191	605	53	16	10786 ± 149	598	49	10675 ± 131	917
14	10	5145 ± 100	317	59	4935 ± 61	471	54	1	4410 ± 18	—	16	4699 ± 113	452
15	6	10725 ± 204	499	8	10655 ± 214	606	154	1	13830 ± 20	—	3	13818 ± 212	367
16	9	9376 ± 161	484	9	9376 ± 161	484	254	1	26612 ± 73	—	1	26612 ± 73	—
17	7	14417 ± 196	519	9	14355 ± 163	490	55	3	9635 ± 39	68	7	9636 ± 29	77
18	7	8353 ± 169	448	11	8206 ± 232	770	155	1	20285 ± 29	—	1	20285 ± 29	—
19	2	9256 ± 511	723	2	9256 ± 511	723	56	3	26280 ± 339	587	3	26280 ± 339	587
20	8	9676 ± 172	486	14	9663 ± 126	472	156	1	7955 ± 15	—	1	7955 ± 15	—
120	1	16841 ± 15	—	1	16841 ± 15	—	256	1	12730 ± 18	—	1	12730 ± 18	—
21	13	7241 ± 210	757	86	7253 ± 72	663	356	1	21558 ± 71	—	1	21558 ± 71	—
22	3	8666 ± 28	48	4	8649 ± 26	53	57	6	9520 ± 107	263	22	9582 ± 132	617
23	9	20400 ± 188	565	9	20400 ± 188	565	157	1	16075 ± 45	—	4	16237 ± 192	384
123	2	11852 ± 34	49	2	11852 ± 34	49	58	16	10943 ± 197	789	99	11106 ± 79	781
24	6	9651 ± 204	499	137	9854 ± 66	776	59	8	12864 ± 242	683	16	12693 ± 215	862
25	10	10999 ± 204	646	12	11021 ± 172	596	60	5	13993 ± 142	318	11	13707 ± 119	393
26	3	8677 ± 297	514	12	8937 ± 189	653	160	2	10474 ± 651	921	6	10730 ± 223	547
27	5	4540 ± 188	420	30	4436 ± 65	355	260	1	18545 ± 26	—	1	18545 ± 26	—
28	5	10030 ± 363	813	17	9884 ± 208	857	61	2	4742 ± 16	11	5	4878 ± 90	202
128	1	16920 ± 27	—	1	16920 ± 27	—	62	5	14918 ± 167	373	6	14991 ± 154	378
29	5	9492 ± 339	759	18	9528 ± 181	769	63	3	9699 ± 92	159	7	9777 ± 116	306
129	1	16483 ± 18	—	1	16483 ± 18	—	163	5	15563 ± 273	610	5	15563 ± 273	610
30	1	12711 ± 19	—	2	12576 ± 135	191	263	1	27486 ± 71	—	1	27486 ± 71	—
130	1	16760 ± 14	—	1	16760 ± 14	—	64	5	9062 ± 154	345	28	9423 ± 113	598
31	3	7288 ± 47	81	5	7289 ± 31	69	65	11	9223 ± 171	568	46	9137 ± 97	659
131	3	9384 ± 171	296	6	9299 ± 102	250	66	15	9156 ± 178	689	73	9014 ± 93	796
32	5	16416 ± 94	210	5	16416 ± 94	210	166	1	17795 ± 44	—	1	17795 ± 44	—
132	3	13499 ± 144	250	3	13499 ± 144	250	67	7	13876 ± 215	570	7	13876 ± 215	570
232	1	8926 ± 9	—	1	8926 ± 9	—	167	1	6086 ± 12	—	1	6086 ± 12	—
332	1	4758 ± 9	—	1	4758 ± 9	—	68	8	11514 ± 97	273	22	11547 ± 67	316
33	7	11652 ± 340	899	7	11652 ± 340	899	69	2	17344 ± 403	570	2	17344 ± 403	570
34	8	14498 ± 265	749	9	14488 ± 234	702	70	13	10320 ± 109	392	23	10396 ± 78	376
35	27	11834 ± 133	694	66	11866 ± 78	630	71	4	8465 ± 147	294	6	8442 ± 113	278
36	6	12861 ± 309	757	69	12732 ± 101	837	72	5	10311 ± 136	305	5	10311 ± 136	305
136	2	8764 ± 234	331	2	8764 ± 234	331	73	3	8121 ± 176	305	7	7935 ± 194	514
37	6	8888 ± 81	199	8	8866 ± 82	231	74	4	14965 ± 267	534	4	14965 ± 267	534
137	1	11399 ± 18	—	1	11399 ± 18	—	174	1	18187 ± 36	—	1	18187 ± 36	—
38	4	12788 ± 280	559	6	12695 ± 252	616	75	5	9784 ± 233	521	5	9784 ± 233	521
138	1	15405 ± 16	—	1	15405 ± 16	—	76	4	15365 ± 320	640	4	15365 ± 320	640
238	1	17845 ± 27	—	1	17845 ± 27	—	77	9	7639 ± 124	373	20	7728 ± 112	502
338	1	10873 ± 35	—	5	10438 ± 236	528	177	1	23926 ± 12	—	2	23992 ± 66	93
39	12	8882 ± 69	238	18	8832 ± 61	257	277	1	27282 ± 13	—	1	27282 ± 13	—
139	3	19044 ± 183	317	4	19076 ± 133	266	78	6	11504 ± 133	327	9	11571 ± 136	408
239	1	17438 ± 55	—	2	17443 ± 5	7	79	10	12692 ± 328	1036	40	12441 ± 149	944
339	1	15651 ± 30	—	1	15651 ± 30	—	80	24	12331 ± 133	649	48	12399 ± 107	740
40	3	11612 ± 217	376	4	11846 ± 280	560	180	1	27823 ± 8	—	1	27823 ± 8	—
140	2	9806 ± 160	226	2	9806 ± 160	226	82	12	9771 ± 238	824	43	9573 ± 113	741
240	1	17291 ± 26	—	1	17291 ± 26	—	83	9	12157 ± 253	759	23	12252 ± 156	748
340	1	28973 ± 28	—	1	28973 ± 28	—	84	5	8345 ± 224	500	24	7962 ± 90	442
41	5	8973 ± 135	303	7	8848 ± 127	335	90	29	6663 ± 172	924	435	6942 ± 50	1034
141	2	13685 ± 405	573	3	13794 ± 258	447							

