

# WMAP 5-Year Results: Implications for Dark Energy

**Eiichiro Komatsu** (Department of Astronomy, UT Austin)  
3rd Biennial Leopoldina Conference, October 9, 2008

# WMAP 5-Year Papers

- **Hinshaw et al.**, “*Data Processing, Sky Maps, and Basic Results*” [0803.0732](#)
- **Hill et al.**, “*Beam Maps and Window Functions*” [0803.0570](#)
- **Gold et al.**, “*Galactic Foreground Emission*” [0803.0715](#)
- **Wright et al.**, “*Source Catalogue*” [0803.0577](#)
- **Nolta et al.**, “*Angular Power Spectra*” [0803.0593](#)
- **Dunkley et al.**, “*Likelihoods and Parameters from the WMAP data*” [0803.0586](#)
- **Komatsu et al.**, “*Cosmological Interpretation*” [0803.0547](#)

# WMAP 5-Year Science Team

- C.L. Bennett
- G. Hinshaw
- N. Jarosik
- S.S. Meyer
- L. Page
- D.N. Spergel
- E.L. Wright
- M.R. Greason
- M. Halpern
- R.S. Hill
- A. Kogut
- M. Limon
- N. Odegard
- G.S. Tucker
- J. L. Weiland
- E. Wollack
- J. Dunkley
- B. Gold
- E. Komatsu
- D. Larson
- M.R.olta

Special  
Thanks to  
**WMAP**  
**Graduates!**

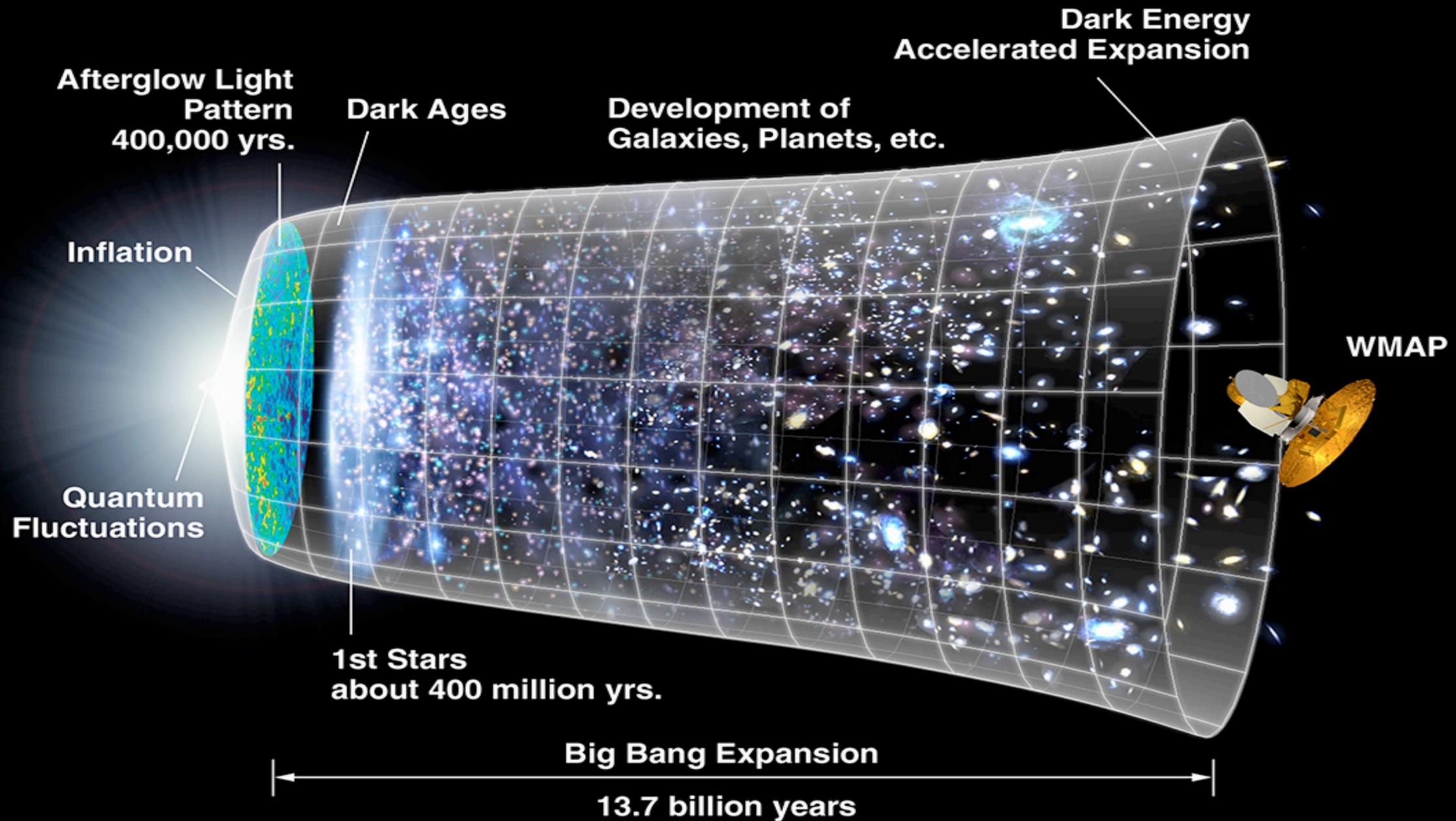
- C. Barnes
- R. Bean
- O. Dore
- H.V. Peiris
- L. Verde

# Need For Dark “Energy”

- First of all, DE does not even need to be energy.
- At present, *anything* that can explain the observed
  - (1) **Luminosity Distances** (Type Ia supernovae)
  - (2) **Angular Diameter Distances** (BAO, CMB)

*simultaneously* is qualified for being called “Dark Energy.”
- The candidates in the literature include: (a) energy, (b) modified gravity, and (c) extreme inhomogeneity.
- Measurements of the (3) **growth of structure** break degeneracy. (The best data right now is the X-ray clusters.)

# Measuring Distances, $H(z)$ & Growth of Structure



# $H(z)$ : Current Knowledge

- $H^2(z) = H^2(0) [\Omega_r(1+z)^4 + \Omega_m(1+z)^3 + \Omega_k(1+z)^2 + \Omega_{de}(1+z)^{3(1+w)}]$
- (expansion rate)  $H(0) = 70.5 \pm 1.3 \text{ km/s/Mpc}$
- (radiation)  $\Omega_r = (8.4 \pm 0.3) \times 10^{-5}$
- (matter)  $\Omega_m = 0.274 \pm 0.015$
- (curvature)  $\Omega_k < 0.008$  (95%CL)
- (dark energy)  $\Omega_{de} = 0.726 \pm 0.015$
- (DE equation of state)  $1+w = -0.006 \pm 0.068$

# H(z) to Distances

- Comoving Distance
  - $\chi(z) = c \int^z [dz'/H(z')]$
- Luminosity Distance
  - $D_L(z) = (1+z)\chi(z) [1 - (k/6)\chi^2(z)/R^2 + \dots]$
  - $R = (\text{curvature radius of the universe}); k = (\text{sign of curvature})$
  - WMAP 5-year limit:  $R > 2\chi(\infty)$ ; justify the Taylor expansion
- Angular Diameter Distance
  - $D_A(z) = [\chi(z)/(1+z)] [1 - (k/6)\chi^2(z)/R^2 + \dots]$

$$D_A(z) = (1+z)^{-2} D_L(z)$$

$D_L(z)$

Type Ia Supernovae

$D_A(z)$

Galaxies (BAO)

CMB

0.02

0.2

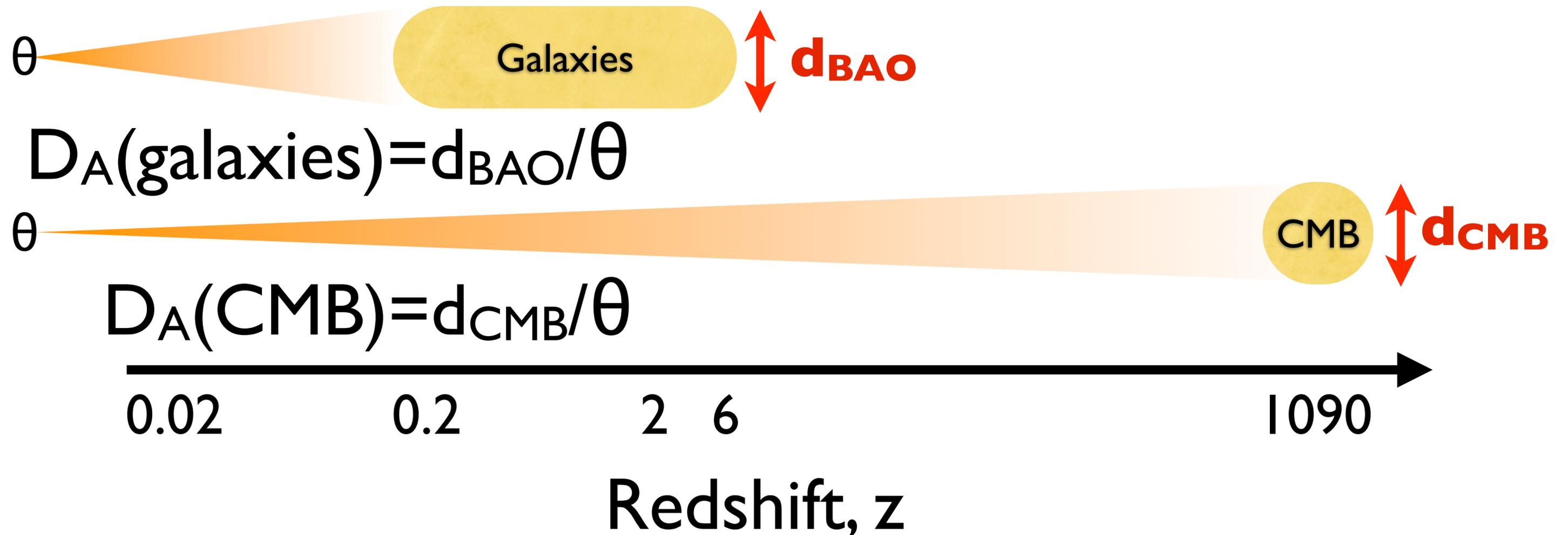
2 6

1090

Redshift,  $z$

- To measure  $D_A(z)$ , we need to know the intrinsic size.
- What can we use as the *standard ruler*?

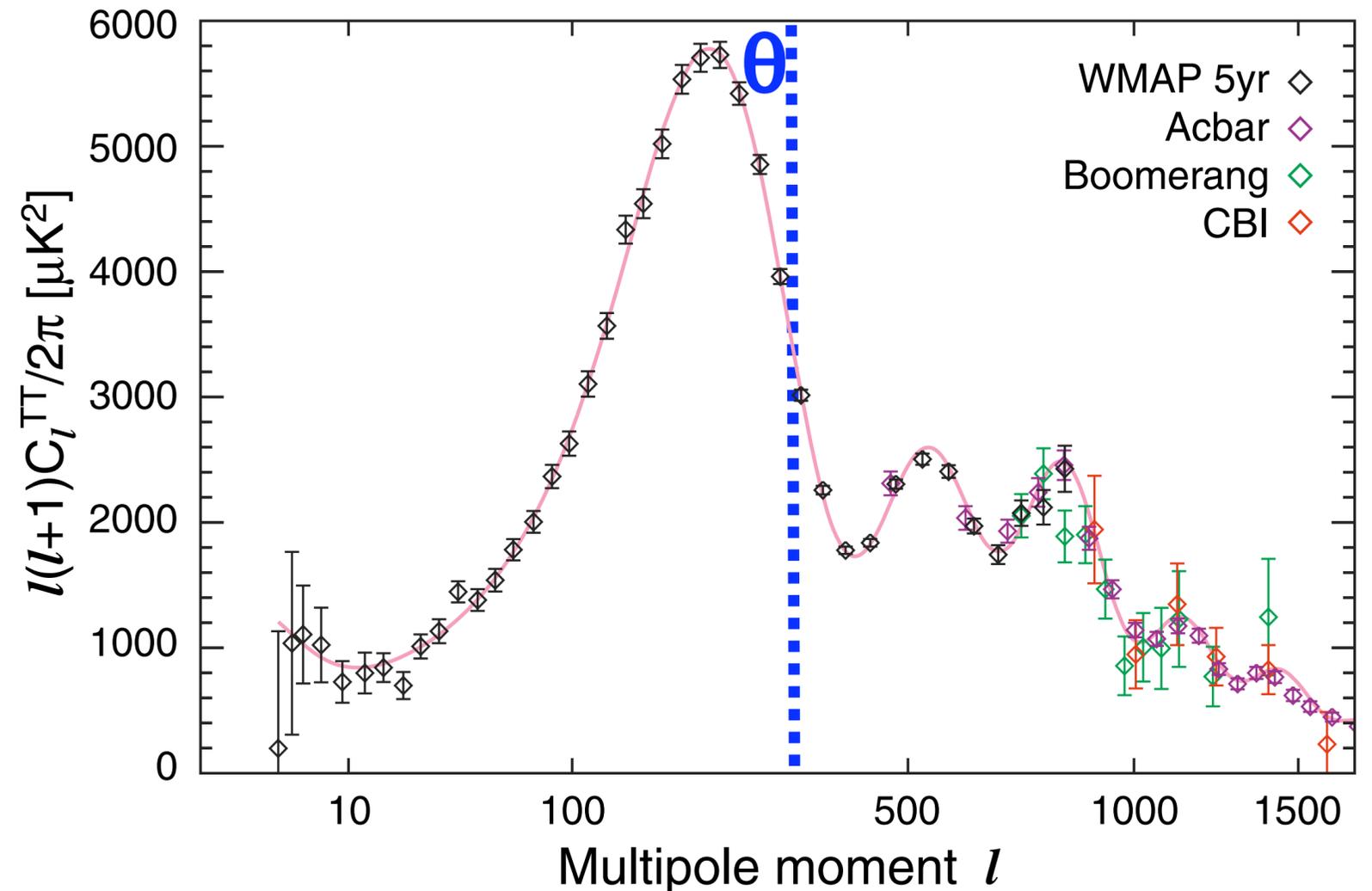
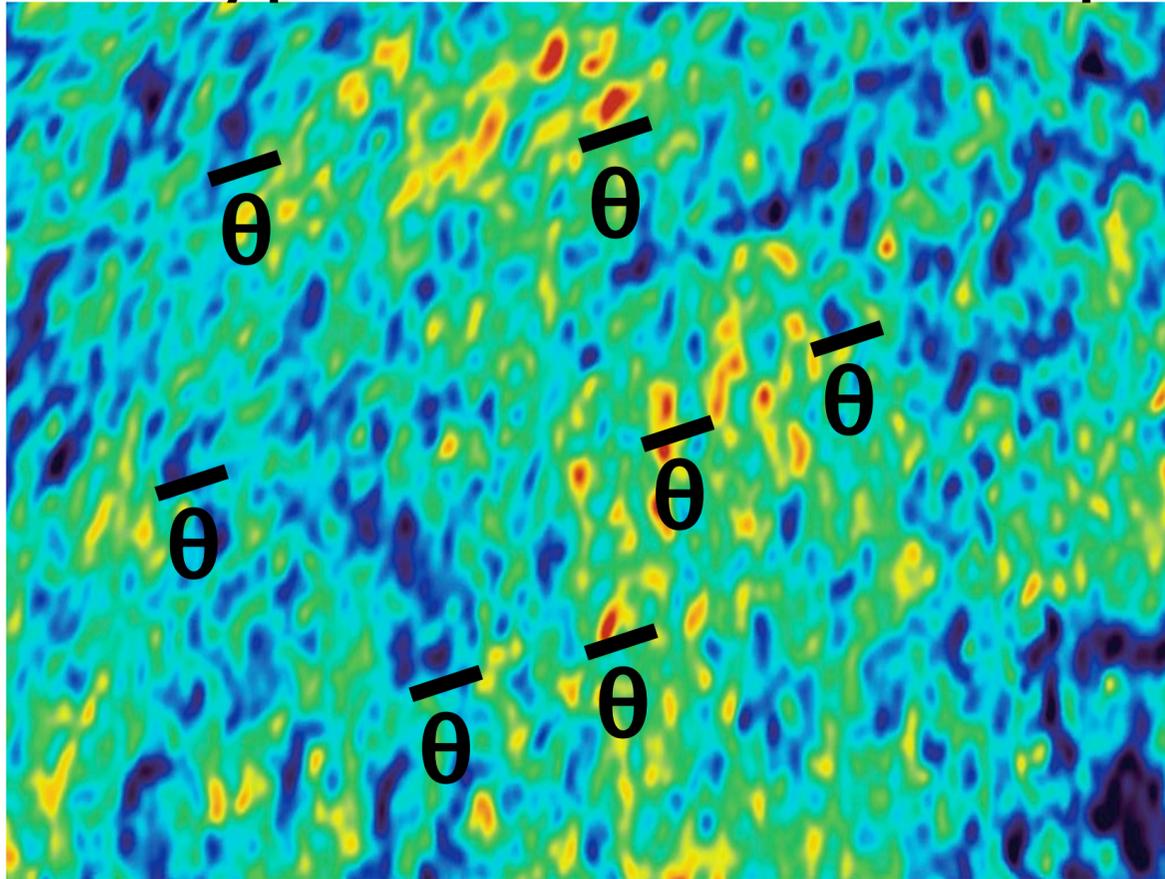
# How Do We Measure $D_A(z)$ ?



