

Dark energy without dark energy

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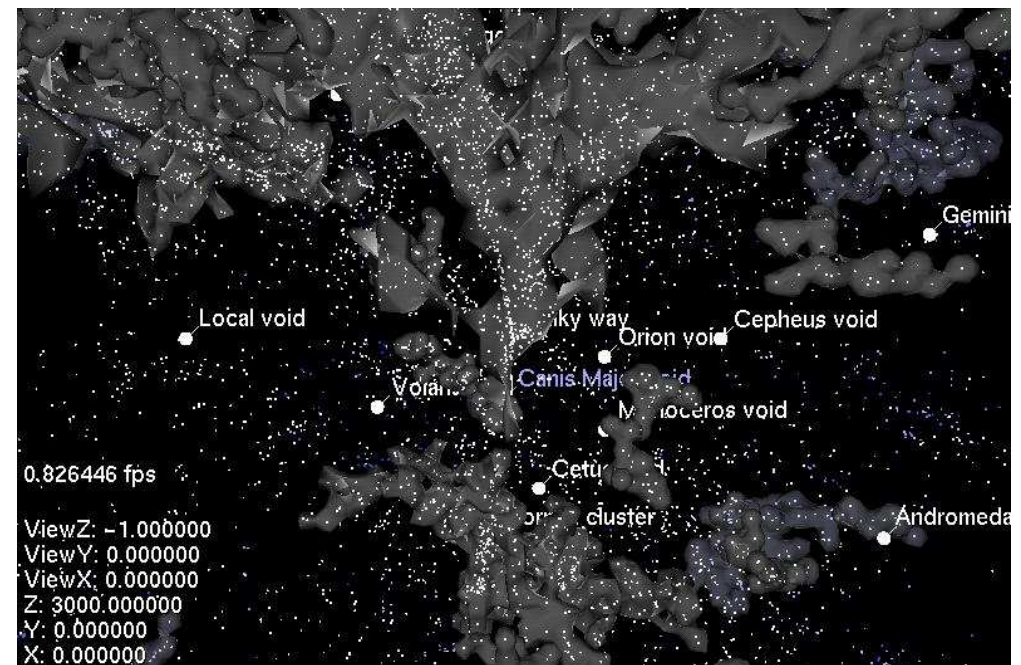
Phys. Rev. D 78 (2008) in press

[arxiv:0809.1183];

new results, to appear

B.M. Leith, S.C.C. Ng and DLW:

ApJ 672 (2008) L91 [arXiv:0709.2535]



What is “dark energy”?

- Usual explanation: a homogeneous isotropic form of “stuff” which violates the strong energy condition. (Locally pressure $P = w\rho c^2$, $w < -\frac{1}{3}$; e.g., for cosmological constant, Λ , $w = -1$.)
- New explanation: in ordinary general relativity, a manifestation of global *variations* of those aspects of gravitational energy which by virtue of the equivalence principle cannot be localised – the *cosmological quasilocal gravitational energy* associated with *dynamical gradients* in spatial curvature generated by a universe as inhomogeneous as the one we observe. [Call this *dark energy* if you like. It involves *energy*, and “nothing” is dark.]

From smooth to lumpy

- Universe was very smooth at time of last scattering; fluctuations in the fluid were tiny ($\delta\rho/\rho \sim 10^{-5}$ in photons and baryons; $\sim 10^{-3}$ in non-baryonic dark matter).
- FLRW approximation very good early on.
- Universe is very lumpy or inhomogeneous today.
- Recent surveys estimate that 40–50% of the volume of the universe is contained in voids of diameter $30h^{-1}$ Mpc. [Hubble constant $H_0 = 100h \text{ km sec}^{-1} \text{ Mpc}^{-1}$] (Hoyle & Vogeley, ApJ 566 (2002) 641; 607 (2004) 751)
- Add some larger voids, and many smaller minivoids, and the universe is *void-dominated* at present epoch.
- Clusters of galaxies are strung in filaments and bubbles around these voids.

