

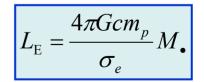
### **Accretion and Accretion Disks**

- Spherical accretion (Bondi Hoyle)
- Standard accretion (Shakura&Sunyaev)
- Advection dominated accretion (ADAF)
- ADIOS
- Corona
- Distortions (ionisation, relativistic effects)

### Spherical Accretion (Bondi Hoyle)

The radiation pressure is 
$$P_{\text{rad}} = \frac{F}{c} = \frac{L}{4\pi r^2 c}$$
, so that  $\vec{F}_{\text{rad}} = \sigma_e \frac{L}{4\pi r^2 c} \hat{r}$ ,  
where  $\sigma_e$  is the Thompson cross-section.  
This has to balance the gravity exerted over an electron-proton pair:  
 $\vec{F}_{\text{grav}} = -\frac{GM_{\bullet}(m_p + m_e)}{r^2} \hat{r}$   
The condition  $\vec{F}_{\text{rad}} \leq \vec{F}_{\text{grav}}$  then implies that  
 $L \leq \frac{4\pi G c m_p}{\sigma_e} M_{\bullet} \approx 6.31 \times 10^4 M_{\bullet} \text{ erg s}^{-1} \approx 1.26 \times 10^{38} (M_{\bullet}/M_{\odot}) \text{ erg s}^{-1}$   
This is known as the Eddington limit.

The Eddington luminosity is the maximum luminosity emitted by a body of mass  $M_{\bullet}$  that is powered by spherical accretion.



### Standard Accretion (Shakura&Sunyaev 1973)

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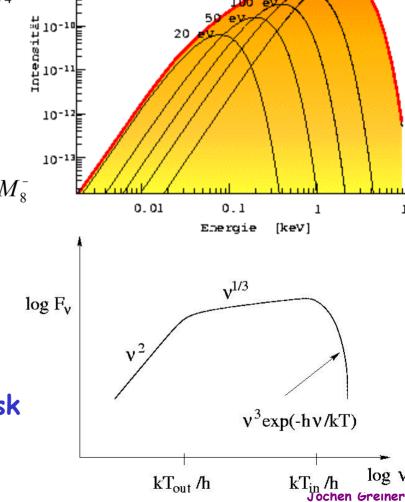
 virial theorem: half of potential is radiated away, other half heats gas

$$L = \frac{1}{2}\dot{U} = \frac{GM_{\bullet}\dot{M}_{\bullet}}{2r} = 2\pi r^2 \sigma T^4 \implies T = \left(\frac{GM_{\bullet}\dot{M}_{\bullet}}{4\pi\sigma r^3}\right)^{1/4}$$

Temperature distribution

$$T(r) \approx \left[\frac{3GM_{\bullet}\dot{M}_{\bullet}}{8\pi\sigma R_{\rm S}^3}\right]^{1/4} \left(\frac{r}{R_{\rm S}}\right)^{-3/4} = 6.3 \times 10^5 \,{\rm K} \,\left(\frac{\dot{M}_{\bullet}}{\dot{M}_{\rm E}}\right)^{1/4} M_{\rm S}$$

- 2 regions:
   radiation pressure > gas pressure
   inner region; torus-like
   gas pressure > radiation pressure
   optically thin, geometrically thick disk
- Innermost radius =  $6 R_g = 6 GM/c^2$



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### **Thin Accretion Disk Models**

- Cooling-Dominated Flows: viscous heating of the gas is balanced by local radiative cooling.
- Advection-Dominated Flows (ADAF): electrons decouple from protons: radiative cooling is very inefficient and most of the dissipated energy is advected into the black hole.
- Advection-Dominated Flows with Inflow and Outflow (ADIOS): outflow is possible based on the Bernoulli number. is interesting as it may describe jets

### **Main Properties of Accretion Disk**

### Inflow of Matter

coming through Lagrange point from companion
 being gas and/or stars from galaxy

### Outwards Transport of Angular Momentum

Through viscous stress

through hydromagnetic winds

Through amplification of local viscous stress through the onset of turbulence resulting from possible hydromagnetic, convective, or shear flow instabilities

through propagating waves

Through torque resulting from the presence of nonaxisymmetric unstable modes in self-gravitating and geometrically thick disks

#### Electron Corona

<sup>©</sup>either spherical around central object

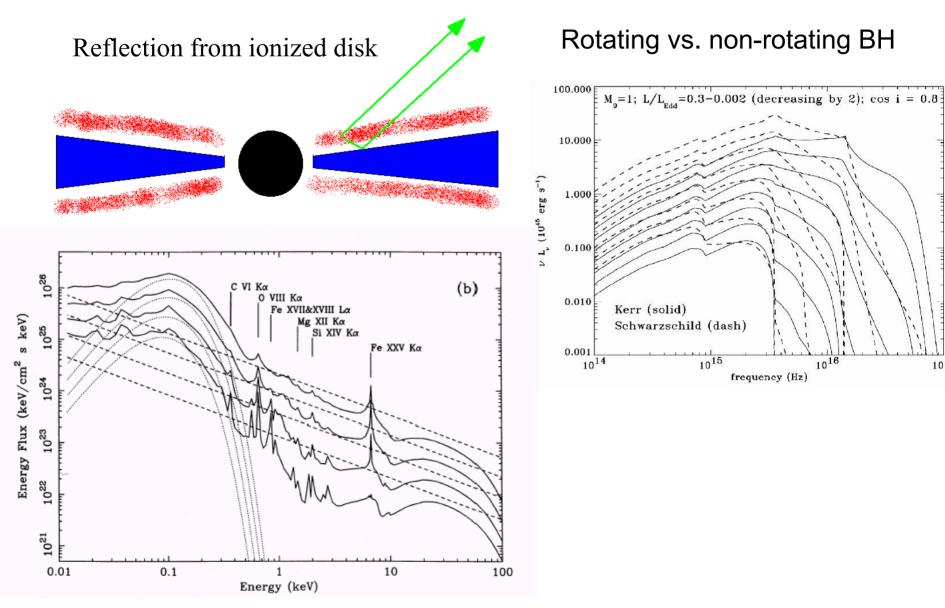
@ or in two slabs (sandwich-like) above/below disk

#### Matter Outflow

<sup>(3)</sup>either spherical wind

@ or directed outflow (jet)