- If we know the intrinsic physical sizes,  $d$ , we can measure  $D_A$ . What determines  $d$ ?

# CMB as a Standard Ruler

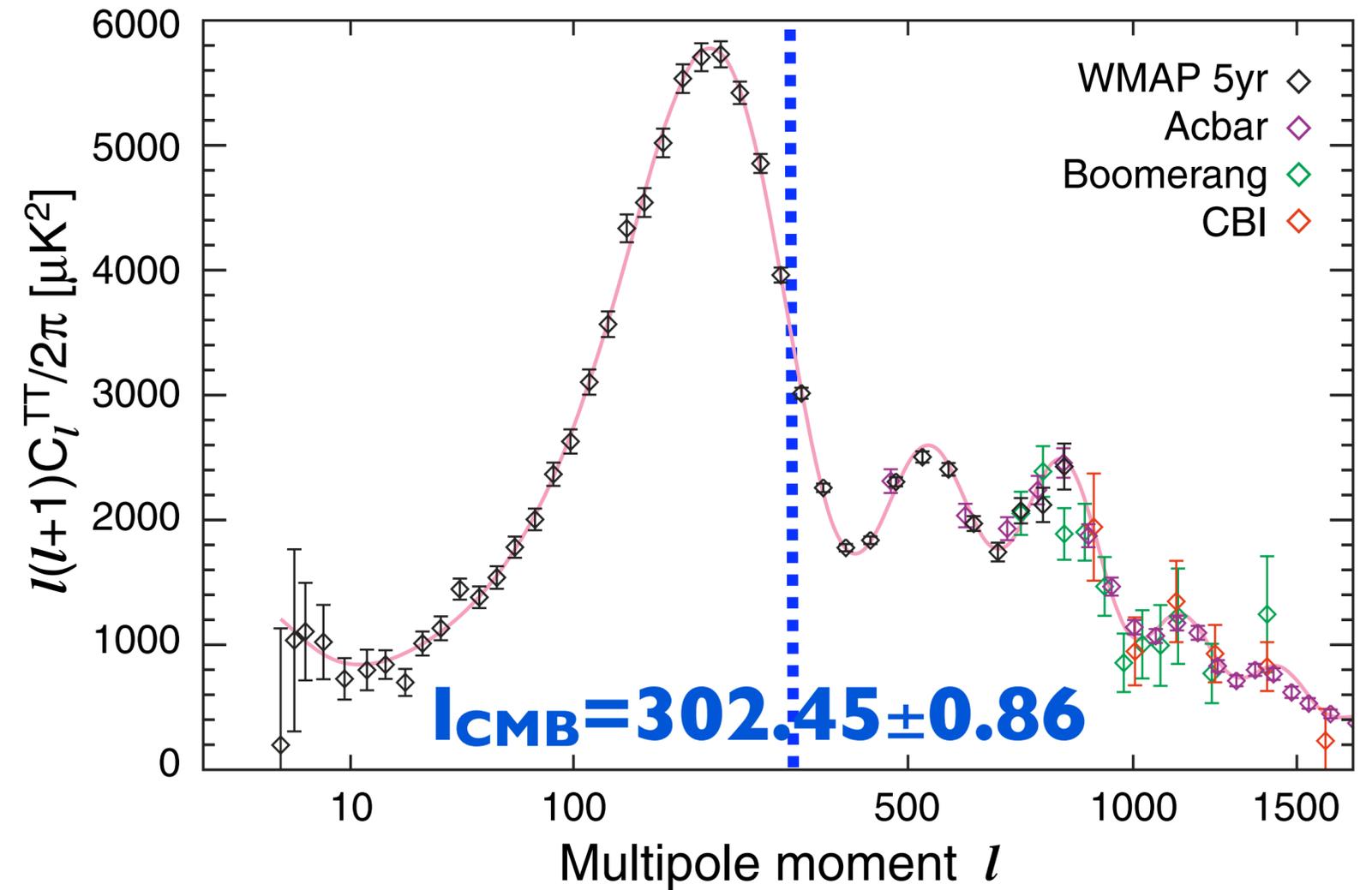
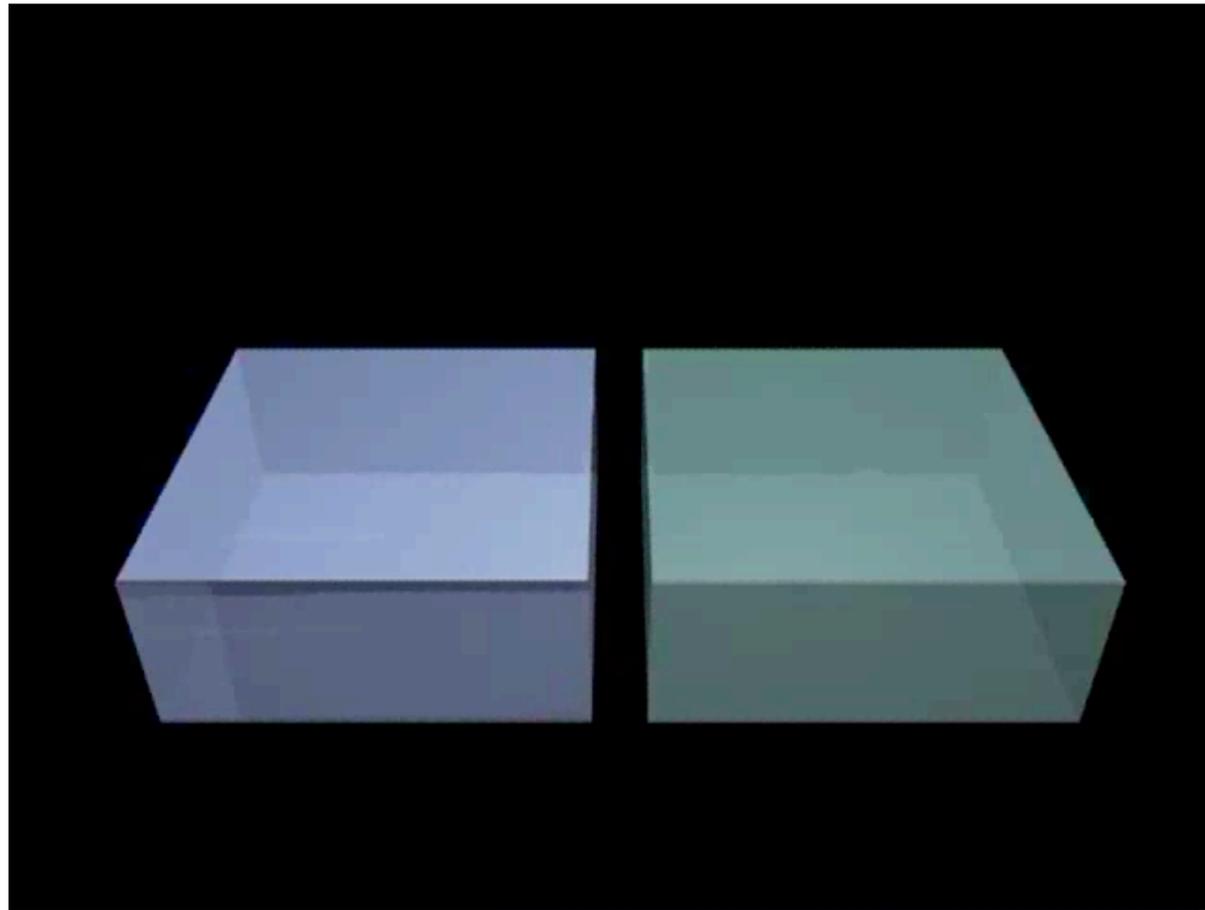
$\theta$  ~ the typical size of hot/cold spots



- The existence of typical spot size in image space yields oscillations in harmonic (Fourier) space. What determines the physical size of typical spots,  $d_{\text{CMB}}$ ? <sup>10</sup>

# Sound Horizon

- The typical spot size,  $d_{\text{CMB}}$ , is determined by the **physical distance traveled by the sound wave** from the Big Bang to the decoupling of photons at  $z_{\text{CMB}} \sim 1090$  ( $t_{\text{CMB}} \sim 380,000$  years).
- The causal horizon (photon horizon) at  $t_{\text{CMB}}$  is given by
  - $d_{\text{H}}(t_{\text{CMB}}) = a(t_{\text{CMB}}) * \text{Integrate}[c \, dt/a(t), \{t, 0, t_{\text{CMB}}\}]$ .
- The sound horizon at  $t_{\text{CMB}}$  is given by
  - $d_{\text{s}}(t_{\text{CMB}}) = a(t_{\text{CMB}}) * \text{Integrate}[c_{\text{s}}(t) \, dt/a(t), \{t, 0, t_{\text{CMB}}\}]$ , where  $c_{\text{s}}(t)$  is the time-dependent **speed of sound of photon-baryon fluid**.



- The WMAP 5-year values:

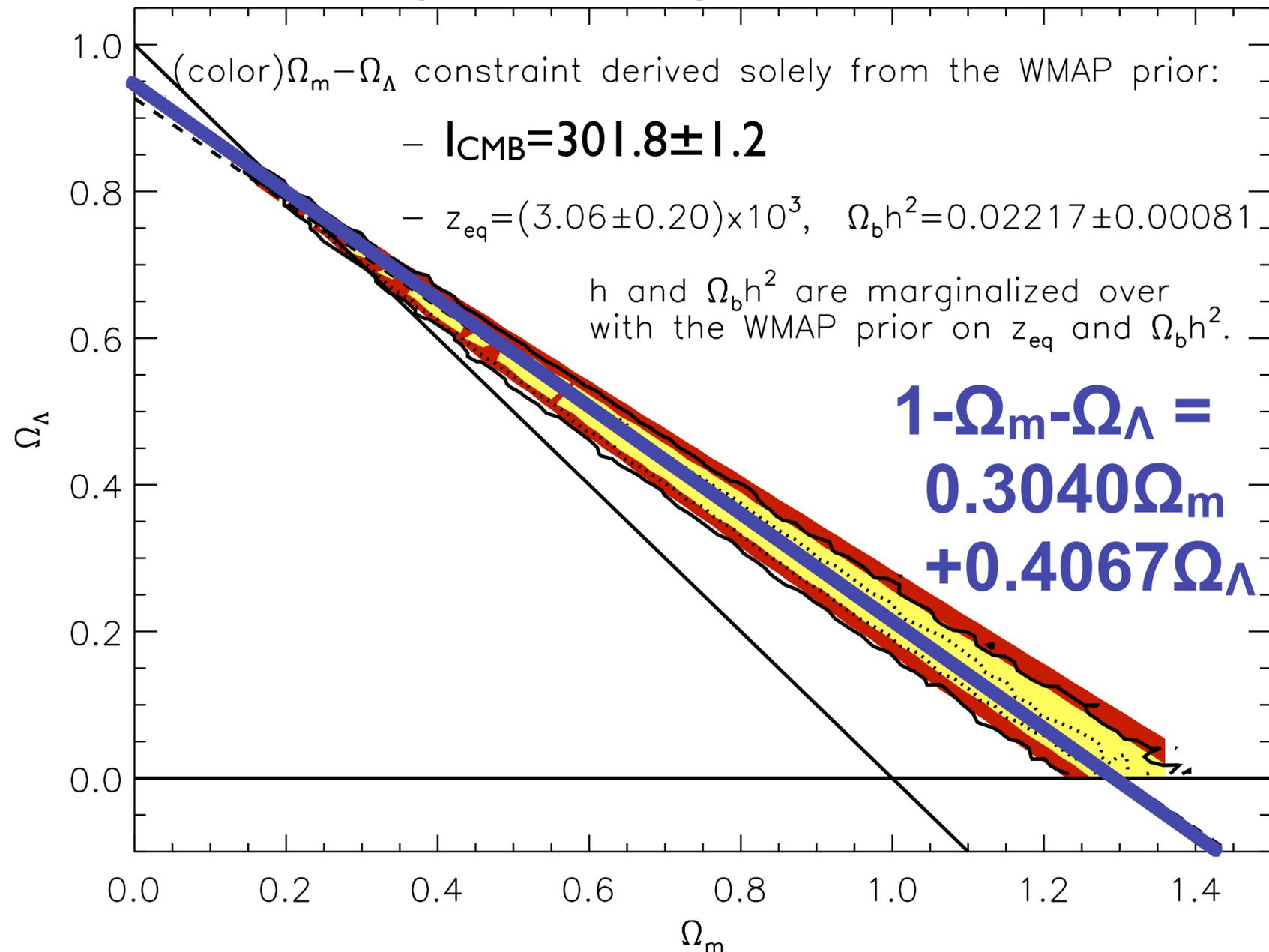
- $l_{\text{CMB}} = \pi/\theta = \pi D_A(z_{\text{CMB}})/d_s(z_{\text{CMB}}) = 302.45 \pm 0.86$

- CMB data constrain the ratio,  $D_A(z_{\text{CMB}})/d_s(z_{\text{CMB}})$ .

