

Stability of Mass Transfer in Eccentric **Compact Binaries**



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Abstract Compact binaries are important progenitors for gamma-ray bursts. Here we present simulations of the onset of mass transfer in compact binaries, focusing on systems containing a neutron star and a white dwarf. We make use of a modified form of smoothed particle hydrodynamics enabling us to model realistically-low mass transfer rates.

Eccentric	WD-NS	Binaries:	Introd	uction

Before the mass transfer (see, for example, Davies et al., 2002):

- ▶ Birth rate of $10^{-4} 10^{-5} \text{yr}^{-1}$ per galaxy
- ► NS is formed second
- ► There are two observed WD-NS binaries, and two proposed subpopulations

During the mass transfer:

- Merger of WD-NS binaries may produce GRBs and SNe (King et al,
 - 2007)
- GRBs possibly originating from these systems have been observed (Gal-Yam et al., 2006; Della Valle et al., 2006; Fynbo et al., 2006) Nuclear burning may be important during the merger (Metzger, 2011)

Notation:	MS – a	main-sec	uence star
			Juchec Stur

- ► WD a white dwarf
- ► NS a neutron star
- ► GRB Gamma-ray burst
- ► SN a supernova
- ► MT mass transfer

Eccentric WD-NS Binaries Formation

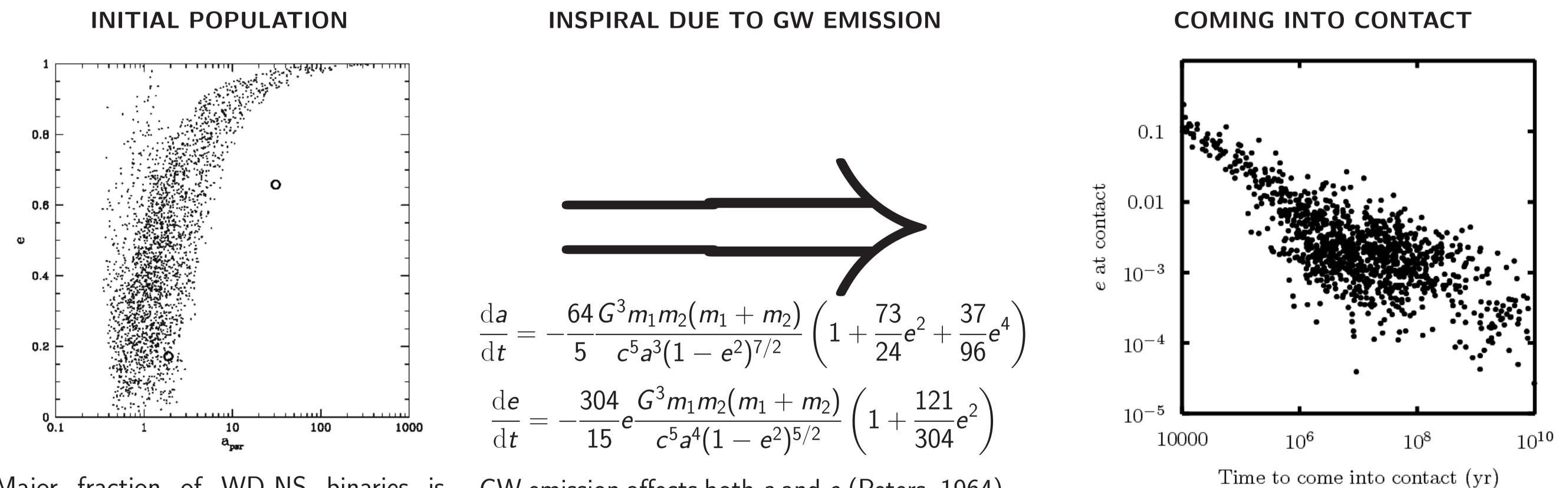
- Primary expands, MT starts
- Only He core left from the primary
- The primary expands and MT starts again
- The primary turns into a WD
- Secondary reaches the red giant phase, and fills its Roche lobe
- Common envelope phase starts, the system gets tight ► He secondary evolves, leading to either MT or mass loss through winds: two populations form

- Merger rate is observationally interesting
- When they come into contact, WD-NS binaries are still interestingly eccentric
- WD-NS binaries make a promising source for GWs

We study the stability of mass transfer in WD-NS binaries

- Secondary explodes as a SN (assuming it has gained enough) mass)
- Newborn NS gets a kick binary becomes eccentric

How Do WD-NS Systems Come Into Contact?



