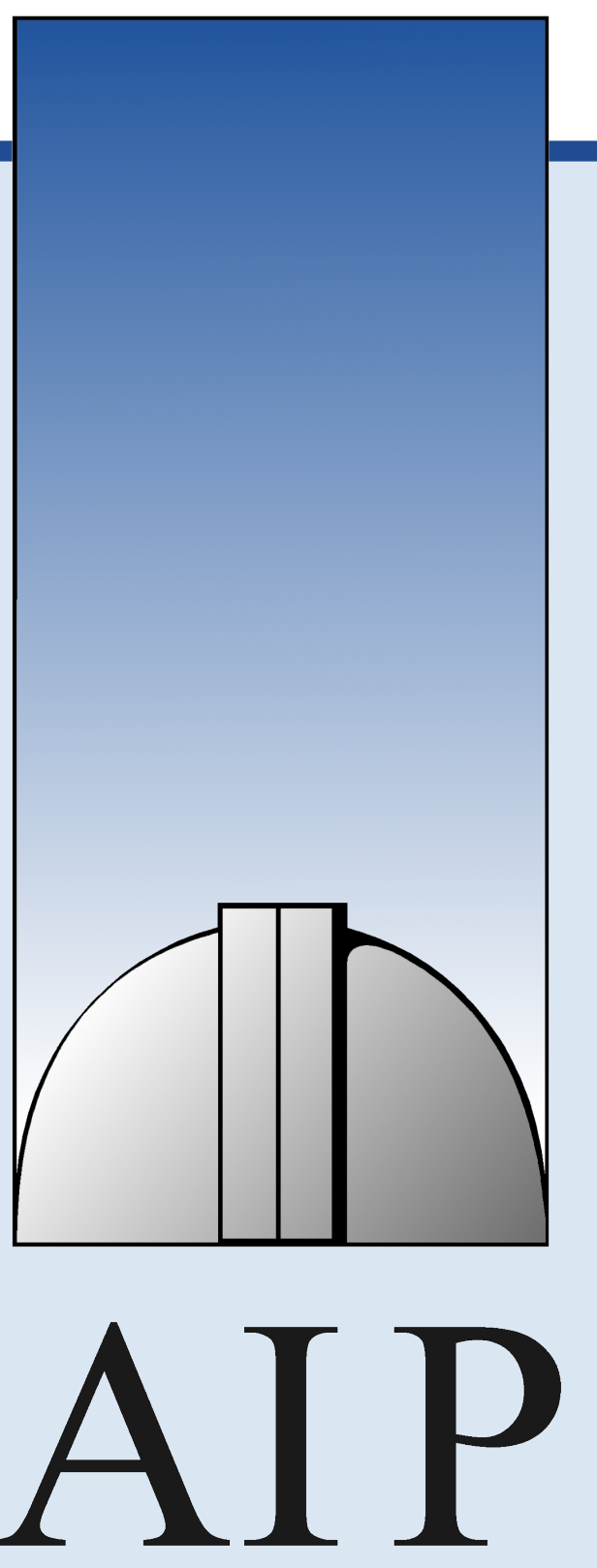


# How to constrain Dark Energy with Baryon Acoustic Oscillations



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## Baryon Acoustic Oscillations

The baryon acoustic oscillations (BAO), imprinted on the matter power spectrum  $P(k)$ , have the same origin as the acoustic anisotropies of the CMB. Therefore their physical scale is known to sufficient accuracy and they can be used as a **standard ruler** to measure the expansion history of the Universe, which depends on the dark energy **equation of state parameter**  $w$ . Using  $\Lambda$ CDM as reference cosmology the BAO are scaled by  $H/H_\Lambda$  and  $D_A^\Lambda/D_A$  along and transverse to the line of sight, respectively. On the right we show the extracted BAO for  $w = -0.8, -1.0, -1.2$ .

With N-Body simulations we are able to analyze nonlinear effects and can predict how well upcoming galaxy redshift surveys, like the Hobby-Eberly Telescope Dark Energy Experiment, **HETDEX**, can constrain dark energy.