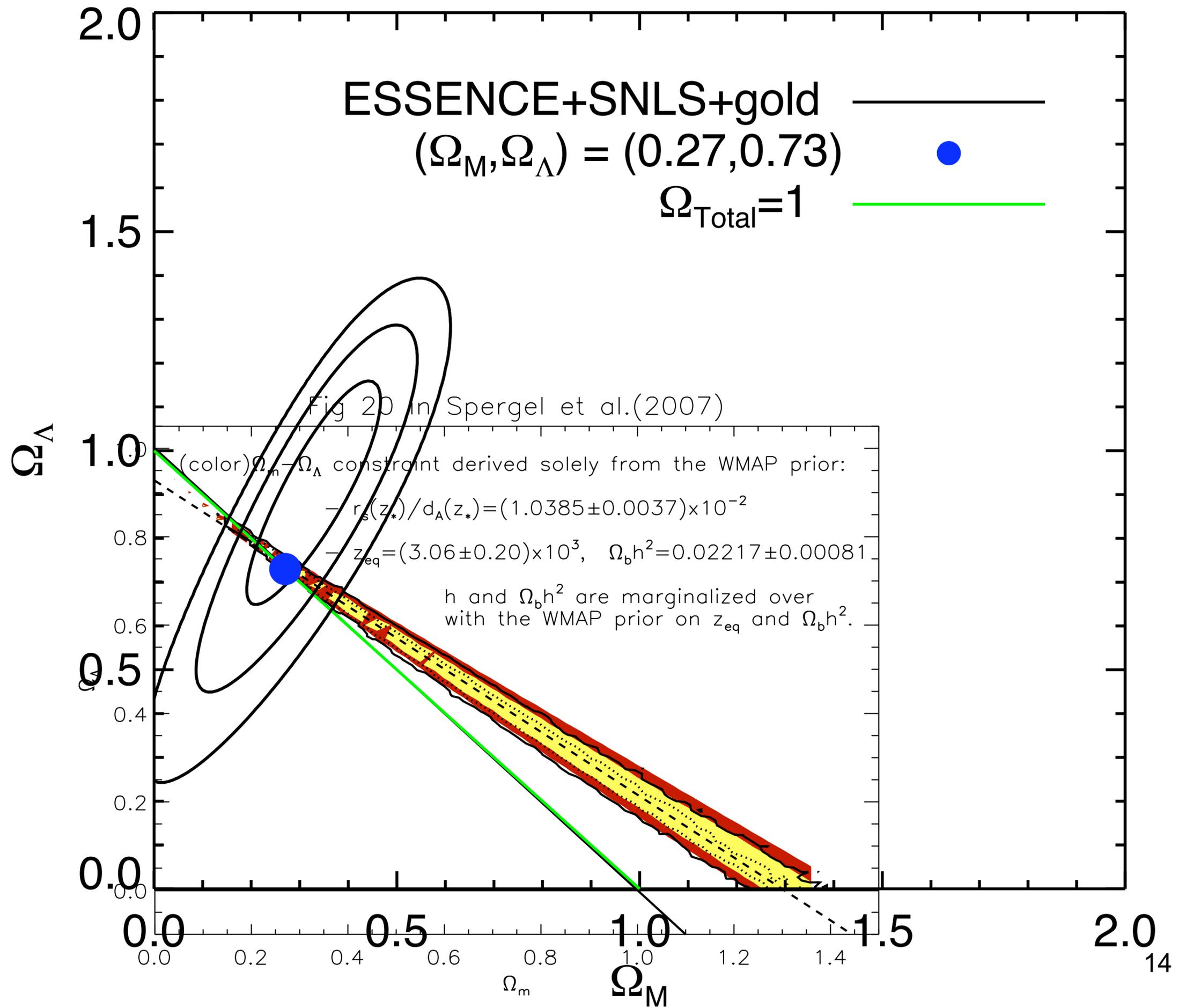
- $r_s(z_{\text{CMB}}) = (1+z_{\text{CMB}})d_s(z_{\text{CMB}}) = 146.8 \pm 1.8$  Mpc (comoving)

# What $D_A(z_{\text{CMB}})/d_s(z_{\text{CMB}})$ Gives You (3-year example)

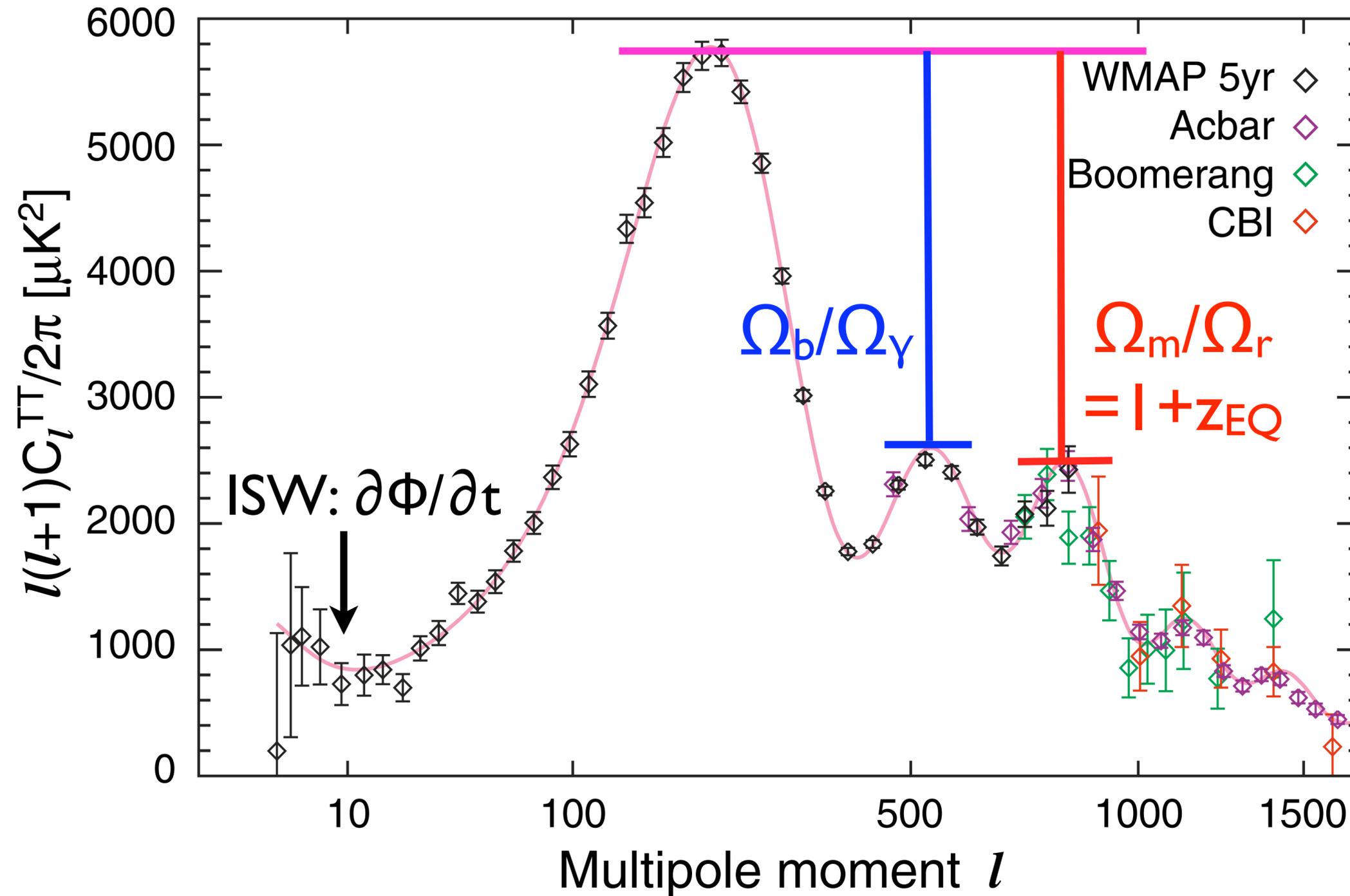
Fig 20 in Spergel et al.(2007)



- **Color**: constraint from  $l_{\text{CMB}} = \pi D_A(z_{\text{CMB}})/d_s(z_{\text{CMB}})$  with  $z_{\text{EQ}}$  &  $\Omega_b h^2$ .
- **Black contours**: Markov Chain from WMAP 3yr (Spergel et al. 2007)



# Other Observables



- $l$ -to-2: baryon-to-photon;  $l$ -to-3: matter-to-radiation ratio
- Low- $l$ : Integrated Sachs Wolfe Effect (more talks later!)

# Dark Energy From Distance Information Alone

- We provide a set of “WMAP distance priors” for testing various dark energy models.

$\Omega_b/\Omega_\gamma$

- Redshift of decoupling,  $z^*=1091.13$  (Err=0.93)

- Acoustic scale,  $l_A=\pi d_A(z^*)/r_s(z^*)=302.45$  (Err=0.86)

$\Omega_m/\Omega_r$

- Shift parameter,  $R=\text{sqrt}(\Omega_m H_0^2) d_A(z^*)=1.721$  (Err=0.019)

- Full covariance between these three quantities are also provided.