Major fraction of WD-NS binaries is formed in tight and eccentric binaries. The observed systems J1141 + 6545 and B2303 + 46 are shown on the figure. The latter binary belongs to the other, more wide population.

$$\frac{\mathrm{d}e}{\mathrm{d}t} = -\frac{304}{15}e\frac{G^3m_1m_2(m_1+m_2)}{c^5a^4(1-e^2)^{5/2}}\left(1+\frac{121}{304}e^2\right)$$

GW emission affects both *a* and *e* (Peters, 1964). 95 percent of the presented population shall merge in less than a Hubble time.

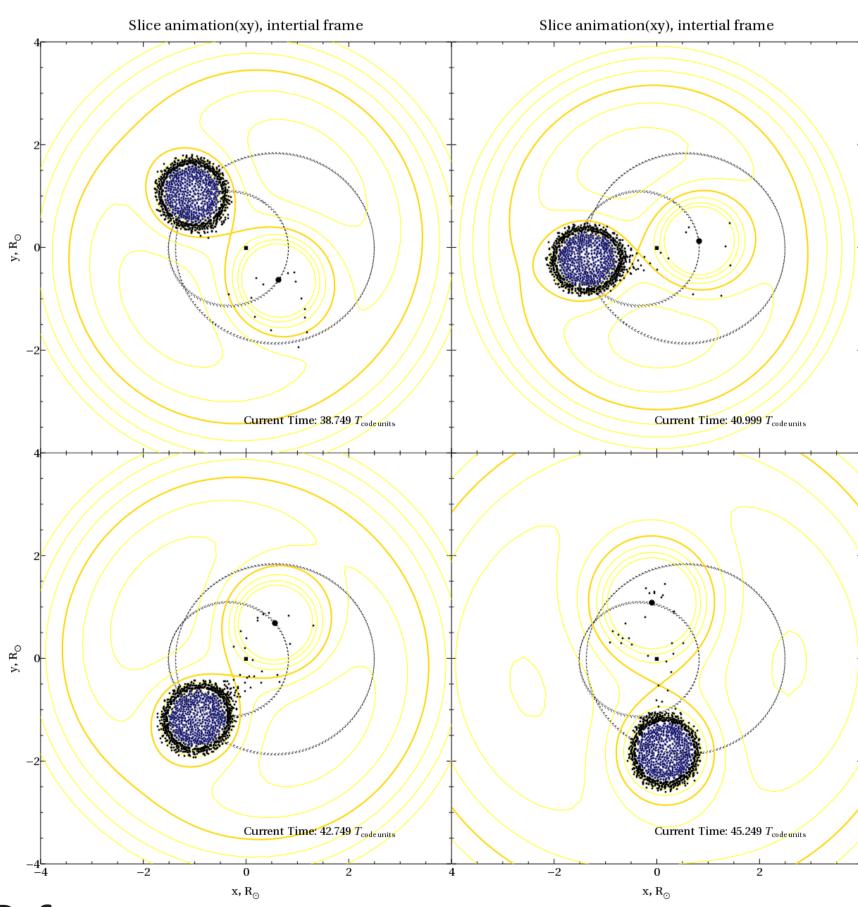
At the moment of coming into contact the binaries are still interestingly eccentric: variation of the binary separation is of order of scale height of the WD. Consequences: periodic mass transfer, Roche lobe formalism is not applicable.

Modelling the Mass Transfer

MODIFIED SPH

Motivation:

- SPH method replaces the bodies by sets of particles of comparable masses
- ► If the binary were circular, the mass transfer rate would be down to of order $10^{-12} M_{\odot}$ per period (Rosswog & Bruggen, 2003). • Hence one would need of order 10^{12} SPH particles to resolve the mass transfer Oil-on-water scheme (Church et al, 2009):



SIMULATIONS

A simulation of $\gamma = 5/3$ polytrope $0.6 M_{\odot}$ star orbiting a $1M_{\odot}$ compact companion, with a resolved phase of episodic mass transfer happening between the stars. We use the units, in which $M_{\odot} = 1, R_{\odot} = 1$ 1, G = 1. The binary separation at its minimum is $2.2R_{\odot}$, the eccentricity is e =0.29.

- The main idea is to artificially separate the atmosphere and the body of the star
- This allows one to use two types of SPH particles of very different masses in a single simulation
- Hence one can resolve realistically low mass transfer rates.

References

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