The fully corrected measurements of each parameter from the individual spectra are given in Table 3; the final parameters for 706 galaxies, computed as the appropriately weighted means of the individual measurements, are listed in Table 5. The median estimated errors in the combined measurements (including measurement errors and run correction uncertainties) are $\Delta cz = 20 \text{ km s}^{-1}$, $\Delta\sigma/\sigma = 9.1$ per cent (i.e. $\Delta \log \sigma = 0.040$ dex), $\Delta M_{\text{gb}}/M_{\text{gb}} = 7.2$ per cent (i.e. $\Delta M_{\text{gb}}' = 0.013$ mag) and $\Delta M_{\text{g}_2} = 0.015$ mag. Comparisons with redshifts and dispersions from the literature show no systematic errors. The linestrengths required only small zero-point corrections to bring them on to the Lick system.

We have assigned galaxies to physical clusters (as opposed to apparent projected clusters) by examining the line-of-sight velocity distributions based on EFAR and ZCAT redshifts, together with the projected distributions on the sky. We derive mean redshifts for these physical clusters, which will be used in estimating distances and peculiar velocities, and also velocity dispersions, which will be used to test for trends in the galaxy population with cluster mass or local environment.

The results presented here comprise the largest single set of velocity dispersions and linestrengths for early-type galaxies published to date. These data will be used in combination with the sample selection criteria of Wegner et al. (1996, Paper I) and the photometric data of Saglia et al. (1997a, Paper III) to analyse the properties and peculiar motions of early-type galaxies in the two distant regions studied by the EFAR project.

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APPENDIX A: OBSERVING DETAILS

This appendix gives further details of the instrumental configurations used on different telescopes.

MDM 2.4-m: the Mark IIIa spectrograph, with a 1.87 arcsec wide slit, was used for all runs up to the end of 1988; from 1989 this was replaced by the Mark IIIb spectrograph, which is identical except that a 1.68-arcsec slit was used (except for run 113, when the slit width was 2.36 arcsec). For runs 101–103 a 600 line mm^{-1} grism blazed at 4260 Å was used; for all subsequent runs, a 600 line mm^{-1} grism blazed at 5700 Å was employed. The slit was usually oriented north–south. Two-pixel binning perpendicular to the dispersion direction was employed to lower the readout noise.

KPNO 4-m: the RC spectrograph and grating KPC-17B (527 line mm^{-1}) were used with the UV Fast Camera and the TI2 CCD.

KPNO 2.1-m: the Gold spectrograph/camera and grating #240 (500 line mm^{-1}) were used with the TI5 CCD.

WHT 4.2-m: the blue arm of the ISIS spectrograph was used with the CCD-IPCS imaging photon counting system. Most objects were observed using the R600B grating (600 line mm^{-1}), but one object (J26 A, GIN = 648) was observed with the R300B grating (300 line mm^{-1}).

INT 2.5-m: the Intermediate Dispersion Spectrograph (IDS) and R632V grating (632 line mm^{-1}) were used with the 235-mm camera for all runs.

SSO 2.3-m: both runs used the blue arm of the Double Beam Spectrograph (DBS) with a 600 line mm^{-1} grating. Run 130 used the Photon Counting Array (PCA), while run 132 used a Loral CCD.

MMT Blue: the 'Big Blue' spectrograph was employed with a 300 line mm^{-1} grating (blazed at 4800 Å in first order) and the Reticon detector. The MMT image stacker gave two 2.5-arcsec circular apertures separated by 36 arcsec.

MMT Red: the MMT Red Channel was used with a 600 line mm^{-1} grating (blazed at 4800 Å) and the 800 × 800 TI CCD

binned by two pixels perpendicular to the dispersion. The slit was $1.5 \times 180 \text{ arcsec}^2$, but heavily vignettted in the outer 30 arcsec in one direction.

Calar Alto 2.2-m: the Cassegrain Boller & Chivens slit spectrograph with grating #7 (60 Å mm^{-1}) was used in combination with the TEK#6 CCD.

ESO 3.6-m: the MEFOS fibre feed and the Boller & Chivens spectrograph were used. MEFOS has 58 2.6 arcsec diameter fibres (29 for targets and 29 for sky) positioned within a 1°-diameter field at prime focus. The detector was a Tektronix TK512CB CCD (ESO#32).

CTIO 4-m: the ARGUS 24-object fibre spectrograph was used. ARGUS has a 50-arcmin field at the $f/2.8$ prime focus. Each of the 24 arms holds two 1.9 arcsec diameter fibres which lie 36 arcsec apart on the sky; one arm is positioned on the target and the other on sky. The fibres feed a thermally and mechanically isolated bench spectrograph with a 510 mm focal length Schmidt blue collimator and a 229 mm focal length Schmidt camera. A Reticon II 1200 × 400 CCD detector was used with grating KPGL #3.

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