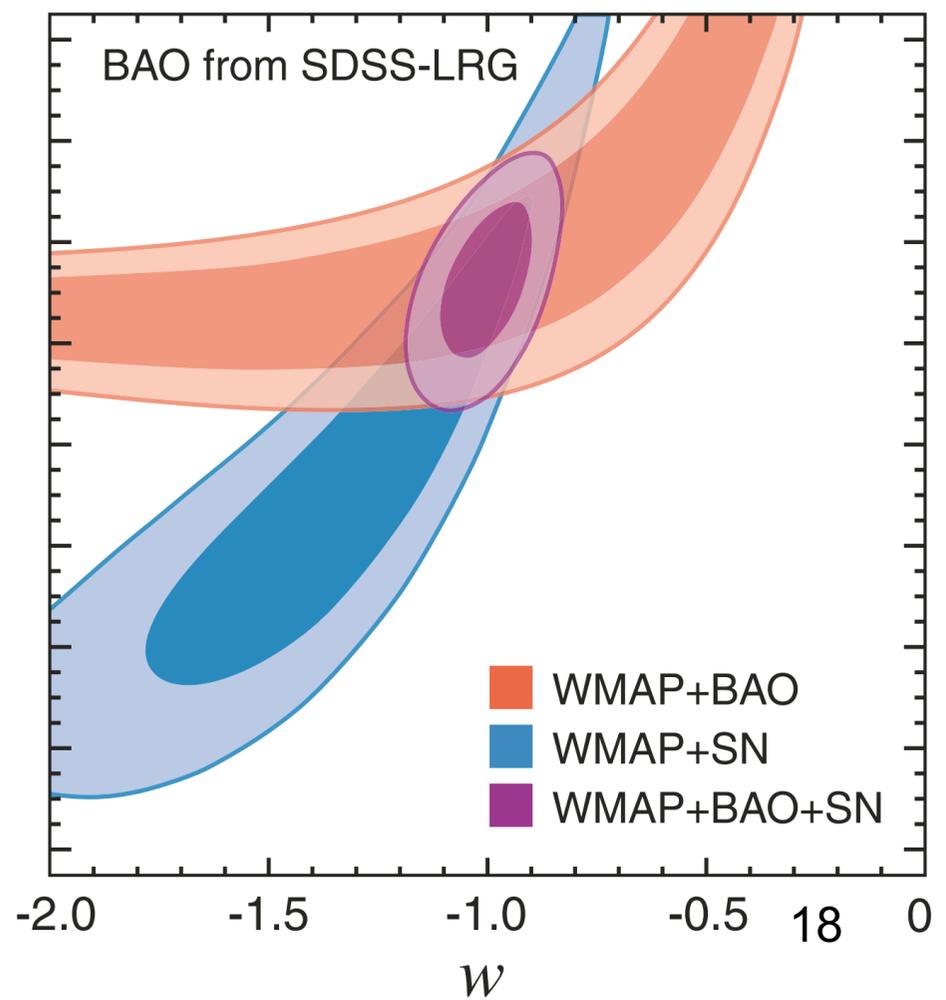
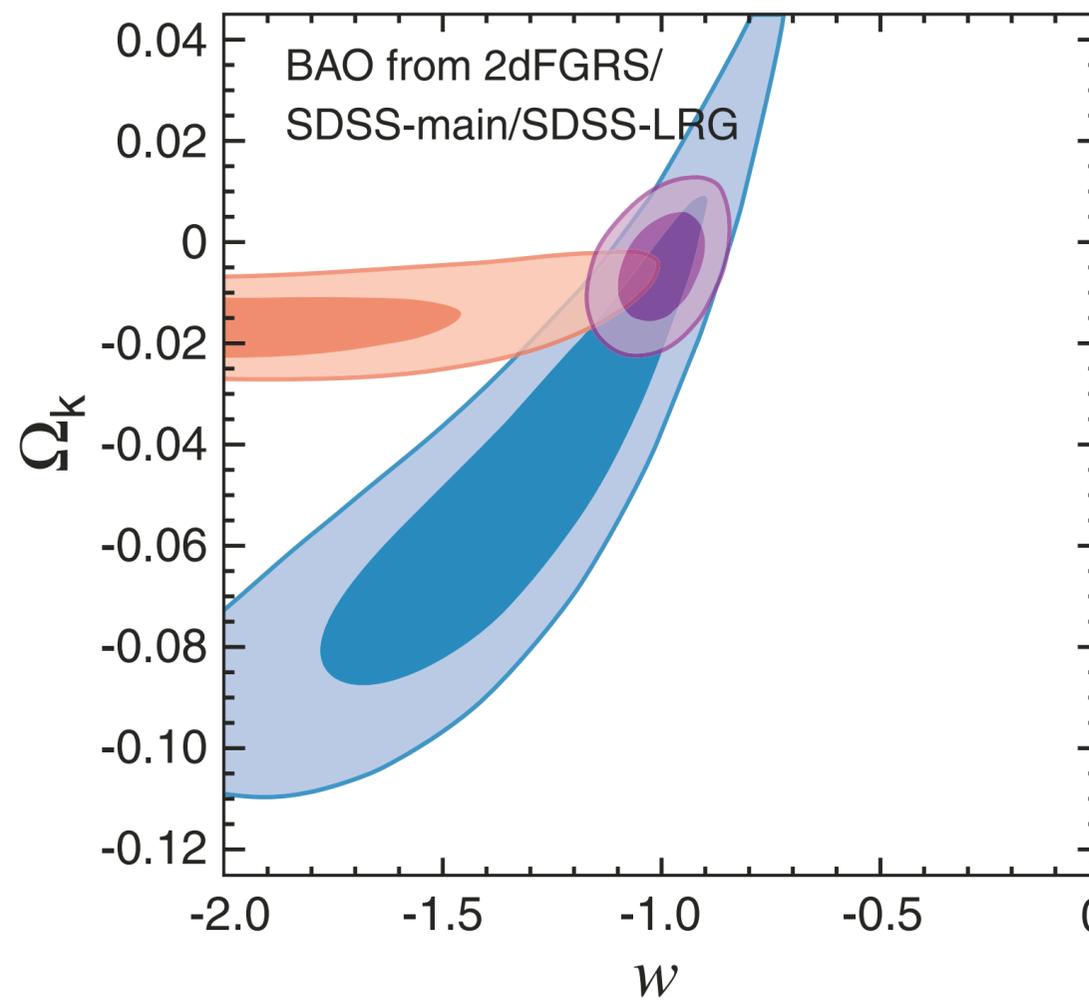
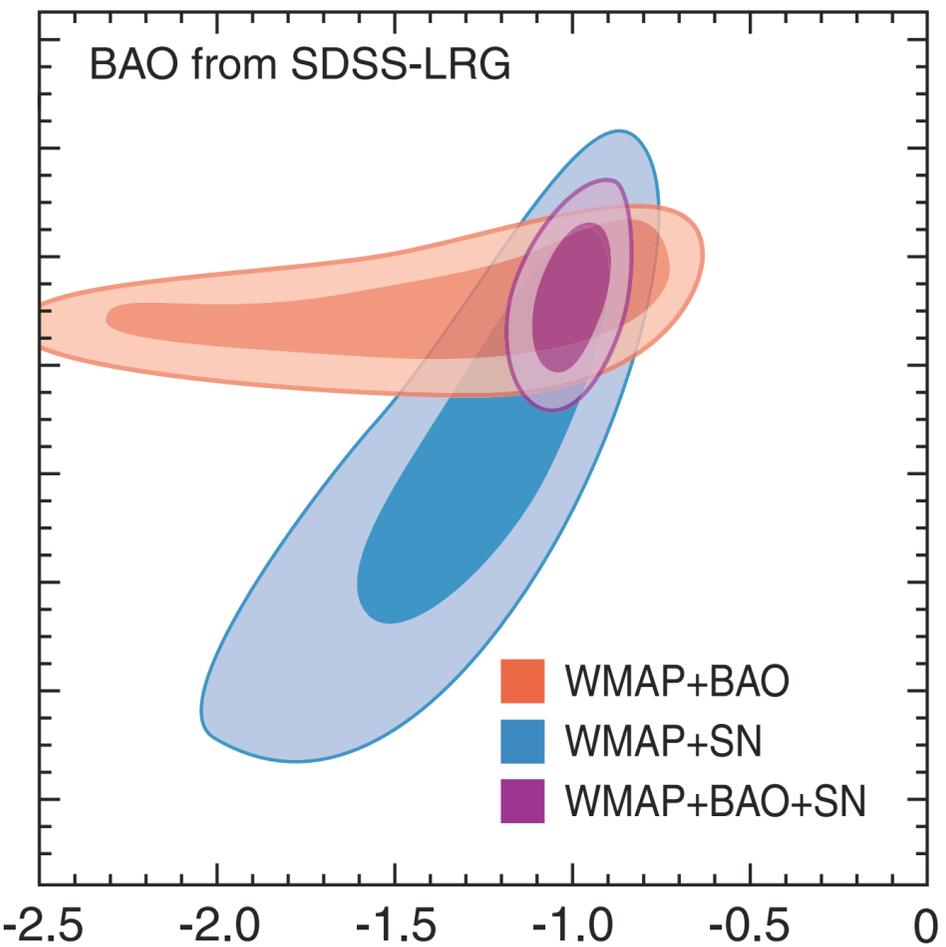
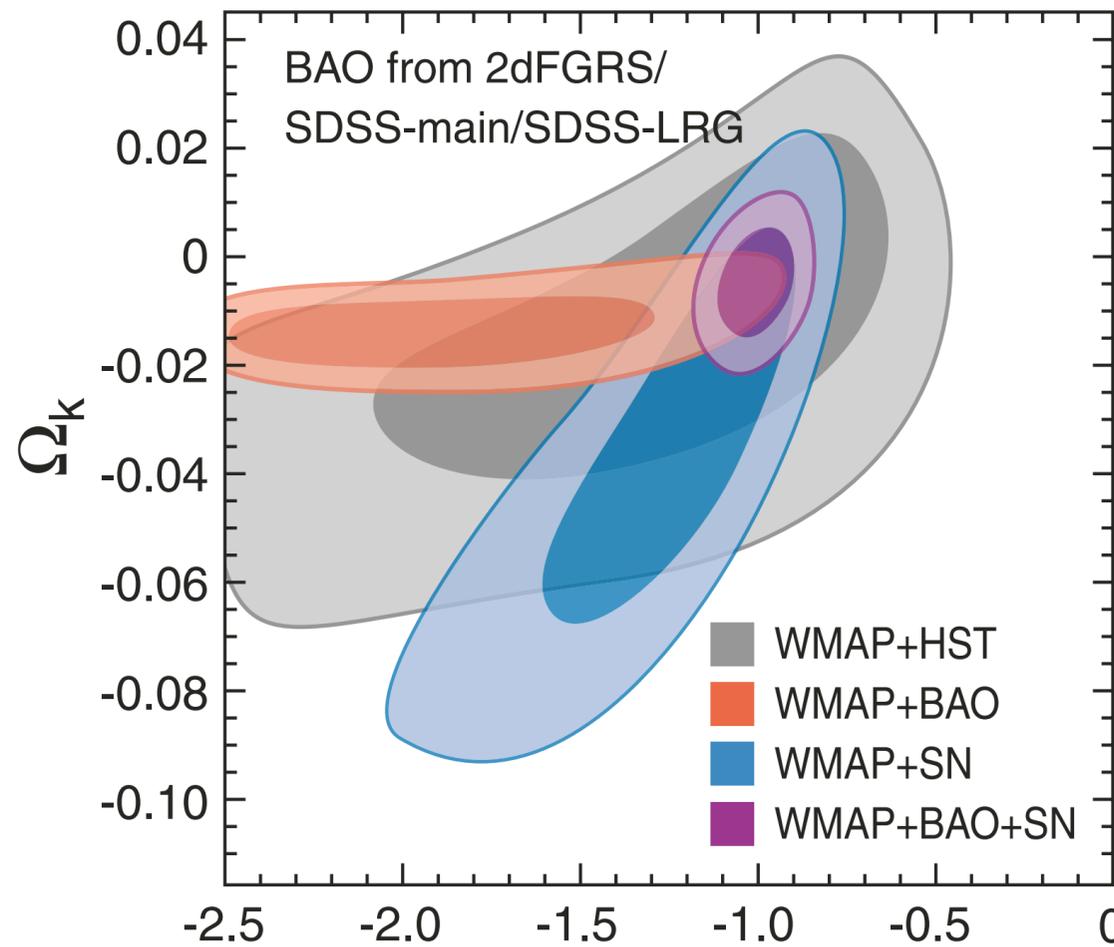
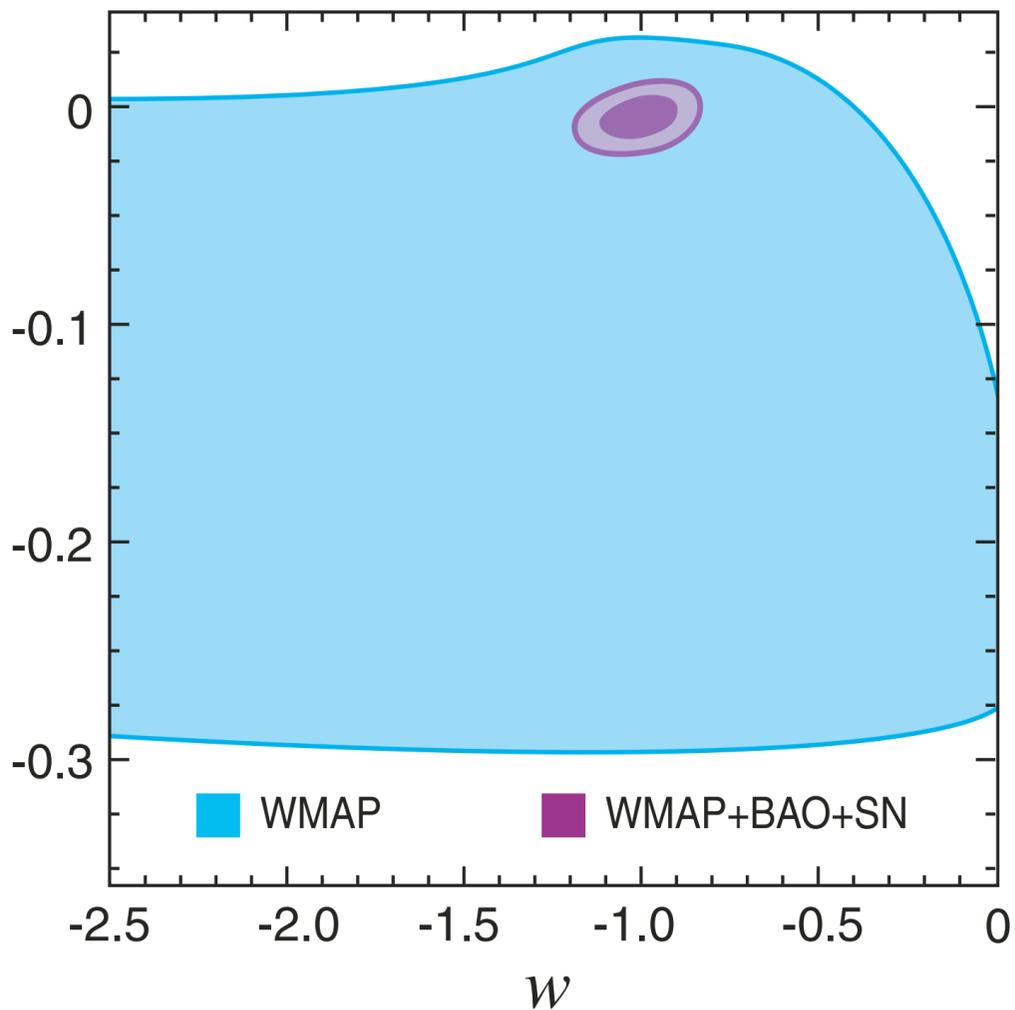
INVERSE COVARIANCE MATRIX FOR THE  
*WMAP* DISTANCE PRIORS

	$l_A(z_*)$	$R(z_*)$	$z_*$
$l_A(z_*)$	1.800	27.968	-1.103
$R(z_*)$		5667.577	-92.263
$z_*$			2.923

- **WMAP 5-Year ML**
- $z^*=1091.13$
- $l_A=302.45$
- $R=1.721$
- $100\Omega_b h^2=2.2765$

INVERSE COVARIANCE MATRIX FOR THE EXTENDED *WMAP*  
DISTANCE PRIORS. THE MAXIMUM LIKELIHOOD VALUE OF  $\Omega_b h^2$   
IS  $100\Omega_b h^2 = 2.2765$ .

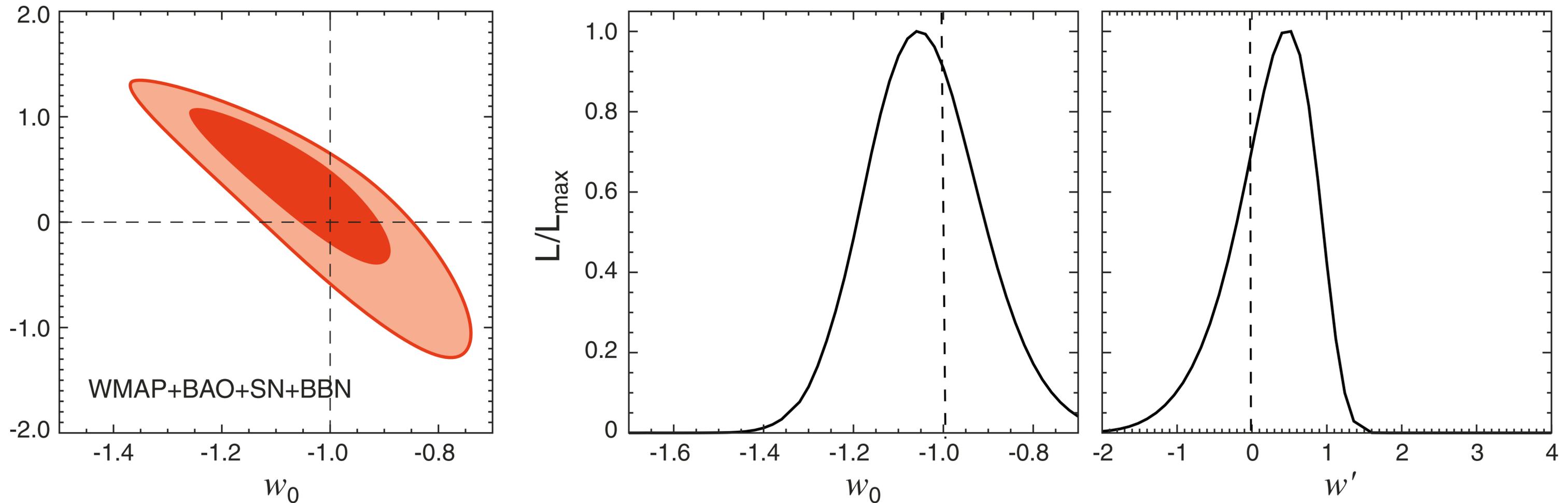
	$l_A(z_*)$	$R(z_*)$	$z_*$	$100\Omega_b h^2$
$l_A(z_*)$	31.001	-5015.642	183.903	2337.977
$R(z_*)$		876807.166	-32046.750	-403818.837
$z_*$			1175.054	14812.579
$100\Omega_b h^2$				187191.186