Buchert's dust equations (2000)

For irrotational dust cosmologies, characterised by an energy density, $\rho(t, \mathbf{x})$, expansion, $\theta(t, \mathbf{x})$, and shear, $\sigma(t, \mathbf{x})$, on a compact domain, \mathcal{D} , of a suitably defined spatial hypersurface of constant average time, t , and spatial 3-metric, average cosmic evolution in Buchert's scheme is described by the exact equations

$$3\frac{\dot{\bar{a}}^2}{\bar{a}^2} = 8\pi G\langle\rho\rangle - \frac{1}{2}\langle\mathcal{R}\rangle - \frac{1}{2}\mathcal{Q}$$

$$3\frac{\ddot{\bar{a}}}{\bar{a}} = -4\pi G\langle\rho\rangle + \mathcal{Q}$$

$$\partial_t\langle\rho\rangle + 3\frac{\dot{\bar{a}}}{\bar{a}}\langle\rho\rangle = 0$$

$$\mathcal{Q} \equiv \frac{2}{3}(\langle\theta^2\rangle - \langle\theta\rangle^2) - 2\langle\sigma^2\rangle$$

Back-reaction

Angle brackets denote the spatial volume average, e.g.,

$$\langle \mathcal{R} \rangle \equiv \left(\int_{\mathcal{D}} d^3x \sqrt{\det {}^3g} \mathcal{R}(t, \mathbf{x}) \right) / \mathcal{V}(t)$$

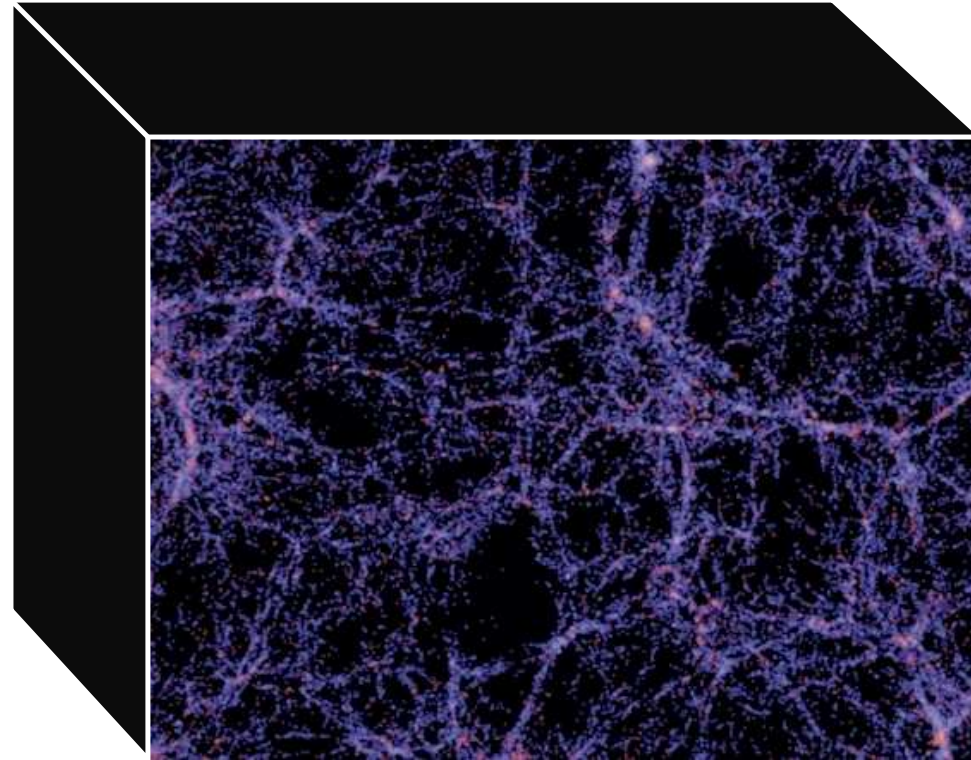
$$\langle \theta \rangle = 3 \frac{\dot{\bar{a}}}{\bar{a}}$$

Generally for any scalar Ψ ,

$$\frac{d}{dt} \langle \Psi \rangle - \left\langle \frac{d\Psi}{dt} \right\rangle = \langle \Psi \theta \rangle - \langle \theta \rangle \langle \Psi \rangle$$

- The extent to which the back–reaction, \mathcal{Q} , can lead to apparent cosmic acceleration or not has been the subject of much debate.

Within a statistically average cell



- Need to consider relative position of observers over scales of tens of Mpc over which $\delta\rho/\rho \sim -1$.
- GR is a local theory: gradients in spatial curvature and gravitational energy can lead to calibration differences between our rods and clocks and volume average ones

The Copernican principle

- Retain Copernican Principle - we are at an average position *for observers in a galaxy*
- Observers in bound systems are not at a volume average position in freely expanding space
- By Copernican principle other average observers should see an isotropic CMB
- BUT *nothing in theory, principle nor observation demands that such observers measure the same mean CMB temperature nor the same angular scales in the CMB anisotropies*
- Average mass environment (galaxy) can differ significantly from volume–average environment (void). Galaxies formed from perturbations that were greater than critical density \Rightarrow natural separation of scales.

Dilemma of gravitational energy...