Nonlinear evolution diminishes the amplitude of the BAO (see right). Using **reconstruction** techniques for the density field we can almost completely reestablish the amplitude and by this improve the measurement of the apparent scale of the BAO, transverse (angular diameter distance  $D_A$ ) and parallel (Hubble parameter  $H$ ) to the line of sight. For surveys at lower redshifts, like **BOSS**, this effect is large and reconstruction becomes very important.

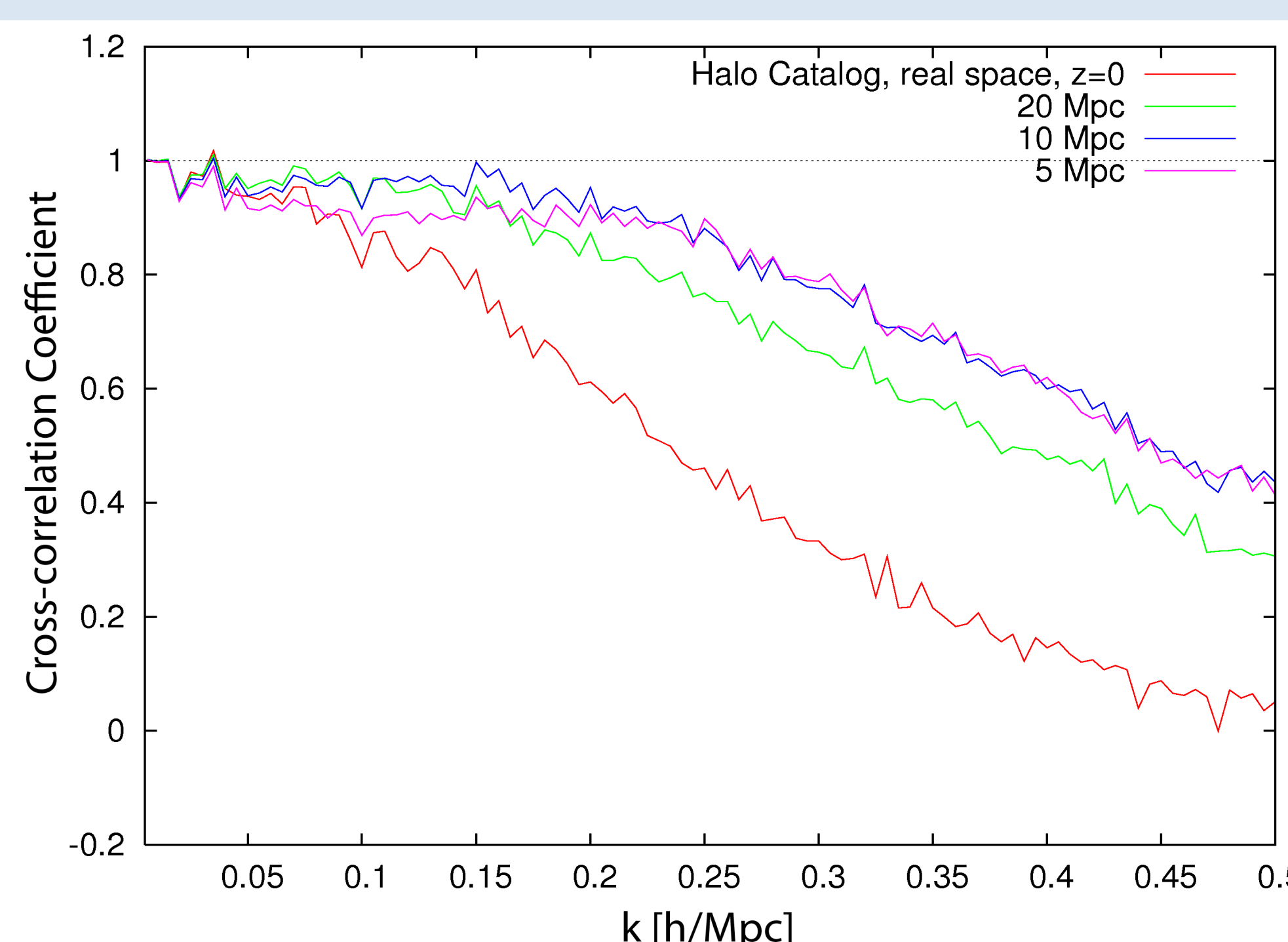
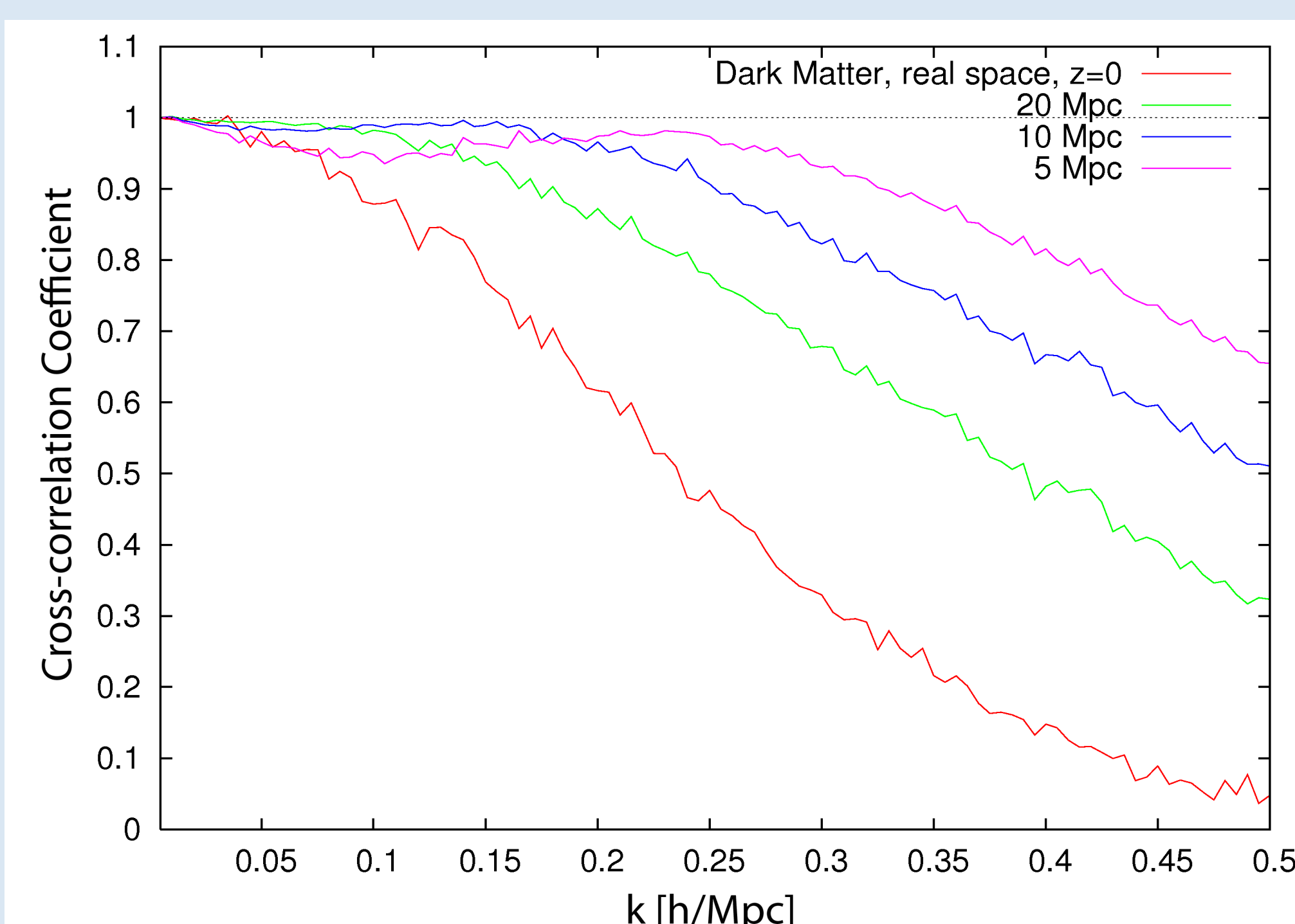
To get constraints on the dark energy parameter  $w$  one has to take other cosmological parameters, like the Hubble constant  $H_0$ , matter fraction  $\omega_m$  and baryon fraction  $\omega_b$ , into account. In particular, the uncertainty in the Hubble constant degrades the constraints on  $w$ .