- Top
- Full WMAP Data
- Bottom
- WMAP Distance Priors

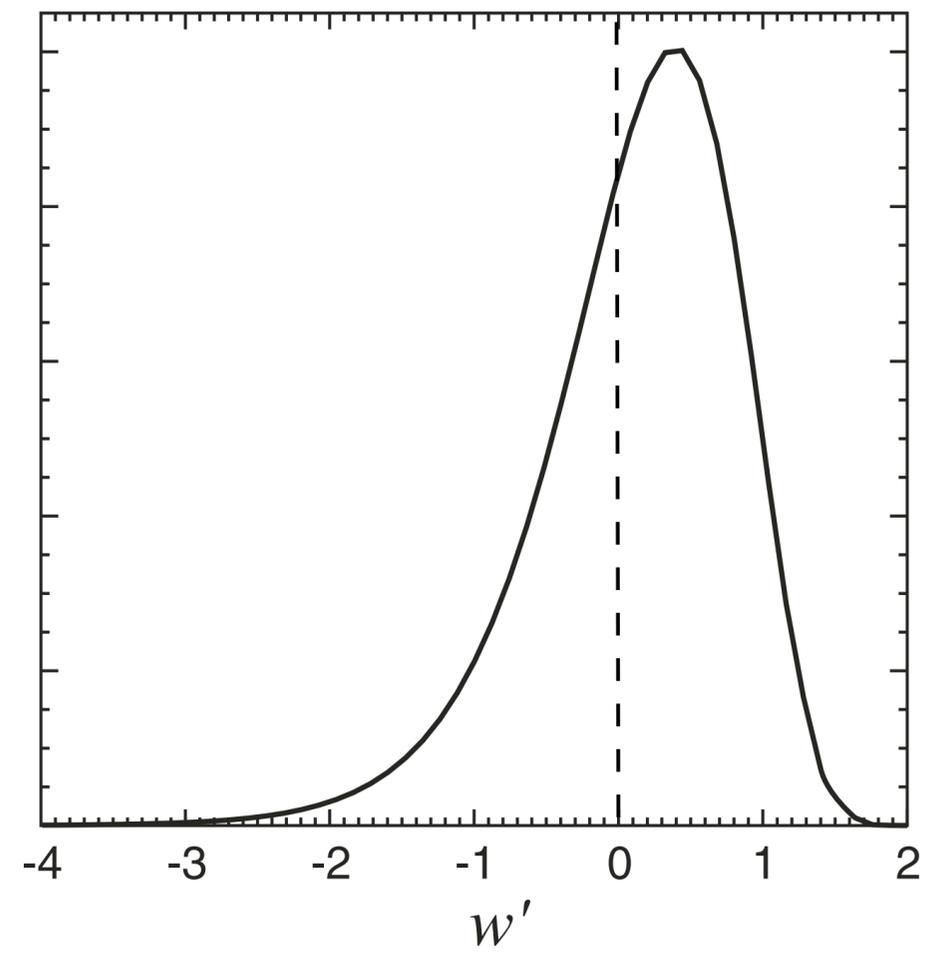
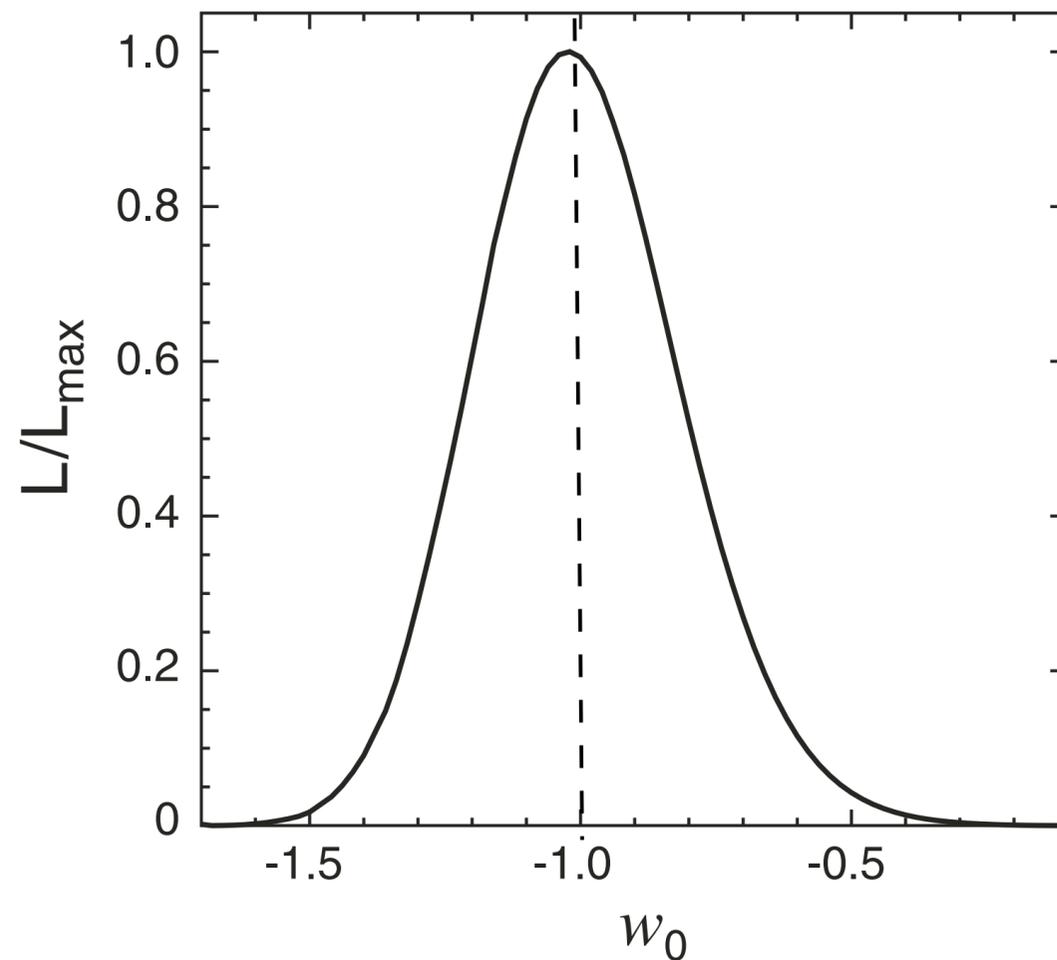
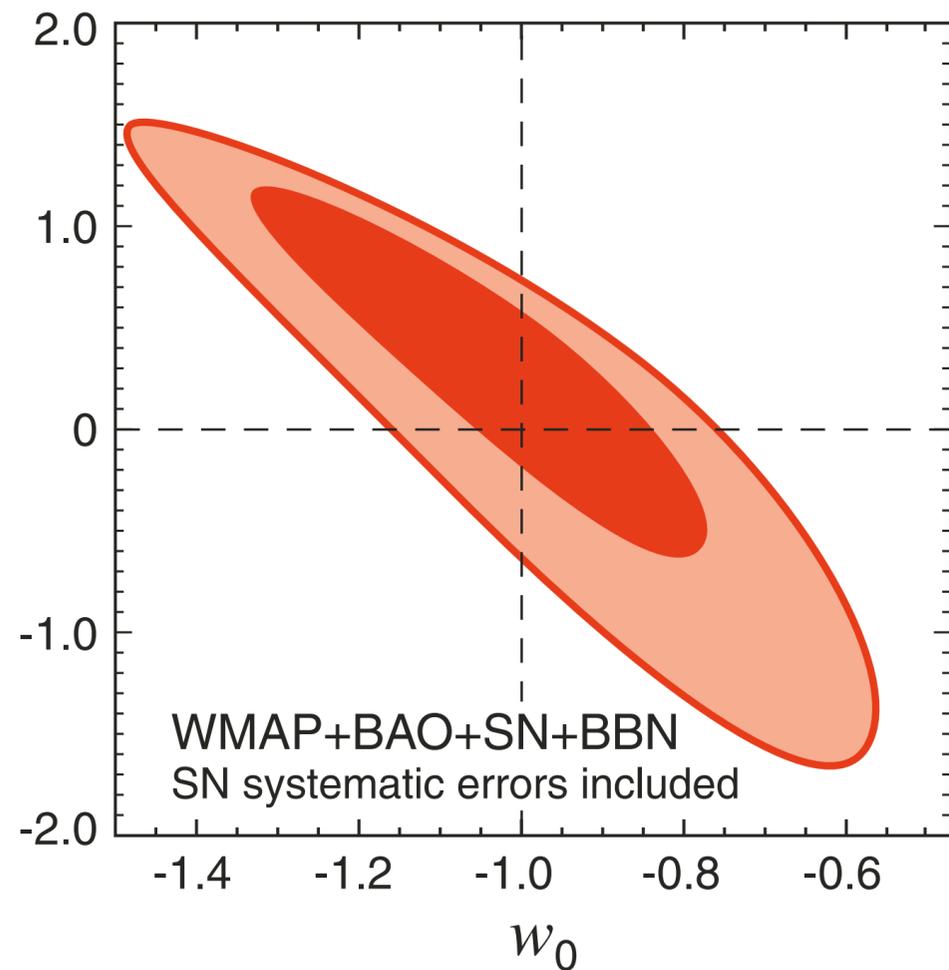
# Dark Energy EOS:

$$w(z) = w_0 + w'z / (1+z)$$



- Dark energy is pretty consistent with cosmological constant:  $w_0 = -1.04 \pm 0.13$  &  $w' = 0.24 \pm 0.55$  (68%CL)

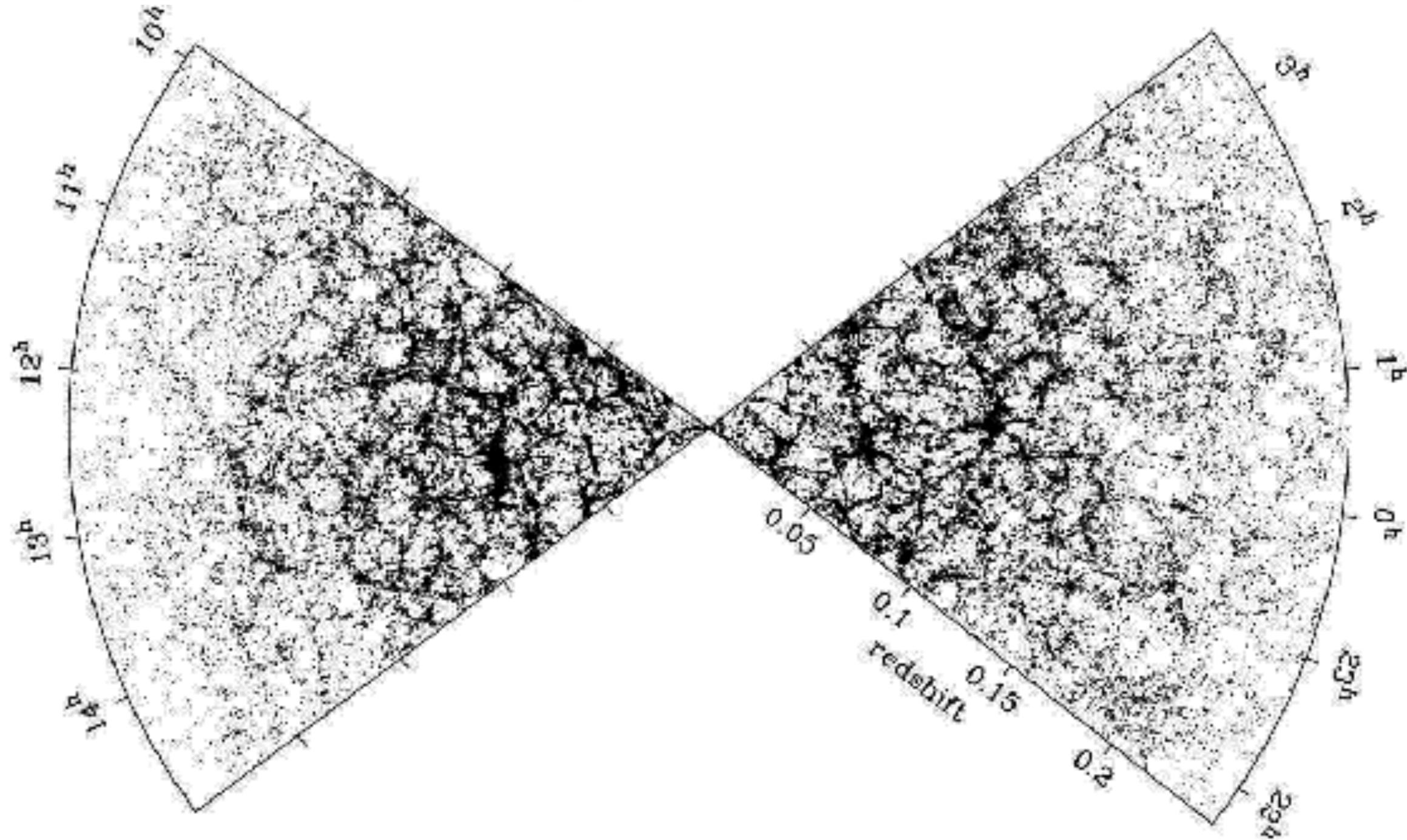
# Dark Energy EOS: <sup>WMAP5+BAO+SN</sup> Including Sys. Err. in SN Ia



- Dark energy is pretty consistent with cosmological constant:  $w_0 = -1.00 \pm 0.19$  &  $w' = 0.11 \pm 0.70$  (68%CL)

# BAO in Galaxy Distribution

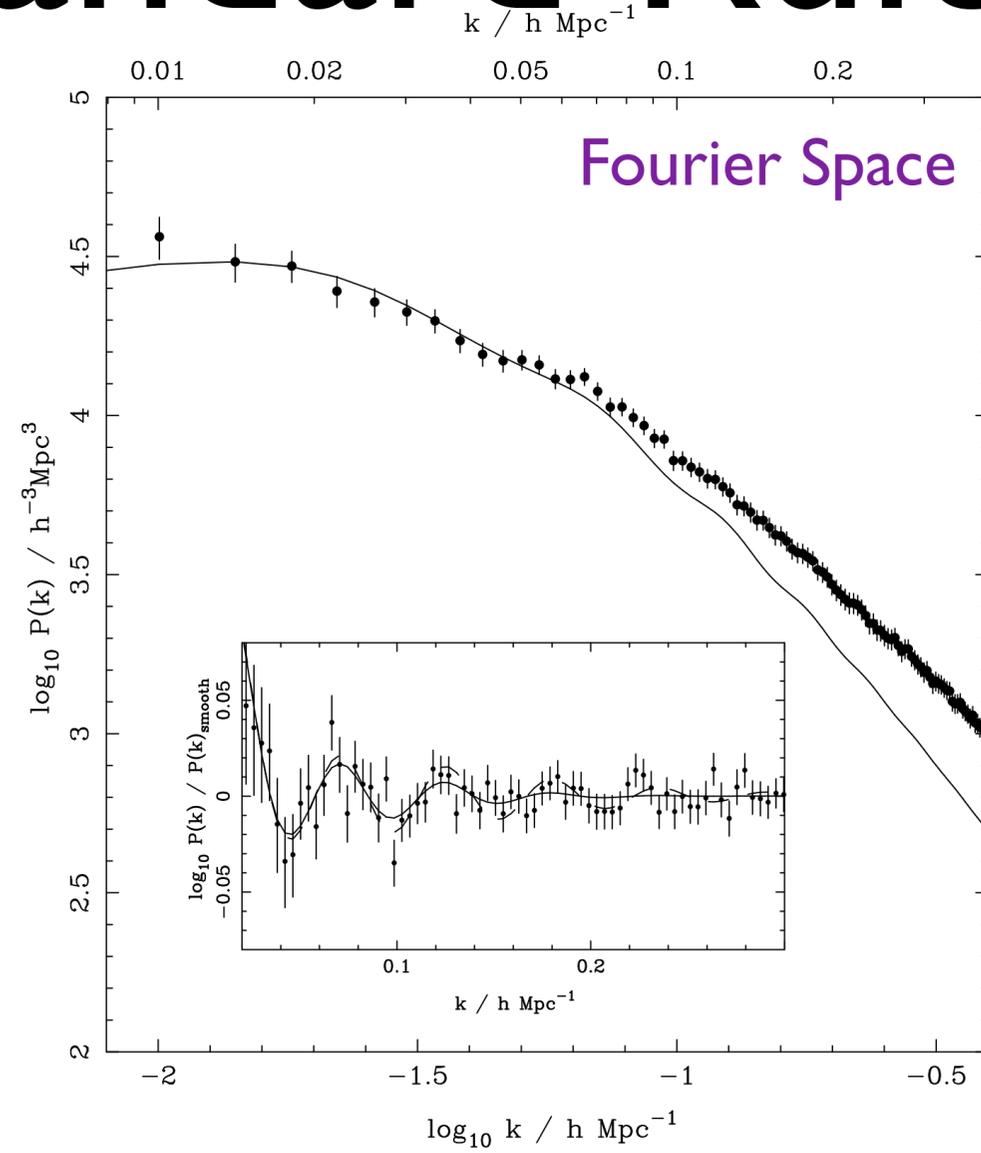
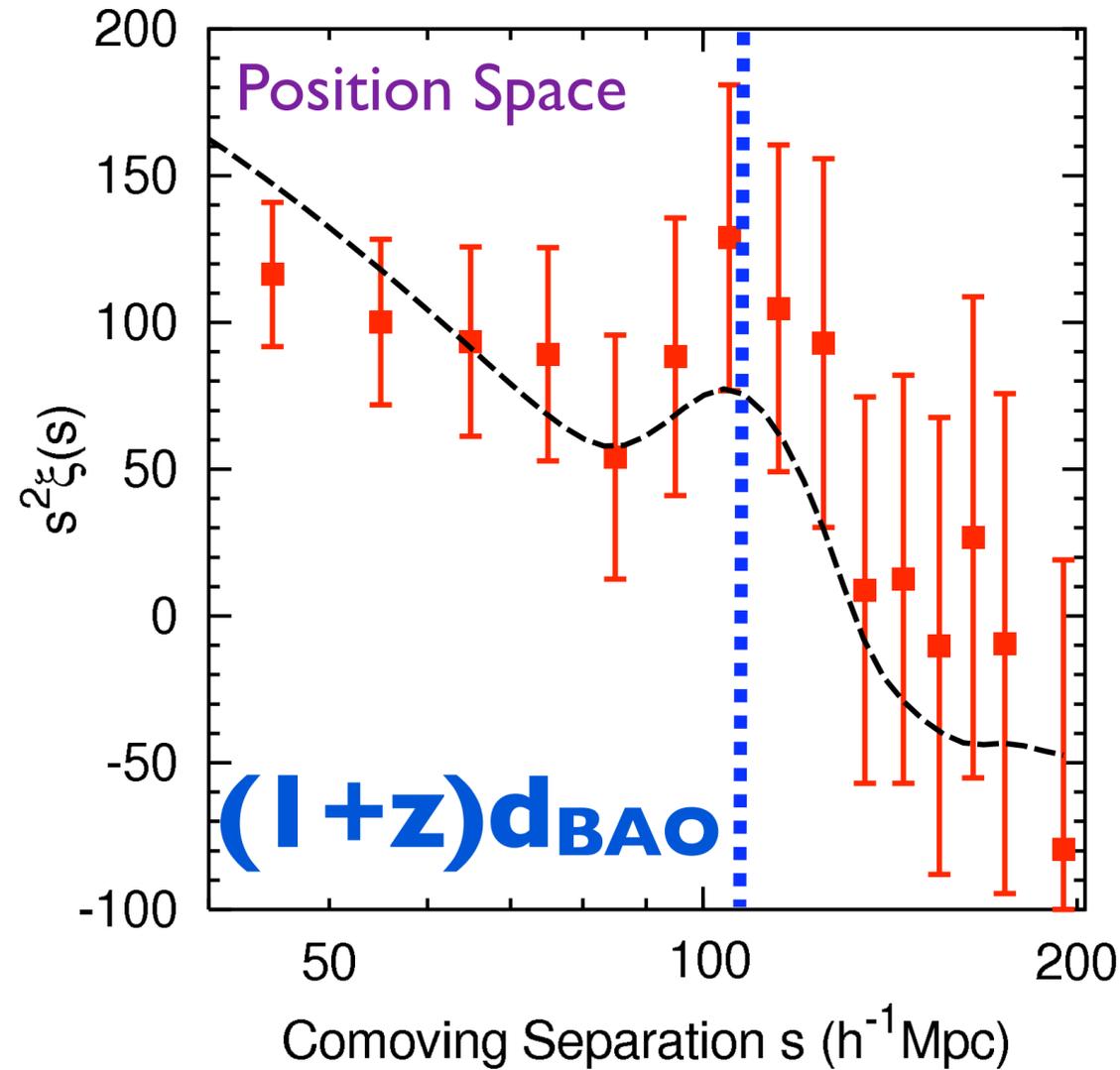
2dFGRS



- The same acoustic oscillations should be hidden in this galaxy distribution...