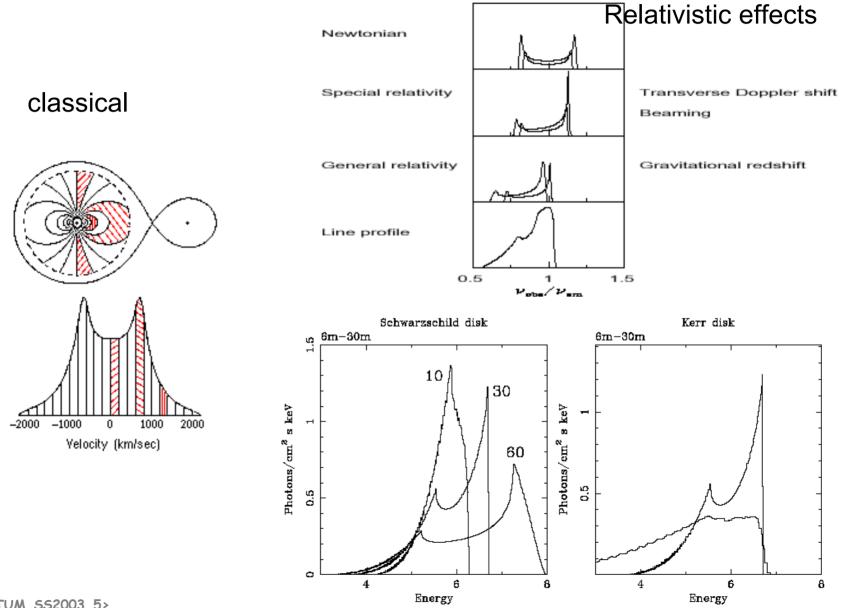
### Accretion Disk: Other Complications (I)



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### Accretion Disk: Other Complications (II)



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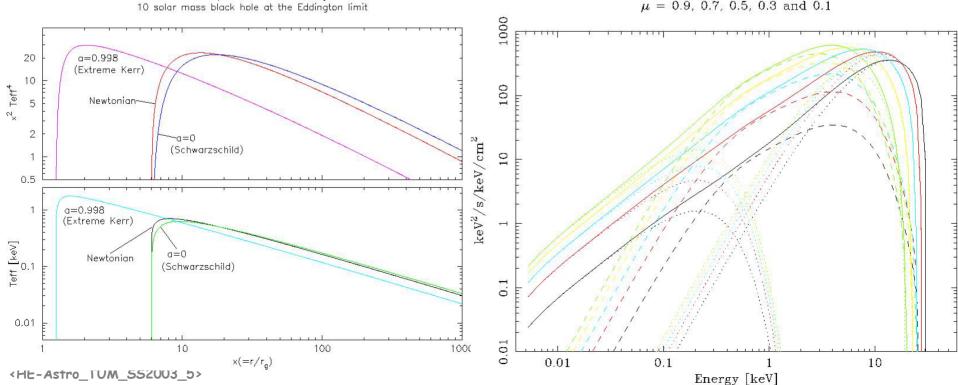
### Accretion Disk: Observable Properties

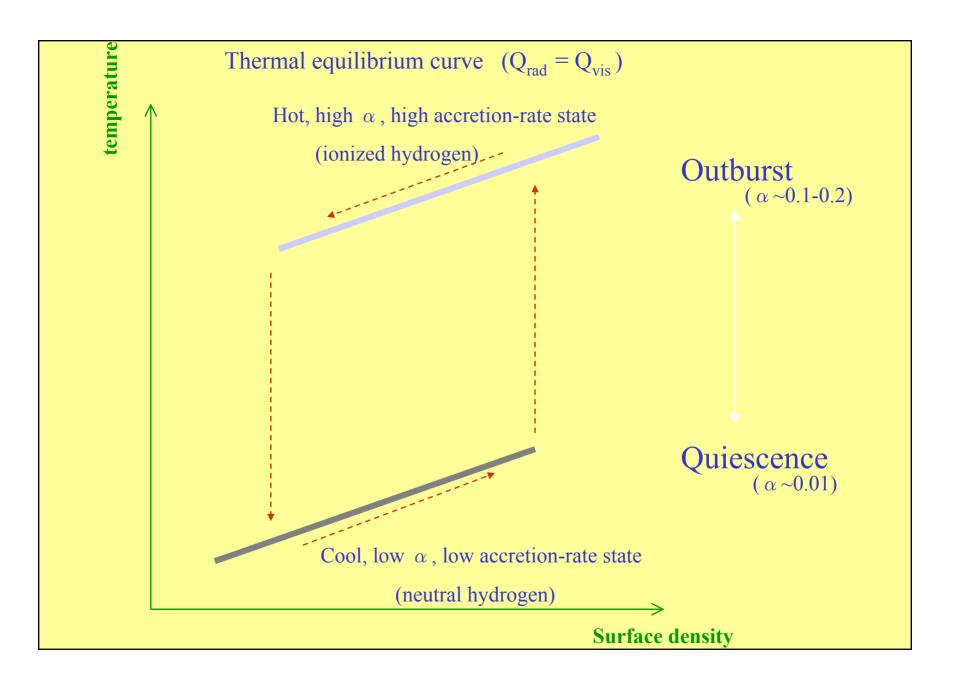
#### • mass, mass accretion rate (luminosity) $\rightarrow$ spectral shape determined

Maximum disk color temperature  $T^{(max)}_{col} \sim 1.2 \text{ keV}$   $((T_{col}/T_{eff})/1.7) (M/M_{Edd})^{1/4} (M/7M_{\odot})^{-1/4}$ is directly measured from observed Xray spectral shape Kerr disk can be much hotter, as innermost radius  $\rightarrow 1.24$  Rg Disk temperature and emissivity Hard emission from very inner parts (1.26  $r_g < r < 7 r_g$ ) is enhanced for inclined Kerr disks (due to Doppler boost)

When the disk is face-on, emission from inner part is weak, and the spectrum is not very different from the Schwarzschild

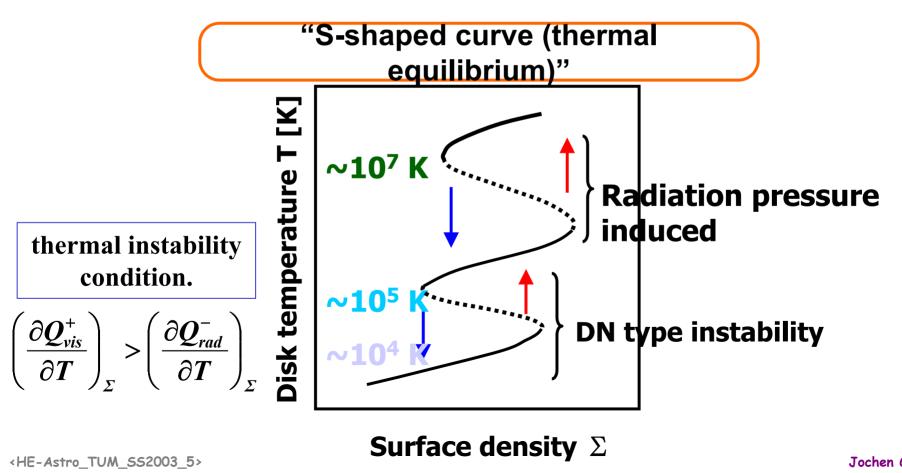
**CASE** Kerr disk at d = 1 kpc, M= 1 M<sub> $\odot$ </sub>. Eddington limit  $\mu$  = 0.9, 0.7, 0.5, 0.3 and 0.1