## Reconstruction

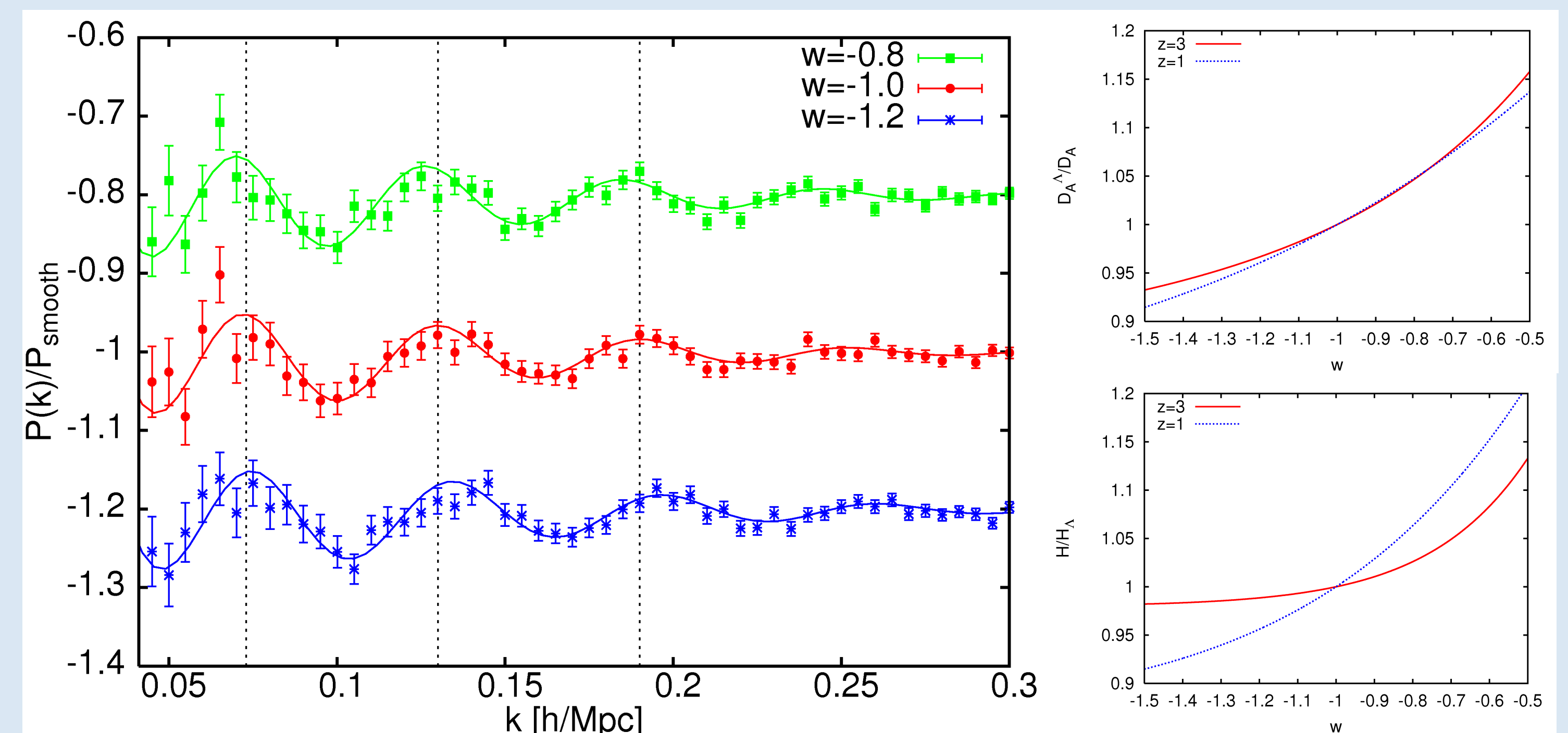
We implemented two different reconstruction techniques. **PIZA** (Path Interchange Zeldovich Approximation) is particle based and derives the displacement field by minimizing iteratively the sum of the displacements squared.