# BAO as a Standard Ruler

Okumura et al. (2007)



Percival et al. (2006)

- The existence of a localized clustering scale in the 2-point function yields oscillations in Fourier space. What determines the physical size of clustering,  $d_{\text{BAO}}$ ?

# Sound Horizon Again

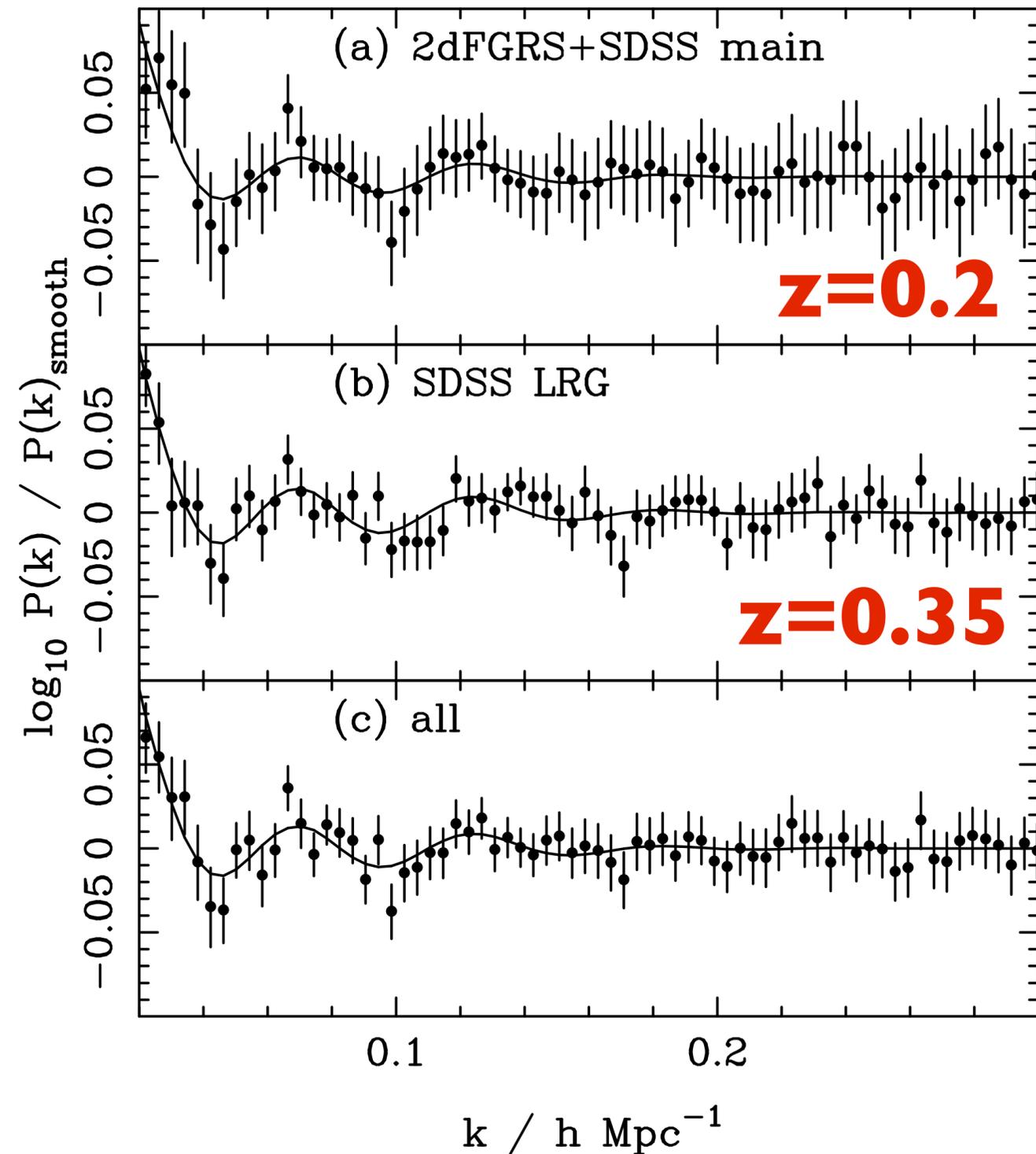
- The clustering scale,  $d_{\text{BAO}}$ , is given by the physical distance traveled by the sound wave from the Big Bang to the **decoupling of baryons** at  $z_{\text{BAO}} = 1020.5 \pm 1.6$  (c.f.,  $z_{\text{CMB}} = 1091 \pm 1$ ).
- The baryons decoupled slightly later than CMB.
  - By the way, this is not universal in cosmology, but *accidentally* happens to be the case for our Universe.
  - If  $3\rho_{\text{baryon}}/(4\rho_{\text{photon}}) = 0.64(\Omega_b h^2/0.022)(1090/(1+z_{\text{CMB}}))$  is greater than unity,  $z_{\text{BAO}} > z_{\text{CMB}}$ . Since our Universe happens to have  $\Omega_b h^2 = 0.022$ ,  $z_{\text{BAO}} < z_{\text{CMB}}$ . (ie,  $d_{\text{BAO}} > d_{\text{CMB}}$ )

# Standard Rulers in CMB & Matter

	Quantity	Eq.	5-year WMAP
CMB	$z_*$	(66)	$1090.51 \pm 0.95$
CMB	$r_s(z_*)$	(6)	$146.8 \pm 1.8$ Mpc
Matter	$z_d$	(3)	$1020.5 \pm 1.6$
Matter	$r_s(z_d)$	(6)	$153.3 \pm 2.0$ Mpc

- For flat LCDM, but very similar results for  $w \neq -1$  and curvature  $\neq 0$ !

# The Latest BAO Measurements



- 2dFGRS and SDSS main samples at  $z=0.2$
- SDSS LRG samples at  $z=0.35$
- These measurements constrain the ratio,  **$D_A(z)/d_s(z_{\text{BAO}})$** .

# Not Just $D_A(z)$ ...

- A really nice thing about BAO at a given redshift is that it can be used to measure not only  $D_A(z)$ , but also the expansion rate,  $H(z)$ , directly, at **that** redshift.

- BAO perpendicular to l.o.s

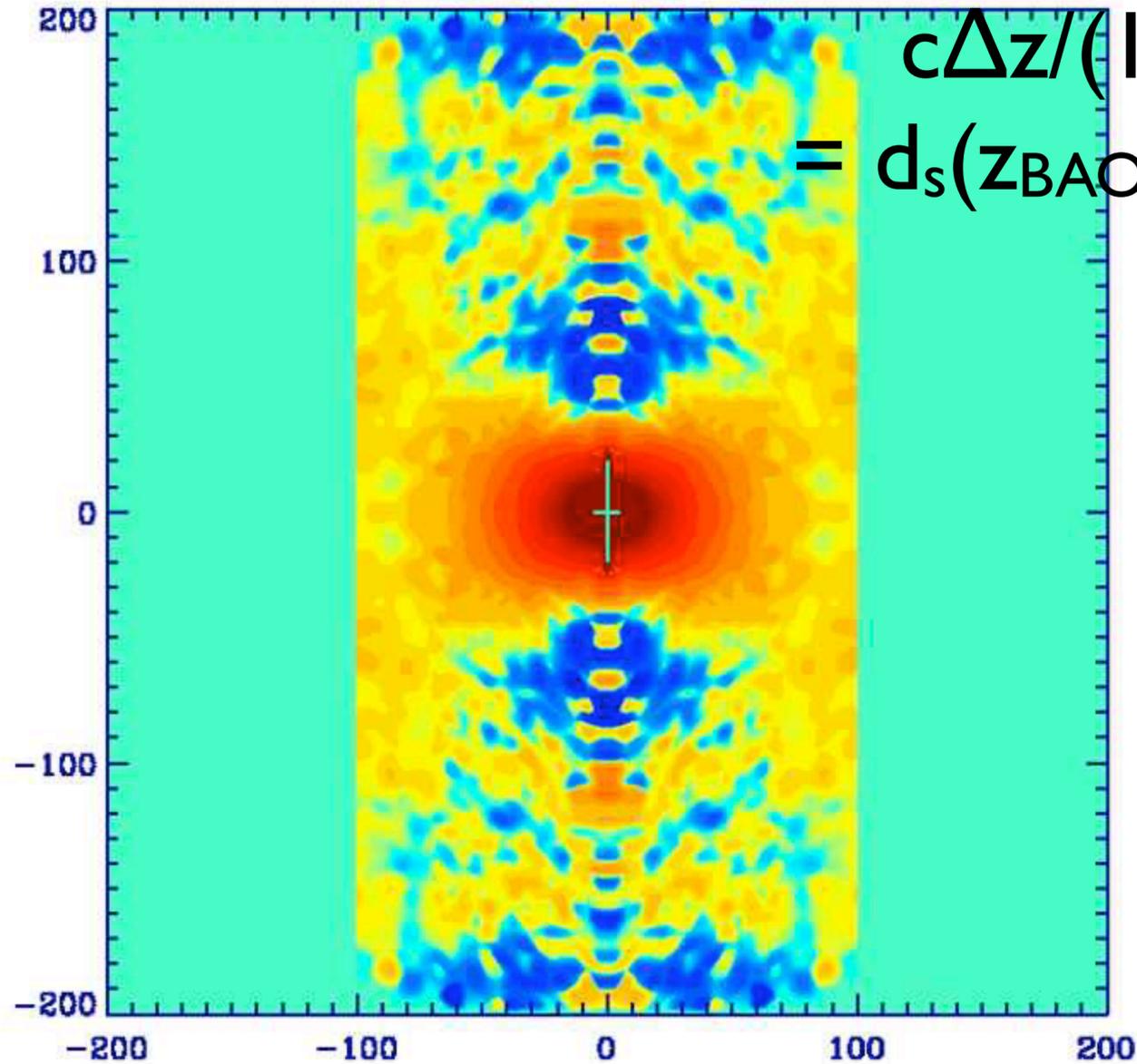
$$\Rightarrow D_A(z) = d_s(z_{\text{BAO}})/\theta$$