### **Time-dependent Properties of Accretion Disk**

- Disk instability ( $Q_{vis} > Q_{rad} > \sim Q_{adv}$ ,  $P \sim P_{rad} \sim T^4$ )
- Dwarf Nova type instability
- Different ignition mechanism for radiation pressure-induced instability



Surface density  $\Sigma$ 

### Compact Sources (WD, NS, BH)

#### As Single Stars

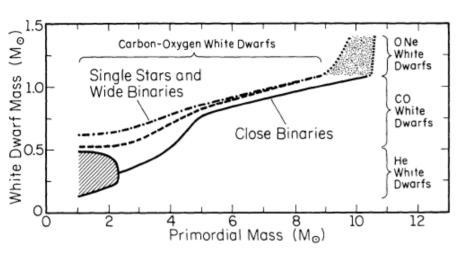
Equation of State
Density/Size
Rotation
Magnetic Field

# In Binary Systems Types according to Central Object

Types according to Companion

### Singles: Only neutron stars observable at X/Gamma-rays. Binaries: All 3 types observable at X/Gamma-rays.

### Formation of Compact Objects



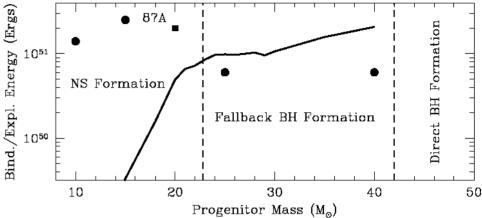


Figure 3. Binding energy/explosion energy vs. progenitor mass assuming no mass loss. The solid line is the binding energy of all but the inner  $3 M_{\odot}$  of the stellar core. If the explosion energy is not at least this powerful, the star will collapse to a black hole. The four circular dots denote the explosion energies from core-collapse simulations and the square is the predicted explosion energy observed from SN 1987A.

#### Type Depends on Mass of Companion star

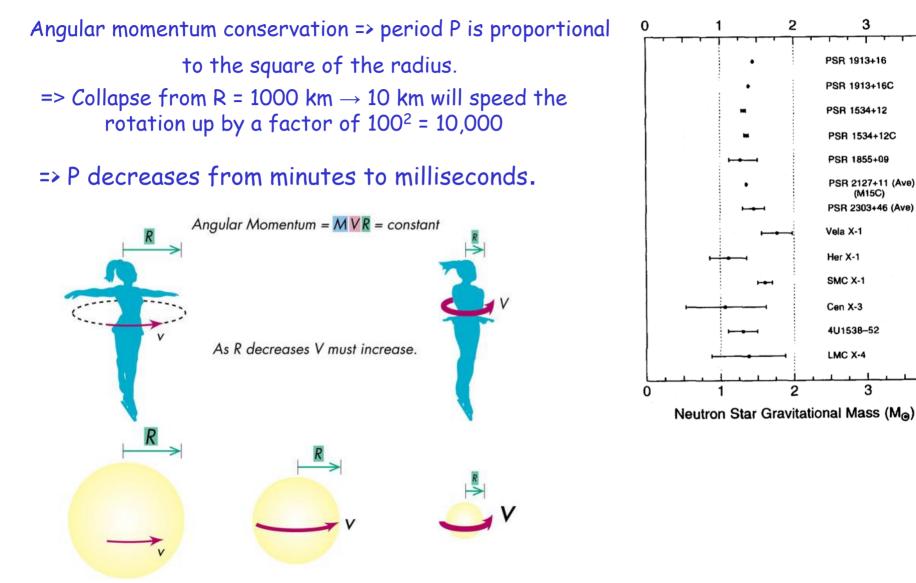
@ 1-8 M $_{\odot}$   $\blacksquare$  White Dwarf <1.4 M $_{\odot}$ @ 8-22  $M_{\odot}$   $\implies$  Neutron Star 1.4-3  $M_{\odot}$ Signature Strate St

- >3 M<sub>o</sub>
- WD Mass Spectrum

Peak at ~0.5-0.6 M<sub>o</sub>

Ρ	roperties of	<b>Compact Objects</b>	
V	Vhite dwarf	Neutron Star	Black Hole
Mass	0.4-1.4	1.4 (-3)	>3
Pressure	>electron pressure	>neutron pressure p + e <sup>-</sup> $\rightarrow$ n + v <sub>e</sub>	
Density	10 <sup>6</sup> g/cm <sup>3</sup>	$10^{14} \text{ g/cm}^3$	arbitrary
Radius	1000 km	10 km	horizon
Surface T	<10 <sup>5</sup> K	10 <sup>6</sup> к	disk size
Moment of Inertia	$10^{51} \text{ g cm}^2$	10 <sup>45</sup> g cm <sup>2</sup>	
Spin Period	10-1000 sec	0.001 sec	mass
Magnetic Field	10 <sup>8</sup> - 10 <sup>9</sup> G	$10^{10} - 10^{14} G$	none
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### What is a Neutron Star?