The **linear reconstruction** scheme uses the density field to determine the displacement field. It applies linear theory to the observed density field after smoothing it on a sufficiently large scale.

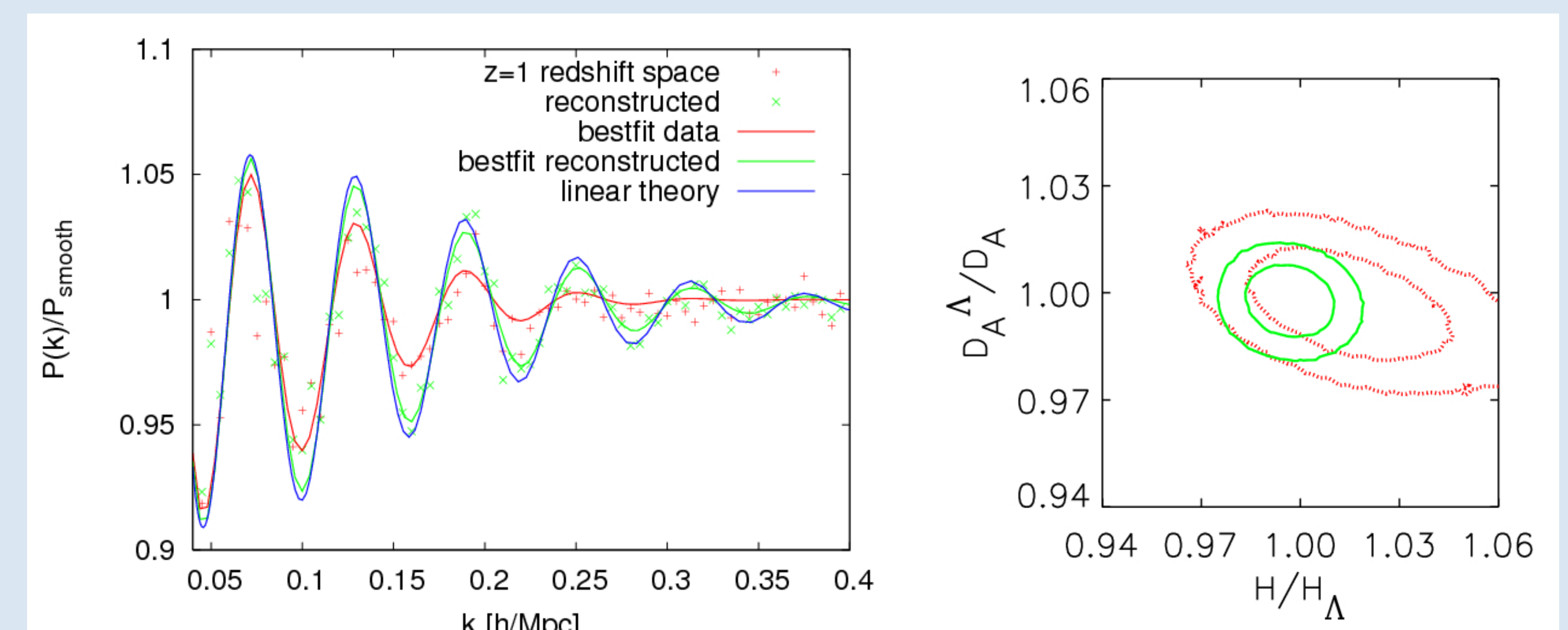
A measure how well the initial density field is reconstructed is the cross-correlation coefficient of the initial and the reconstructed density field  $\langle \delta_{\text{in}}(k) \delta_{\text{rec}}^*(k) / (|\delta_{\text{in}}(k)|^2 D) \rangle$ , with  $\delta(k)$  the Fourier mode of the density and  $D$  the growth function.