- BAO parallel to l.o.s

$$\Rightarrow \mathbf{H(z) = c\Delta z / [(1+z)d_s(z_{\text{BAO}})]}$$

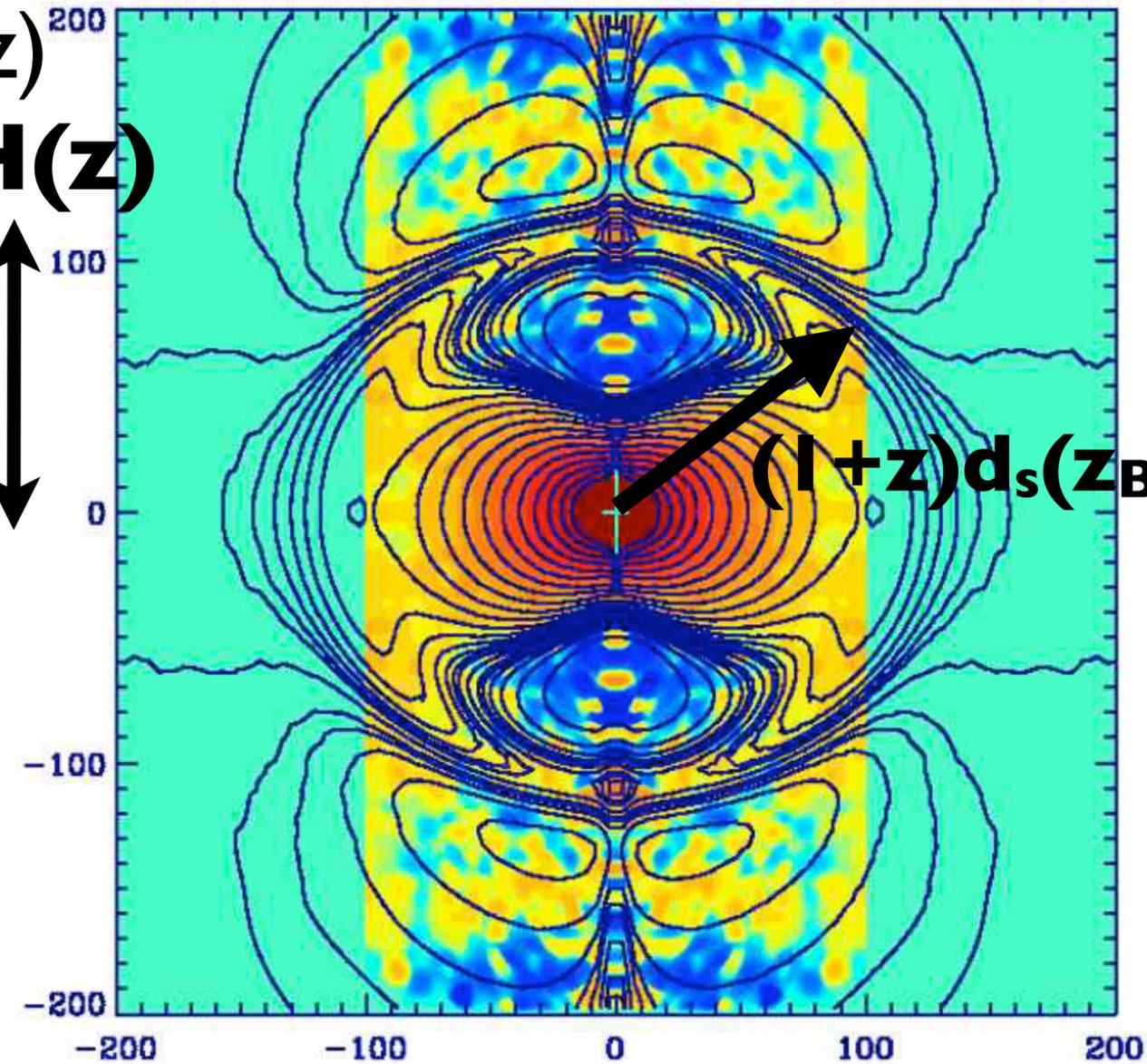
# Transverse= $D_A(z)$ ; Radial= $H(z)$

*SDSS Data*  
DR6



$$\frac{c\Delta z}{(1+z)} = d_s(z_{\text{BAO}}) \mathbf{H}(\mathbf{z})$$

*Linear Theory*  
DR6 + best model



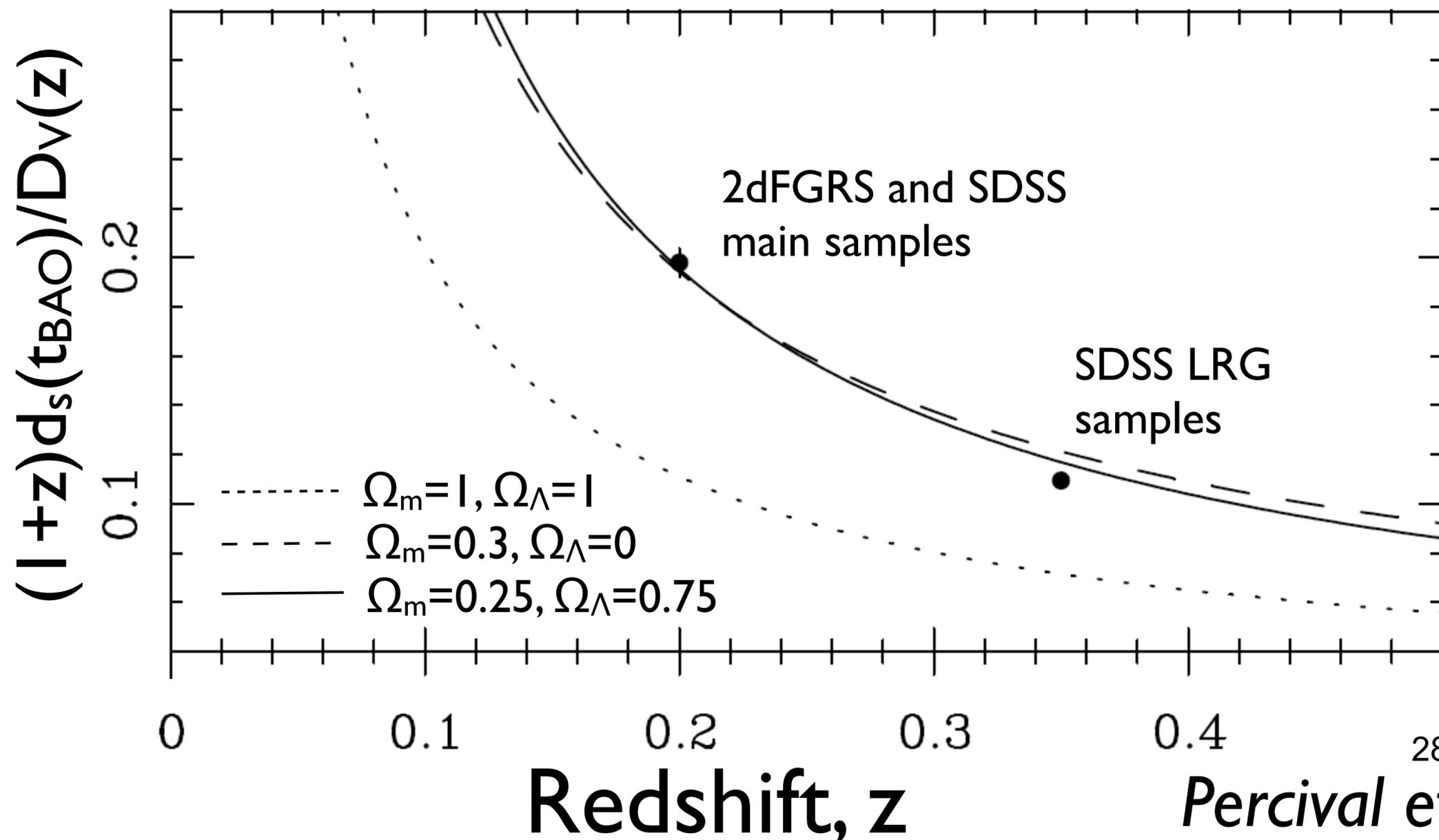
$$(1+z)d_s(z_{\text{BAO}})$$

Two-point correlation function measured from the SDSS Luminous Red Galaxies (Gaztanaga, Cabre & Hui 2008)

$$\theta = d_s(z_{\text{BAO}}) / \mathbf{D}_A(\mathbf{z})$$

$$D_V(z) = \left\{ (1+z)^2 D_A^2(z) [cz/H(z)] \right\}^{1/3}$$

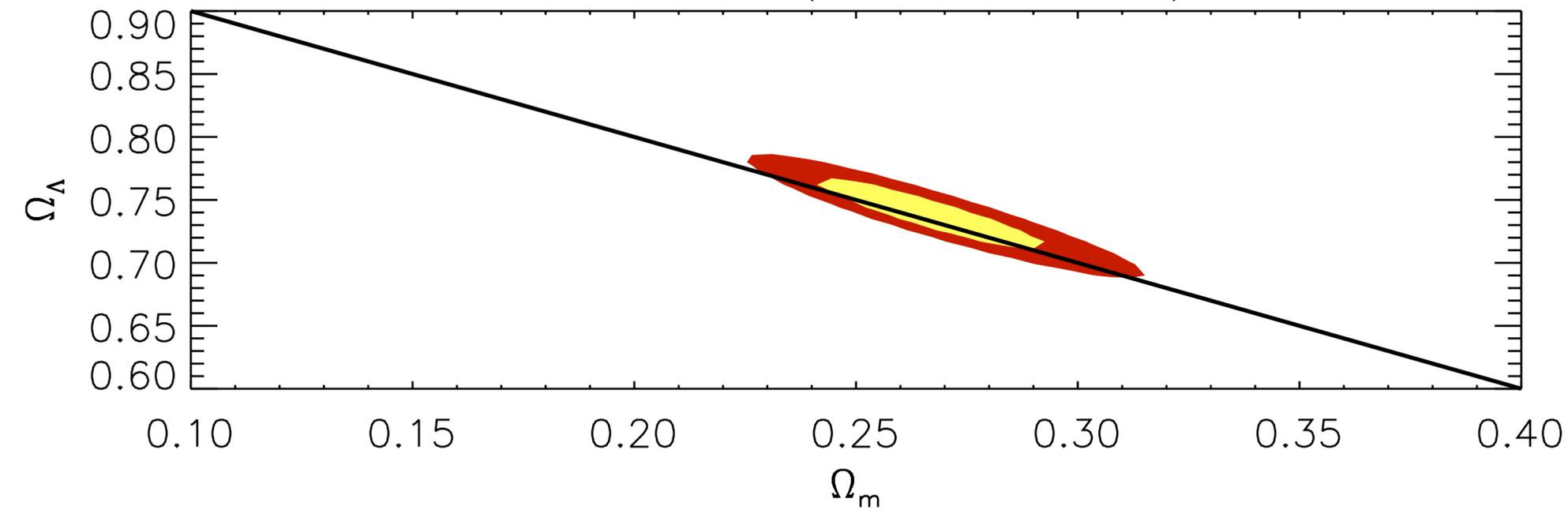
Since the current data are not good enough to constrain  $D_A(z)$  and  $H(z)$  separately, a combination distance,  $D_V(z)$ , has been constrained.



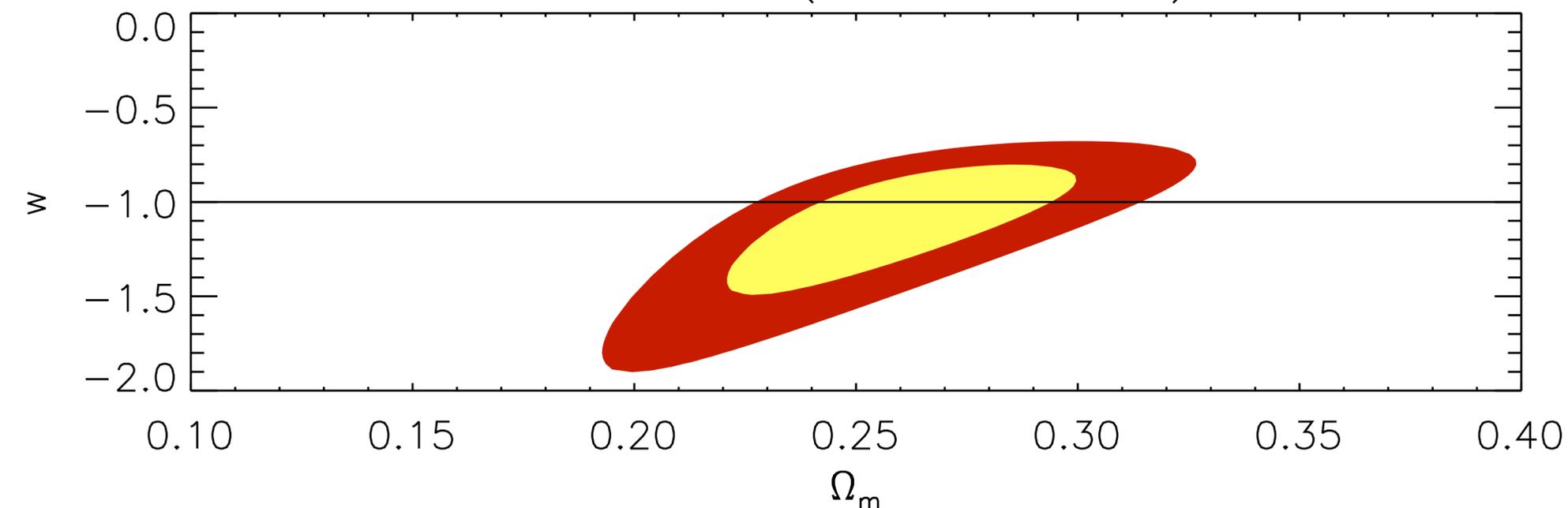
*Percival et al. (2007)*

# CMB + BAO $\Rightarrow$ Curvature

WMAP+BAO(Percival et al.)



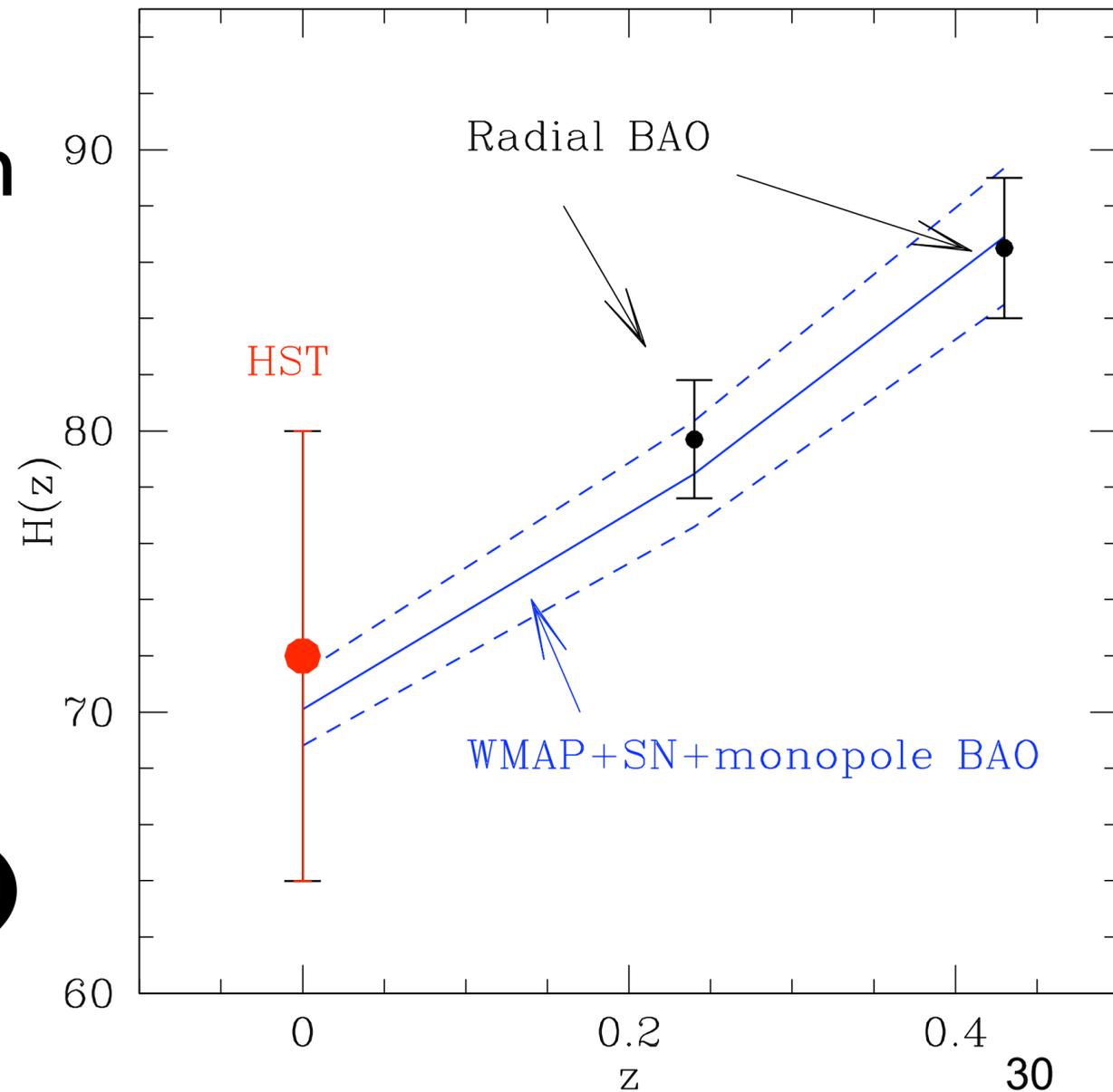
WMAP+BAO(Percival et al.)



- Both CMB and BAO are **absolute** distance indicators.
- Type Ia supernovae only measure relative distances.
- CMB+BAO is the winner for measuring spatial curvature.