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PSR 1913+16 PSR 1913+16C

PSR 1534+12

PSR 1534+12C PSR 1855+09

Vela X-1 Her X-1 SMC X-1

Cen X-3 4U1538-52

LMC X-4

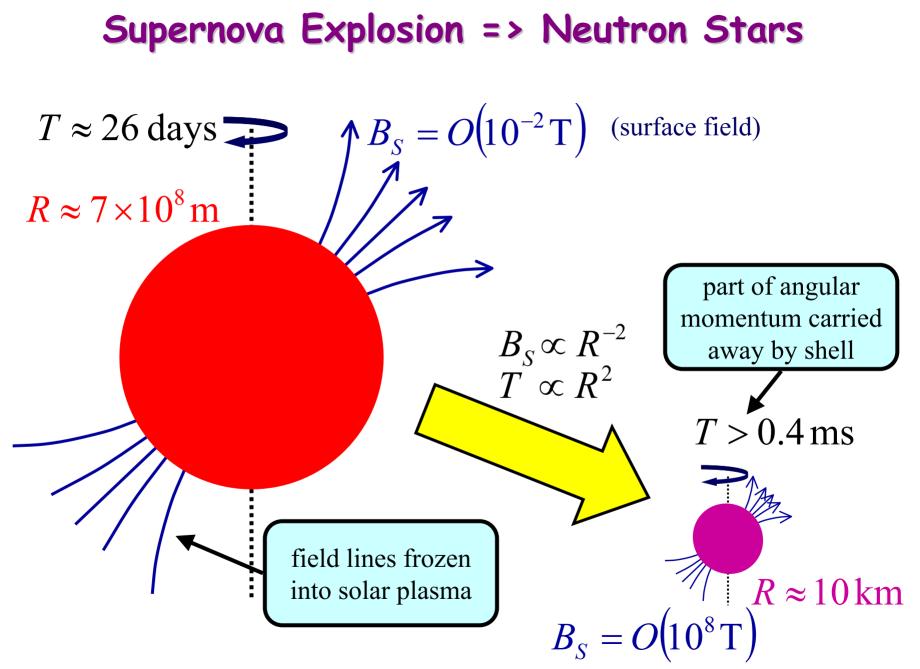
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PSR 2127+11 (Ave) (M15C) PSR 2303+46 (Ave)

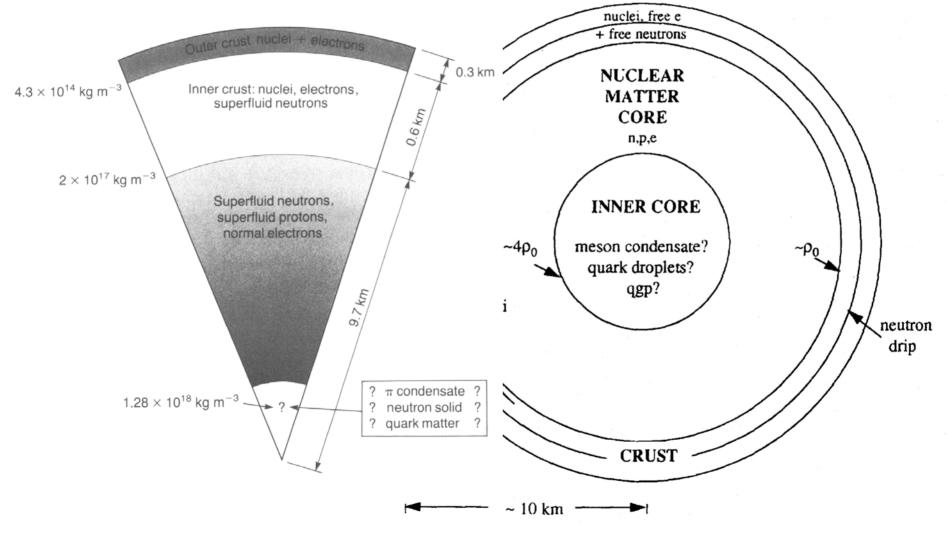
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### Structure

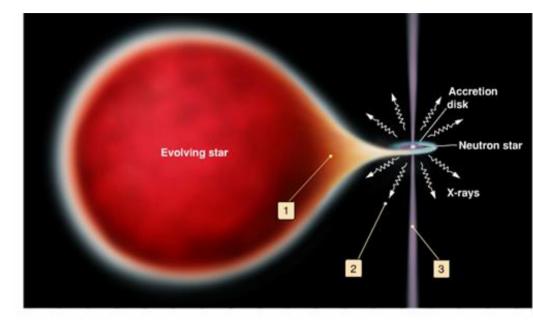


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# X-ray Binary

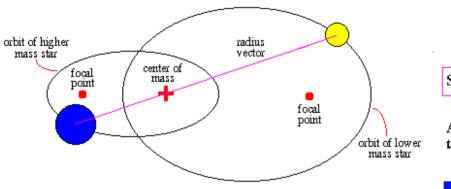
- When a compact object is part of a binary system
- Compact object can be neutron star, white dwarf or black hole



- When the other star fills its Roche limit
   starts feeding matter to neutron star
- The compact object has an accretion disk
   A heated by matter falling onto it
- The accretion disk heats enough to glow in the x-ray part of the spectrum

## X-ray Binary

#### Binary Star Orbit



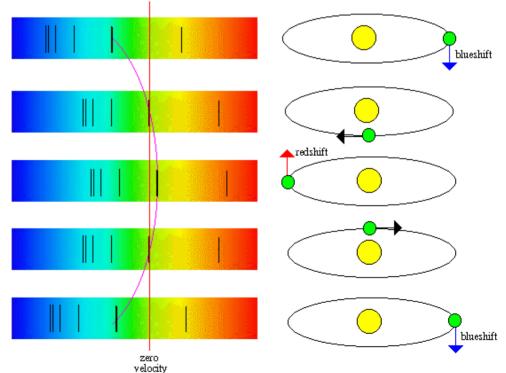
Spectral lines of companion allow velocity and period determination

Mass function

$$f(M) = \frac{K^{3}P}{2\pi G} = \frac{M_{2}^{3} \sin^{3} i}{(M_{1} + M_{2})^{2}}$$

Spectroscopic Binary

A spectroscopic binary is where there is evidence of orbital motion in the spectral features due to the Doppler effect



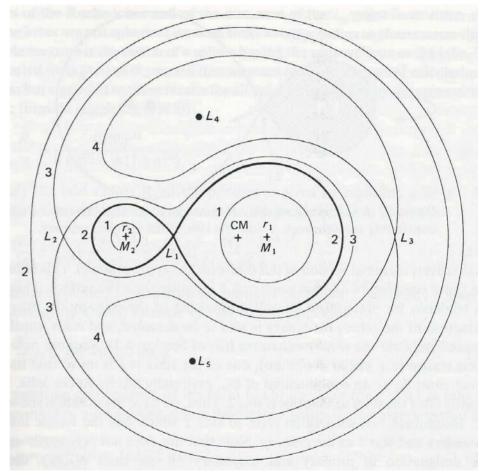
### Properties of the Roche Lobe

### Roche lobe Potential:

$$\Phi_{R}(r) = \frac{-GM_{1}}{|r-r_{1}|} - \frac{GM_{2}}{|r-r_{2}|} - \frac{1}{2}(\omega \times r)^{2}$$

where  $r_1$  and  $r_2$  are the position vectors of the centers of the two stars and  $\omega$  is the angular velocity of the binary.