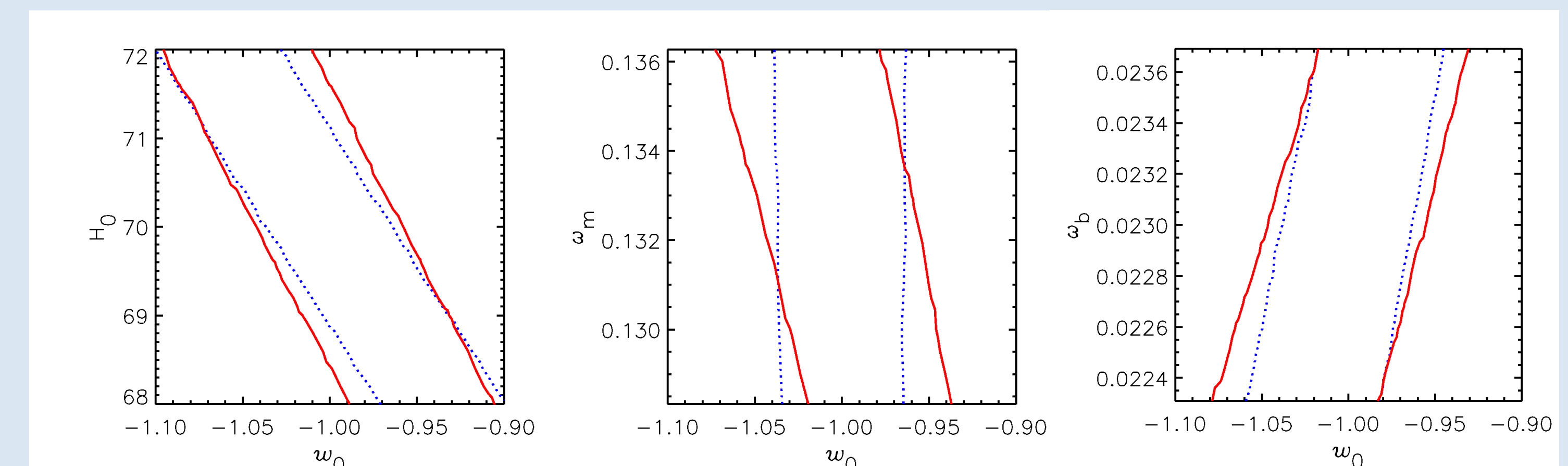
Cross-correlation coefficient of the initial and reconstructed density field at redshift  $z=0$  for different smoothing lengths. The reconstruction was performed from the dark matter (left) and halo (right) distribution by applying the linear reconstruction scheme.