# $H(z)$ also determined recently!

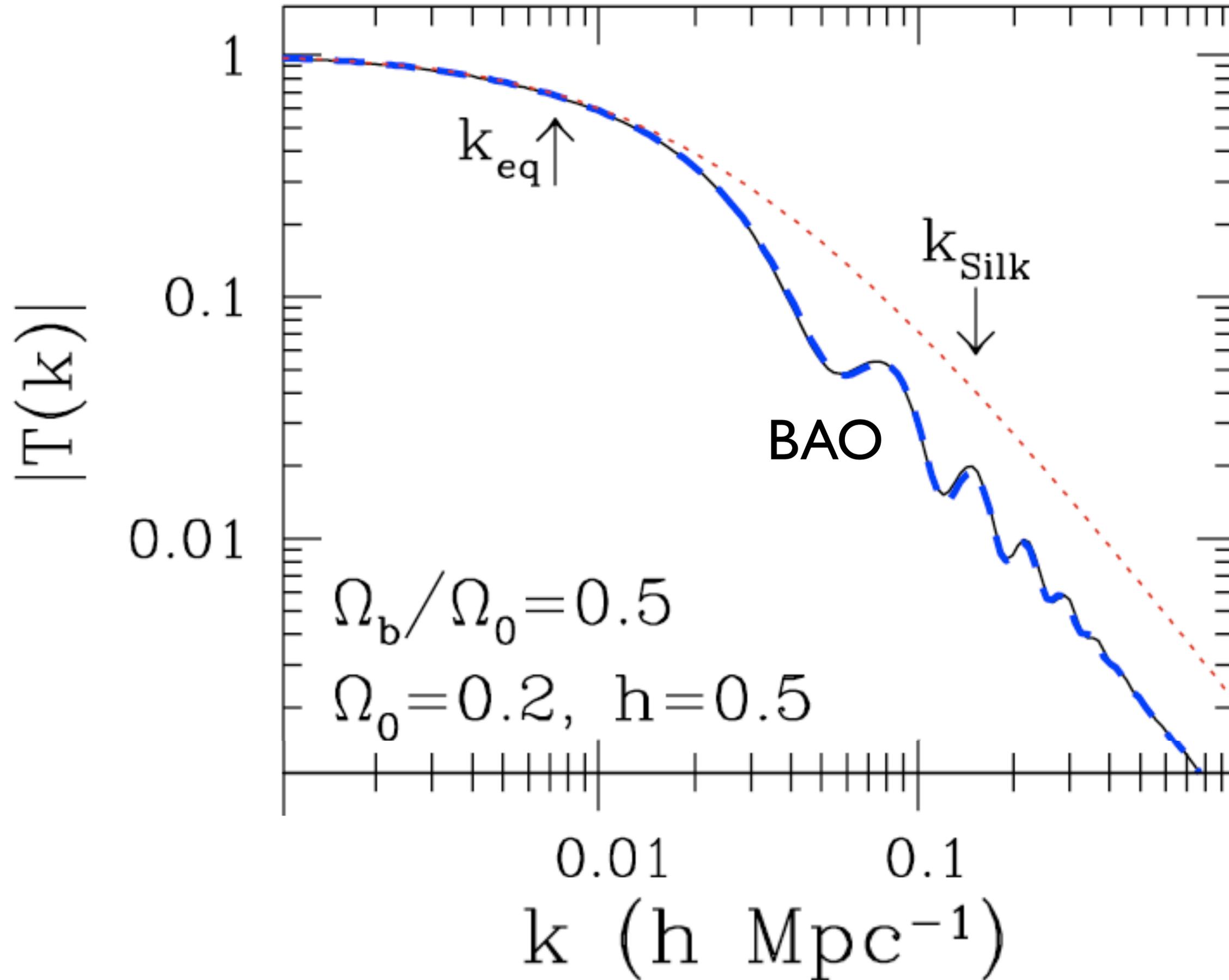
- SDSS DR6 data are now good enough to constrain  $H(z)$  from the 2-dimension correlation function *without spherical averaging*.
- Made possible by WMAP's measurement of  $r_s(z_{\text{BAO}}) = (1+z_{\text{BAO}})d_s(z_{\text{BAO}}) = 153.3 \pm 2.0 \text{ Mpc}$  (comoving)



*Gaztanaga, Cabre & Hui (2008)*

# Beyond BAO

- BAOs capture only a **fraction** of the information contained in the galaxy power spectrum!
- BAOs use the sound horizon size at  $z \sim 1020$  as the standard ruler.
- However, there are other standard rulers:
  - Horizon size at the matter-radiation equality epoch ( $z \sim 3200$ )
  - Silk damping scale



# ...and, these are all well known

- Cosmologists have been measuring  $k_{eq}$  over the last three decades.
- This was usually called the “Shape Parameter,” denoted as  $\Gamma$ .
- $\Gamma$  is proportional to  $k_{eq}/h$ , and:
  - The effect of the Silk damping is contained in the constant of proportionality.
  - Easier to measure than BAOs: the signal is much stronger.

# WMAP & Standard Ruler

- **With WMAP 5-year data only**, the scales of the standard rulers have been determined accurately.

- Even when  $w \neq -1$ ,  $\Omega_k \neq 0$ ,

- $d_s(z_{\text{BAO}}) = 153.4^{+1.9}_{-2.0} \text{ Mpc}$  ( $z_{\text{BAO}} = 1019.8 \pm 1.5$ )

- $k_{\text{eq}} = (0.975^{+0.044}_{-0.045}) \times 10^{-2} \text{ Mpc}^{-1}$  ( $z_{\text{eq}} = 3198^{+145}_{-146}$ )

- $k_{\text{silky}} = (8.83 \pm 0.20) \times 10^{-2} \text{ Mpc}^{-1}$

With Planck, they will be determined to higher precision.

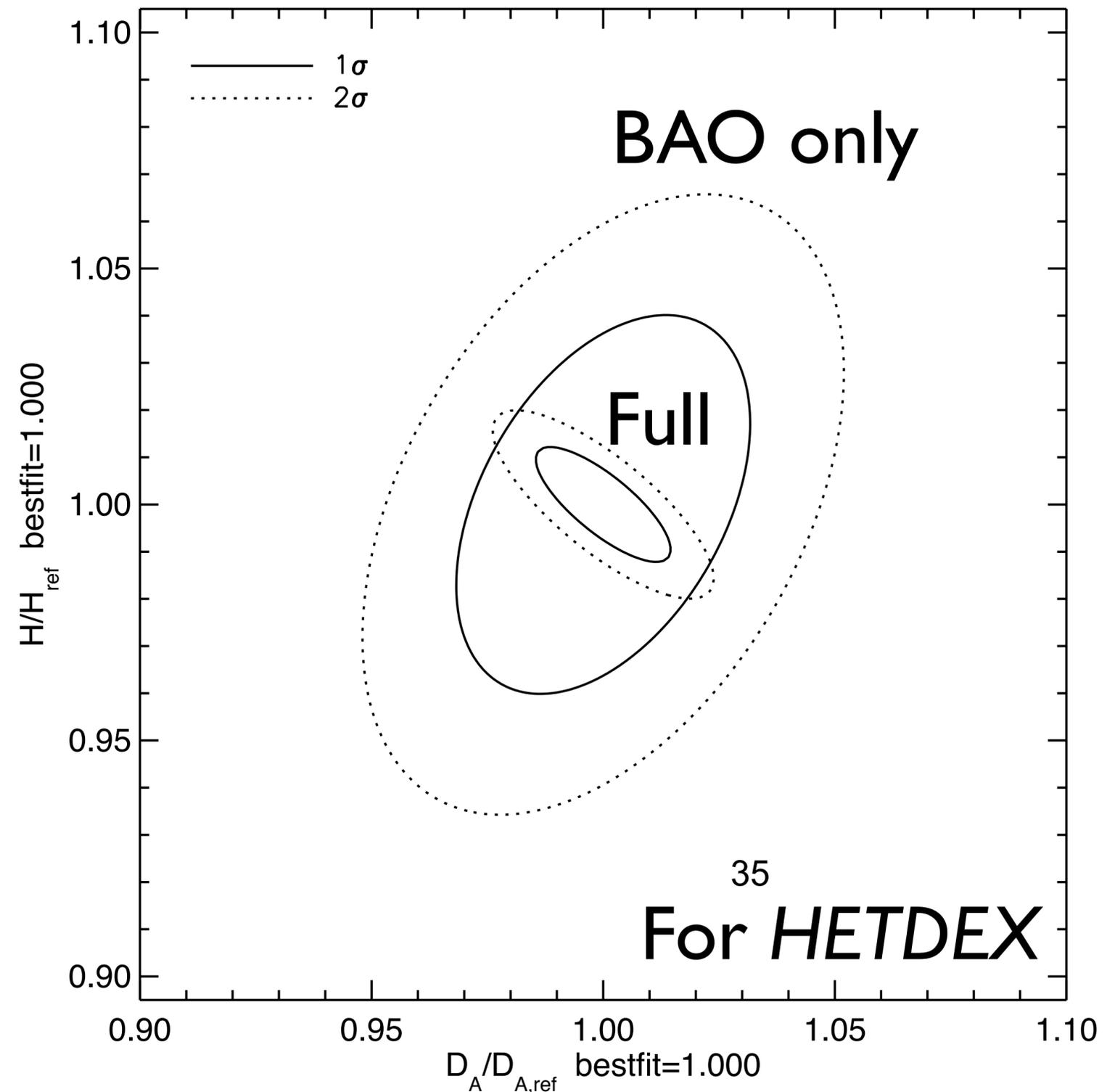
1.3%

4.6%

2.3%

# BAO vs Full Modeling

- Full modeling improves upon the determinations of  $D_A$  &  $H$  by more than a factor of two.
- On the  $D_A$ - $H$  plane, the size of the ellipse shrinks by more than a factor of four.



# Why Not “GPS,” Instead of “BAO”?

- JDEM says, “SN, WL, or BAO at minimum.”
- It does not make sense to single out “BAO”: the observable is the **galaxy power spectrum (GPS)**.
- To get BAO, we need to measure the galaxy power spectrum anyway.
- If we measure the galaxy power spectrum, why just focus on BAO? There is much more information!
- So, it would be better for JDEM to say, perhaps, “use SN, WL, or GPS at minimum”?

# WMAP Amplitude Prior

- WMAP measures the amplitude of curvature perturbations at  $z \sim 1090$ . Let's call that  $R_k$ . The relation to the density fluctuation is

$$\delta_{m,\mathbf{k}}(z) = \frac{2k^3}{5H_0^2\Omega_m} \mathcal{R}_k T(k) D(k, z)$$

- Variance of  $R_k$  has been constrained as:

AMPLITUDE OF CURVATURE PERTURBATIONS,  $\mathcal{R}$ ,  
MEASURED BY WMAP AT  $k_{WMAP} = 0.02 \text{ Mpc}^{-1}$

Model	$10^9 \times \Delta_{\mathcal{R}}^2(k_{WMAP})$
$\Omega_k = 0$ and $w = -1$	$2.211 \pm 0.083$
$\Omega_k \neq 0$ and $w = -1$	$2.212 \pm 0.084$
$\Omega_k = 0$ and $w \neq -1$	$2.208 \pm 0.087$
$\Omega_k \neq 0$ and $w \neq -1$	$2.210 \pm 0.084$
$\Omega_k = 0$ , $w = -1$ and $m_\nu > 0$	$2.212 \pm 0.083$
$\Omega_k = 0$ , $w \neq -1$ and $m_\nu > 0$	$2.218 \pm 0.085$
WMAP Normalization Prior	$2.21 \pm 0.09$

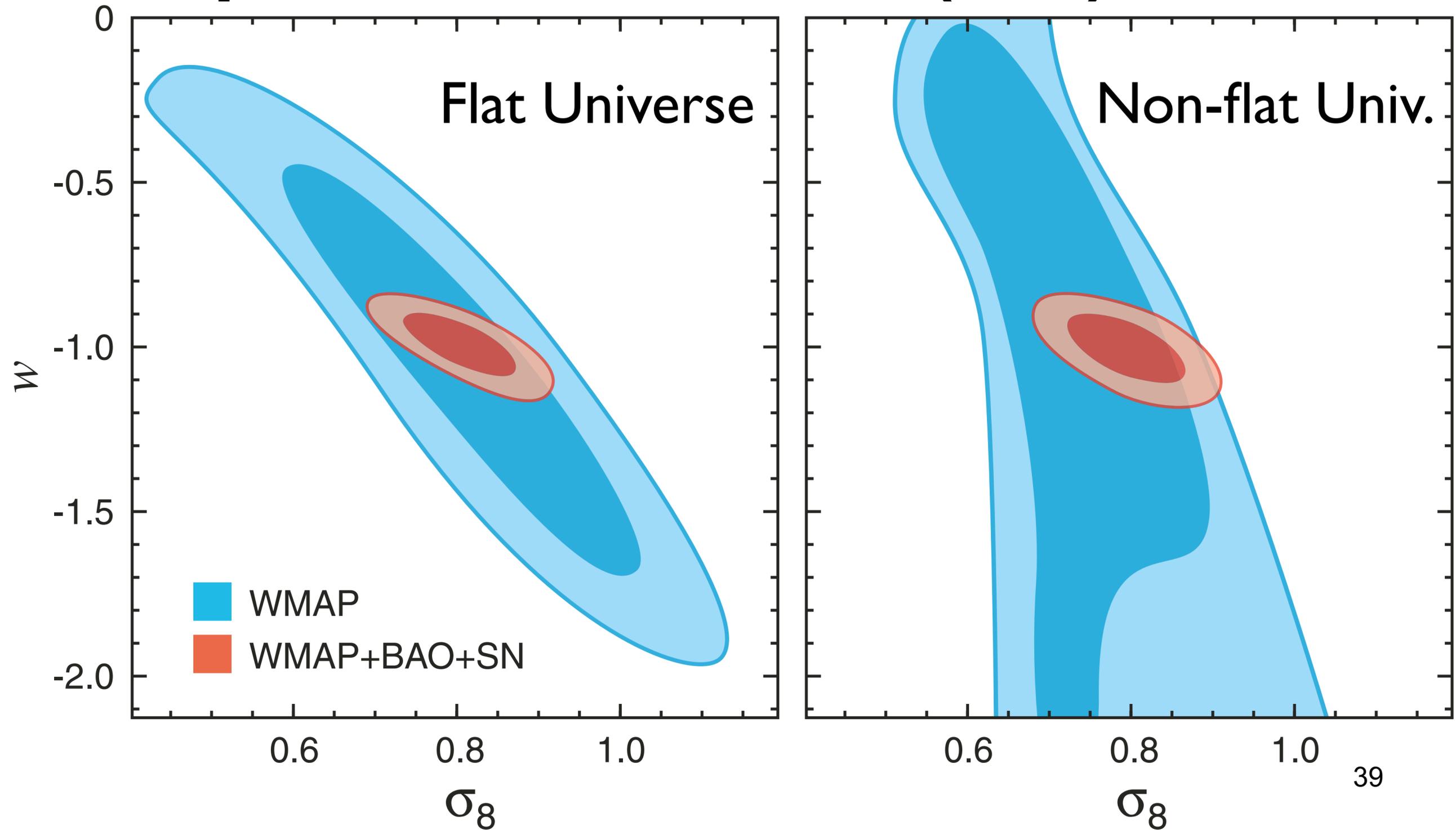
# Then Solve This Diff. Equation...

Ignoring the mass of neutrinos and modifications to gravity, one can obtain the growth rate by solving the following differential equation (Wang & Steinhardt 1998; Linder & Jenkins 2003):  $\mathbf{g}(\mathbf{z})=(\mathbf{I}+\mathbf{z})\mathbf{D}(\mathbf{z})$

$$\frac{d^2 g}{d \ln a^2} + \left[ \frac{5}{2} + \frac{1}{2} (\Omega_k(a) - 3w_{\text{eff}}(a)\Omega_{de}(a)) \right] \frac{dg}{d \ln a} + \left[ 2\Omega_k(a) + \frac{3}{2} (1 - w_{\text{eff}}(a))\Omega_{de}(a) \right] g(a) = 0, \quad (76)$$

- If you need a code for doing this, search for **“Cosmology Routine Library”** on Google 38

# Degeneracy Between Amplitude at $z=0$ ( $\sigma_8$ ) and $w$



# Summary

- WMAP helps constrain the nature of DE by providing:
  - Angular diameter distance to  $z^* \sim 1090$ ,
  - Amplitude of fluctuations at  $z^* \sim 1090$ , and
  - $\partial\Phi/\partial t$  at  $z < 1$  via the Integrated Sachs-Wolfe effect.
- WMAP also measures the sound horizon size for baryons,  $d_{\text{BAO}}$ , which is used by BAO experiments to constrain  $D_A(z)$  and  $H(z)$ .
- Not just BAO! WMAP also provides the other standard rulers,  $k_{\text{eq}}$  and  $k_{\text{silk}}$ , with which the accuracy of  $D_A(z)$  and  $H(z)$  from galaxy surveys can be improved greatly.