Equipotential surfaces: shape of the plot is completely dependent on the mass ratio q ( =  $M_2/M_1$ ), while the scale is dependent on a.



### Types of X-ray Binaries

White Dwarf:

#### CV (cataclysmic variable):

Dwarf novae, polars, VY Scl stars, supersoft X-ray sources, novae

#### Neutron Star or Black Hole: HMXB

Luminous (early OB star) high-mass companion wind accretion (strong winds for M>10 M<sub>☉</sub>) hard X-ray spectra (T>100 million K) often pulsating X-ray eclipses Galactic plane Population I

#### LMXB

Optically faint (blue) counterpart low-mass companion disk accretion (RLOF instable for M>1.4M<sub>☉</sub>) soft X-ray spectra (T~30-80 million K) non-pulsating no X-ray eclipses Gal. Centre + bulge older, population II

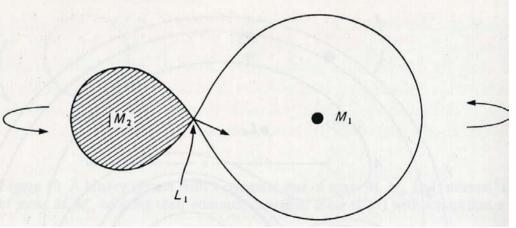
### What is a CV?

A CV is a binary system in which one star (designated the primary) accretes matter from a secondary star that, for whatever reason, has exceeded its Roche lobe. More specifically, it is a semi-detached binary with a white dwarf primary and a secondary that is typically a late-type main sequence star.

Miscellaneous Facts

- $\approx$  1 million in our galaxy
- *a* is typically a few times the Earth-Moon separation
- Typical orbital periods *P* range from 1-10 hours

Geometry of a CV:



The Roche lobe becomes filled in CV's most typically by a decreasing semimajor axis. The driving mechanism for the shrinkage is driven by magnetic breaking of the secondary if P > 3 hours or gravitational radiation if P < 2.

### Differences between NS/BH Accretion

	Neutron star	Black Hole
Magnetic field	yes	no
Solid surface	yes	no
Disk extends	down to NS surface (depending on B)	innermost stable orbit
Spin	does not matter	strong dependence
	spectra of NS accretors should be more complicated!	

Observations: spectra rather similar!!! Why?

## **Typical-Source Energy Spectra**

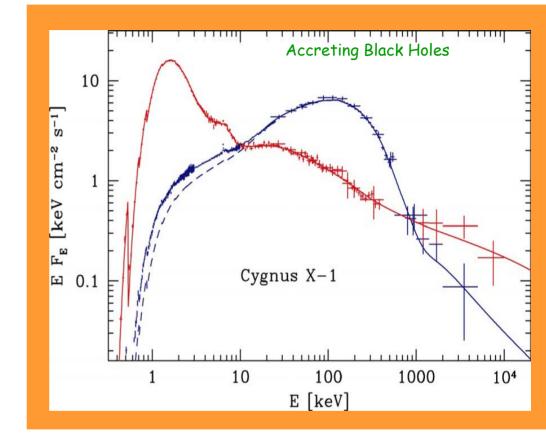
#### Thermal Components

Stellar Surfaces
Accretion Disks
Plasma Bubbles

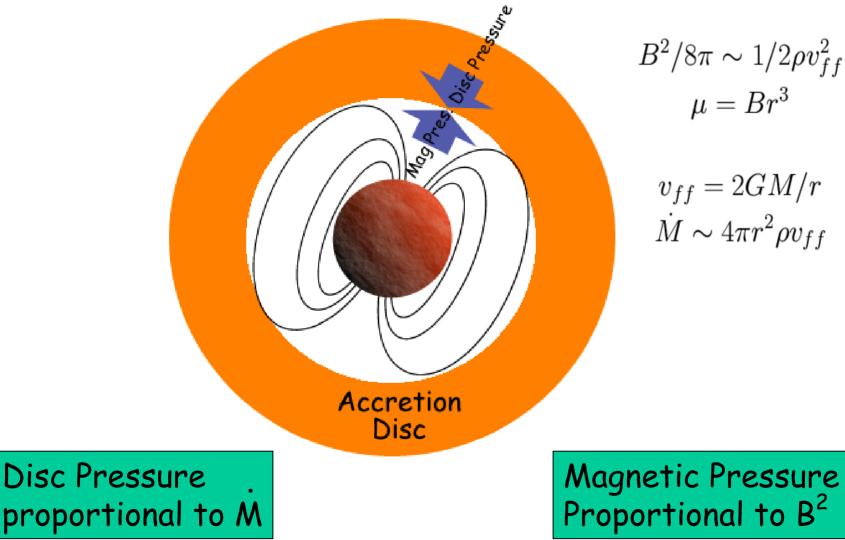
#### Non-Thermal Components

#### ☆ Accelerated Particles

- Synchrotron Radiation
- Bremsstrahlung
- Inverse Comptonization



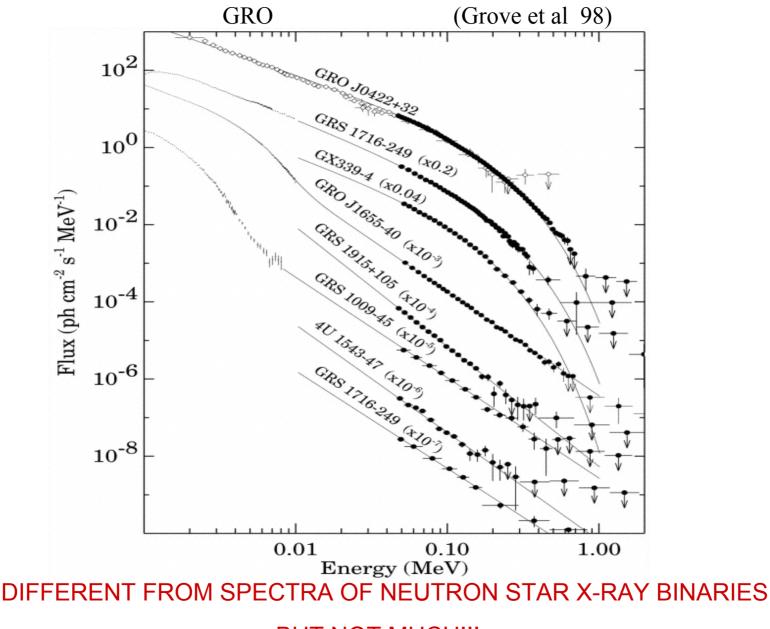
### **Disc - Magnetic Field Interaction**