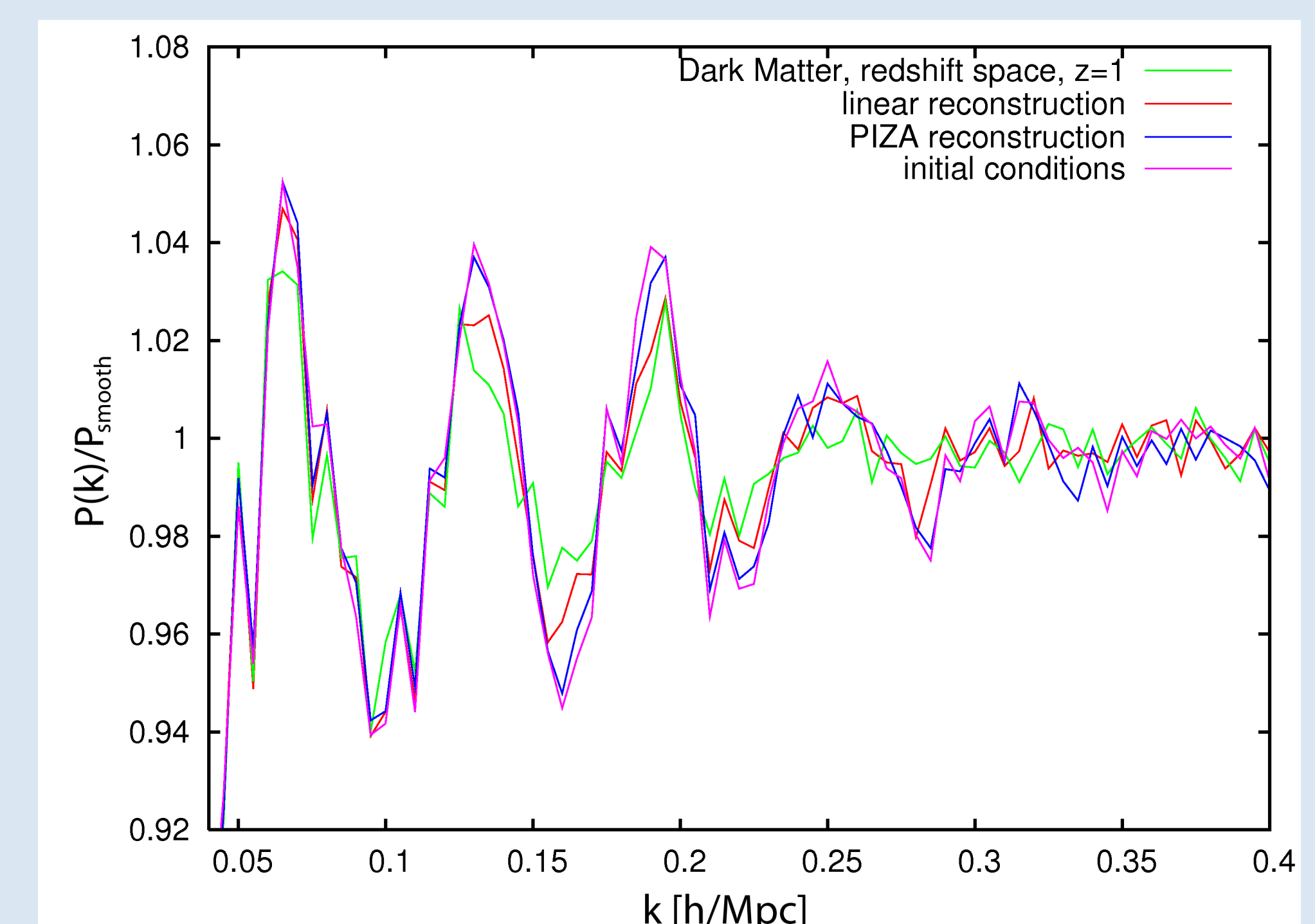
Left: The extracted BAO for  $w = -0.8, -1.0, -1.2$  at  $z=3$  (shifted by  $-0.8, -1.0, -1.2$  along the y-axis, respectively). Right: Scaling factors as a function of  $w$ .



Left: The extracted BAO imprinted in the dark matter distribution. Right: Contour lines (68% and 95% c.l.) for the scaling factors derived from the BAO **before** and **after** the reconstruction of the density field.



Degeneracies (95% c.l.) of the cosmological parameters  $H_0$ ,  $\omega_m$ , and  $\omega_b$  with respect to a constant  $w=w_0$  at redshift  $z=1$  (dotted) and  $z=3$  (solid) derived from the dark matter distribution in a 1.5 Gpc/h computational box.



Reconstruction of the BAO in redshift space at  $z=0$  using different reconstruction techniques.