 $B^2/8\pi \sim 1/2\rho v_{ff}^2$  $\mu = Br^3$ 

 $v_{ff} = 2GM/r$  $\dot{M} \sim 4\pi r^2 \rho v_{ff}$ 

### X/Gamma-Ray Spectra of Black Hole Binaries



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**BUT NOT MUCH!!!** 

Jochen Greiner

## **Dynamically Established BH Systems**

Measure mass of candidate black hole via mass function

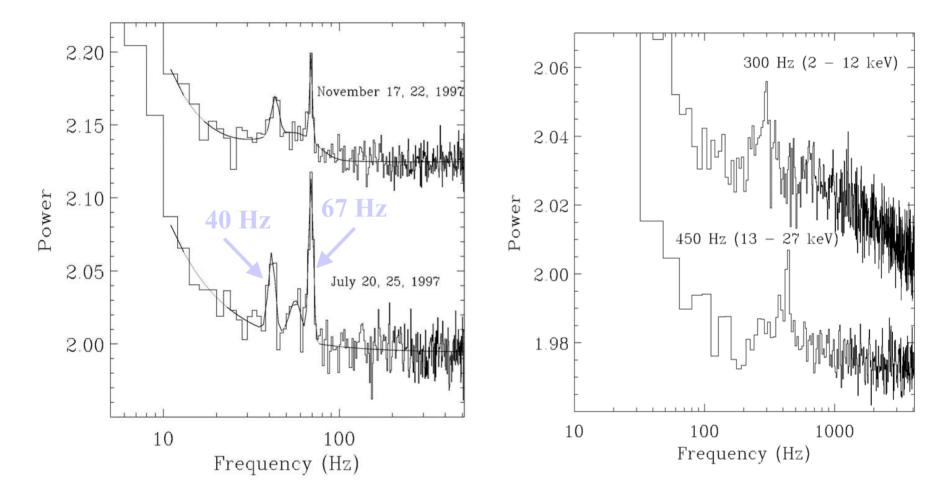
This constitutes the strongest evidence for the existence of stellar-mass black holes in binary systems. Currently there are at least 10 systems for which dynamical mass determinates result in the compact object having mass greater than 3  $M_0$ , the theoretical upper limit of a neutron star.

Source Name	Alternate Name	BH Mass (M <sub>sun</sub> )
GRS 1915+105	_	10-18
0538-641	LMC X-3	7 - 14
0540-697	LMC X-1	4 - 10
GRO J 0422+32	XN Per 1992	$3.57 \pm 0.34$
A0620-00	XN Mon 1975	4.9 - 10
GRS1124–683	XN Mus 1991	5.0 - 7.5
GRO J1655–40	XN Sco 1994	$7.02\pm0.22$
Н 1705–250	XN Oph 1977	$4.9 \pm 1.3$
1956+350	Cyg X-1	7 - 20
GS2000+25	XN Vul 1988	$8.5 \pm 1.5$
GS2023+338	V404 Cyg	$12.3 \pm 0.3$

### Accretion Disks Exhibit "Quasi-Periodic" Oscillations (QPO)

- Rich spectrum of modes: coherence up to Q ~20
- Highest frequency QPOs must come from close to horizon tied to marginally stable orbit?
  - Detailed origins of different modes still unclear
  - ☆ Period of circular orbits
  - ☆ Precessional modes
    - Lense-Thirring precession tied to BH spin dragging of inertial frames
    - Precession of periastron
  - ☆ Pulsational modes: "diskoseismology"
    - Disk acts as resonant cavity under influence of GR effects
  - ☆ Bending modes
    - Redundant probes of spacetime structure
  - ... all depend on metric in different ways

### QPOs in two Microquasars



**GRS 1915+105** 

**GRO J1655-40** 

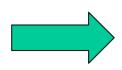
### **Transient Phenomena**

#### Different causes

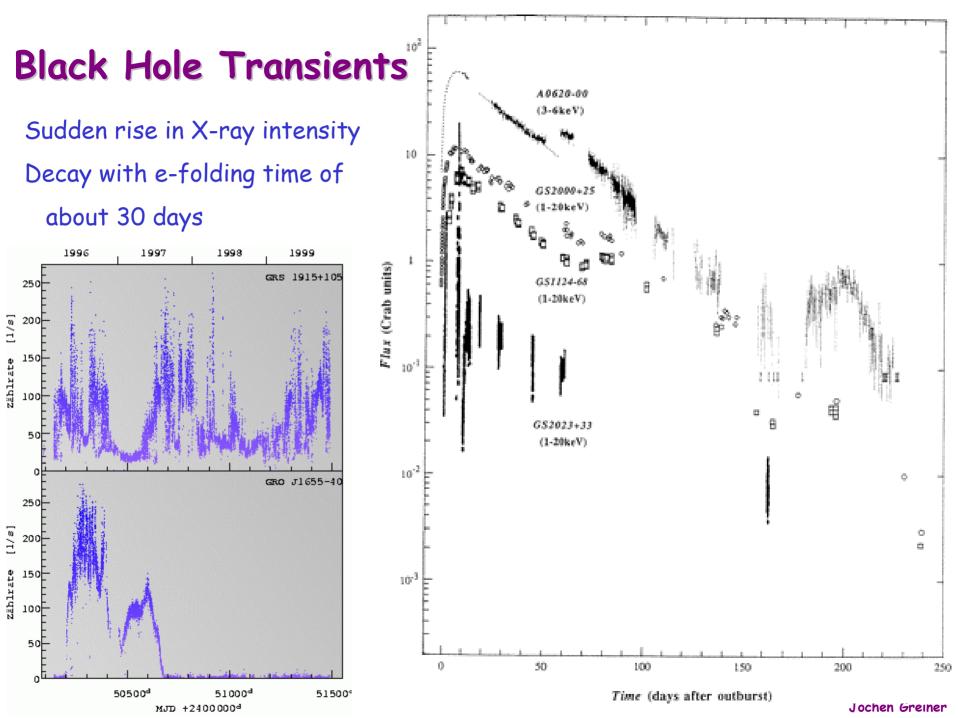
Accretion disk instabilities (Dwarf novae, BH transients)
Variable mass overflow (VY Scl stars)
Variable density (NS/Be systems)
Variable irradiation

#### Various types of transient sources

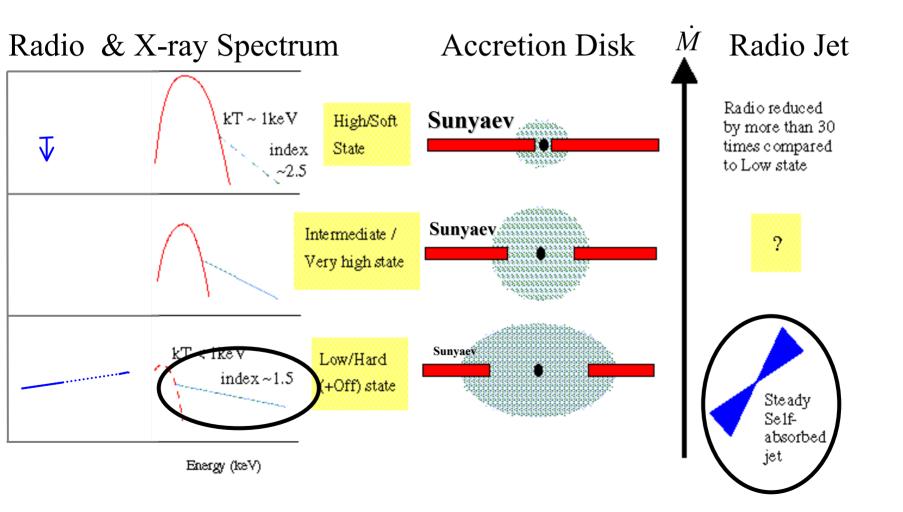
NS/Be systems
Dwarf novae
BH transients
Microquasars



Complicated: time dependent phenomena in accretion disk Short-lived: can investigate certain states only ones

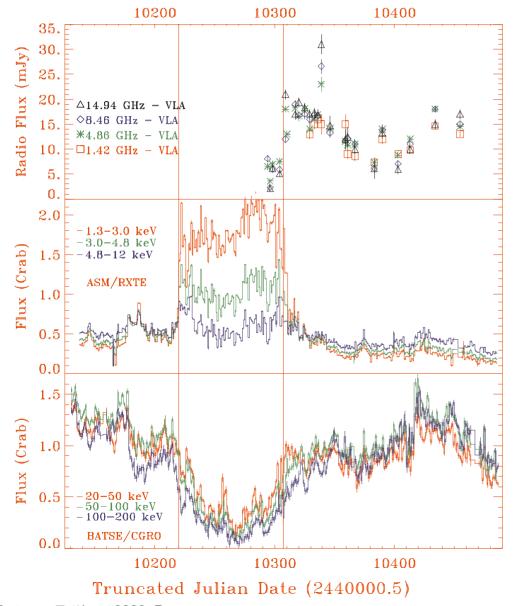


### The Power-Evolution of Black Hole Transients



#### Jochen Greiner

### High-/Low States



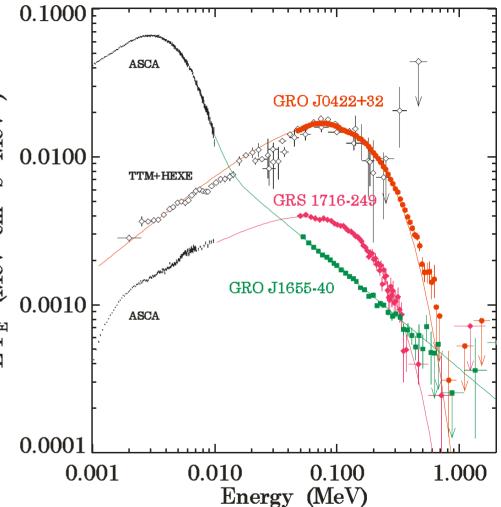
Spectral state changes in the persistent source Cyg X-1. Soft X-ray emission is anticorrelated with hard X-ray and gamma-ray emission. The transition from the X-ray high state is marked by a radio flare.

### High-Energy Emission during High-/Low States

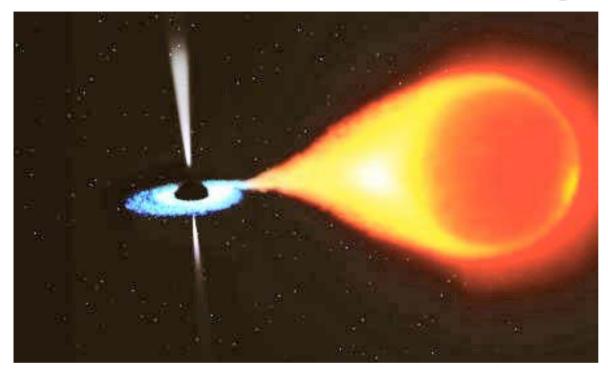
Several high-energy states:

• "X-ray high/soft" state, the soft component is dominant. It is emitted from the inner accretion disk and can be well modeled by a multi-color disk blackbody spectrum, where the temperature falls as radius r<sup>-3/4</sup>. The hard component is a simple power law and may be the signature of inverse-Compton scattering from a non-thermal relativistic electron population.

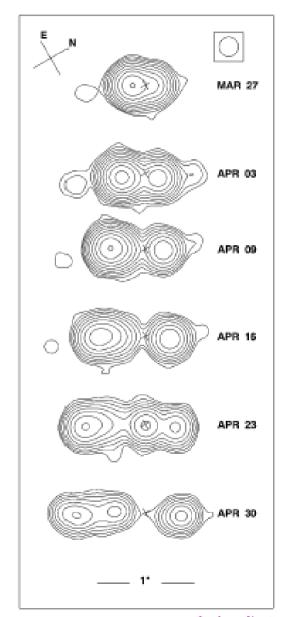
•"X-ray low/hard" state, the black-body component is weak or absent, having been Compton up-scattered to gamma-ray energies by a thermal plasma with temperature ~100 keV. The plasma likely exists either as a hot inner disk or as a patchy corona above a cold disk.



### Plasma Jets from Accreting Compact Stars

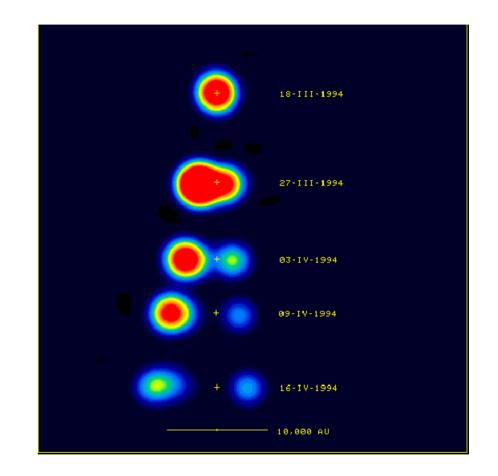


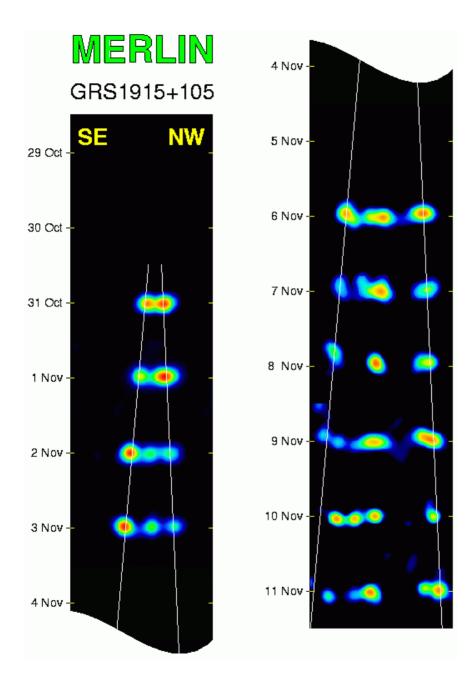
- "Micro-Quasars": Jets from Accreting Stellar-Mass Compact Objects (NS + BH)
- X-Ray and Radio/γ Emission from Hot Plasma & Relativistic Particles
- "Superluminal" Plasma Blobs Traced in Radio Emission



### Superluminal Motion

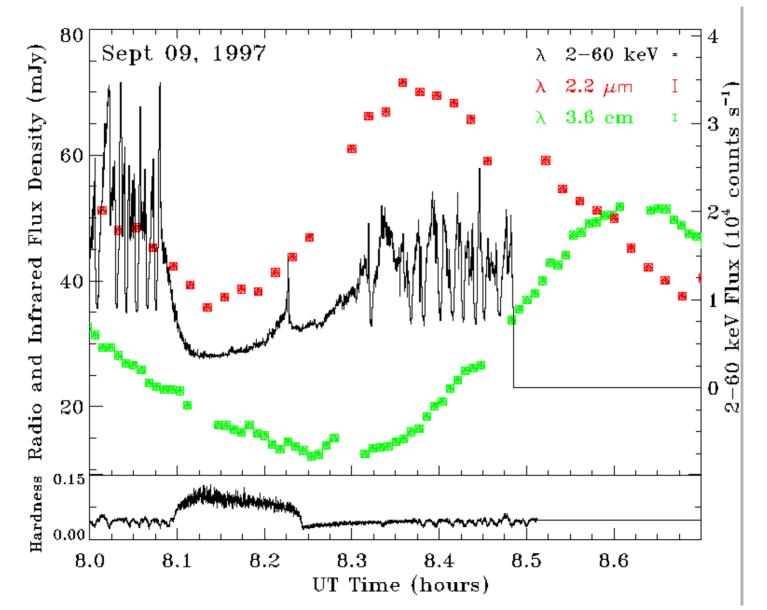
Relativistic, radio-loud jets from the cores of external Active Galaxies have been known for some years. Recently, such apparent superluminal motion has been detected from several objects within our Galaxy for the first time. A sequence of 3.6-cm radio images from the black hole candidate GRS 1915+105 shows radio-luminous ejecta moving in the sky. The images were taken with the Very Large Array with an angular resolution of 0.2 arc seconds. The illusion of superluminal motion arises when a relativistic blob of emitting matter is ejected close to our line of sight to the object.



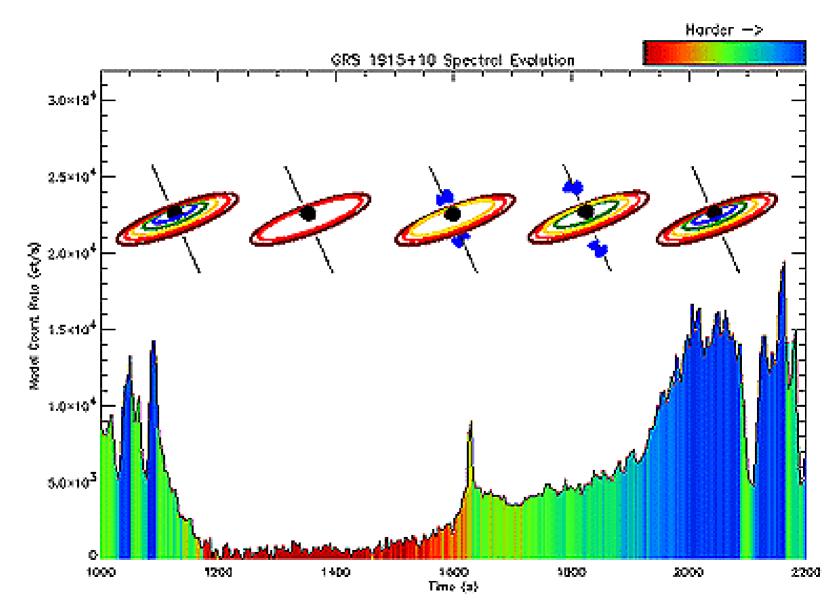


The radio jets of the Galactic blackhole candidate GRS 1915+105

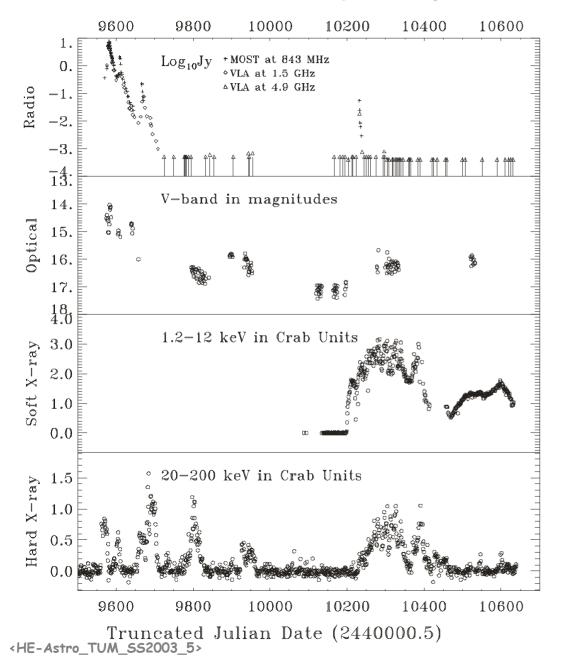
**Disk-Jet Connection (I)** 



### Disk-Jet Connection (II)



### **Multifrequency** Observations



Flaring episodes of the transient GRO J1655-40, which exhibits super-luminal radio jets. Strong hard X-ray and gamma-ray flares not accompanied by radio flares indicate that the high-energy emission is not produced in the jets.