- In GR spacetime carries *energy* & *angular momentum*

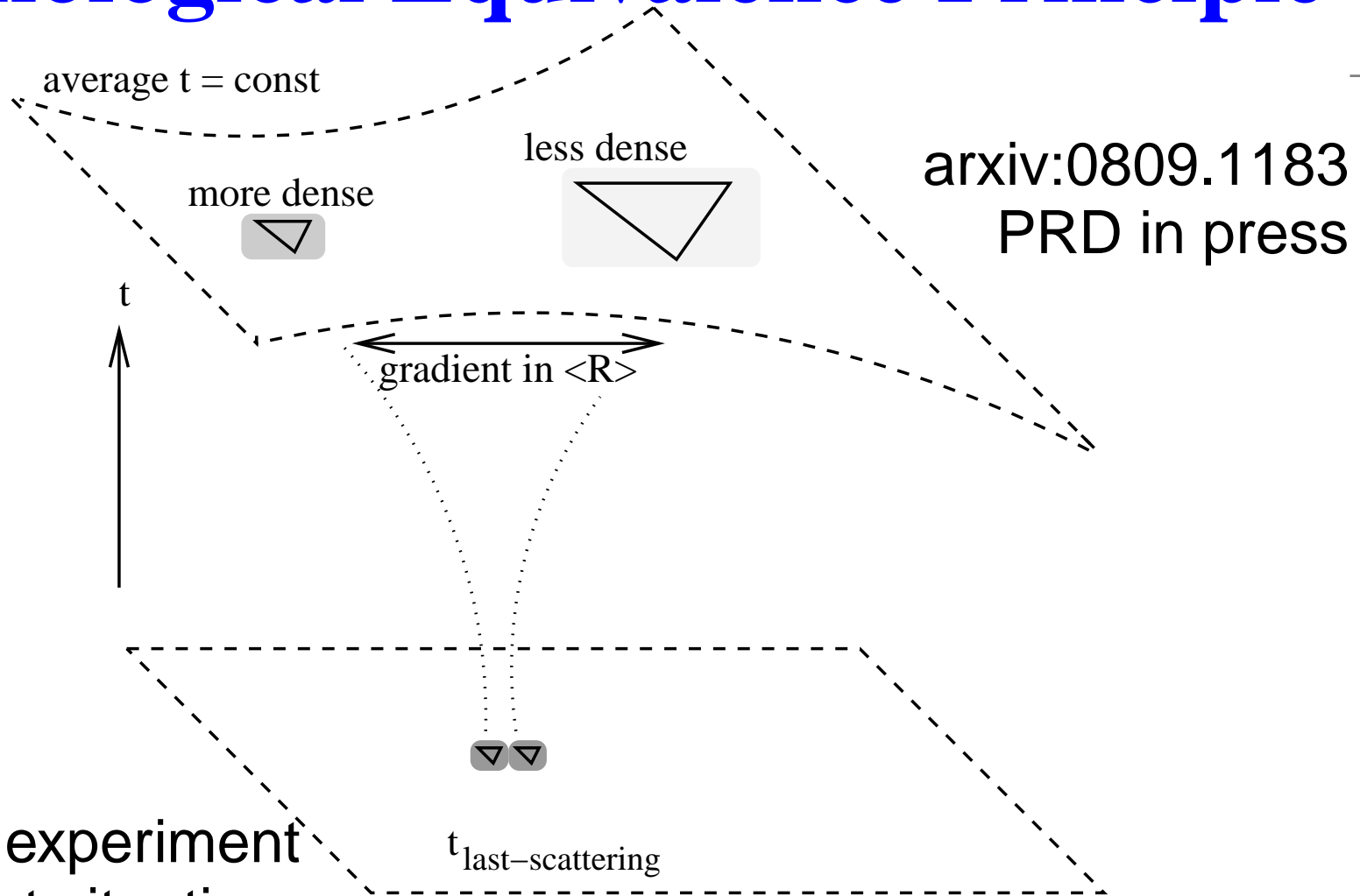
$$G_{\mu\nu} = \frac{8\pi G}{c^4} T_{\mu\nu}$$

- On account of the strong equivalence principle, $T_{\mu\nu}$ contains localizable energy–momentum only
- Kinetic energy and energy associated with spatial curvature are in $G_{\mu\nu}$: variations are “quasilocal”!
- Newtonian version, $T - U = -V$, of Friedmann equation

$$\frac{\dot{a}^2}{a^2} + \frac{kc^2}{a^2} = \frac{8\pi G\rho}{3}$$

where $T = \frac{1}{2}m\dot{a}^2x^2$, $U = -\frac{1}{2}kmc^2x^2$, $V = -\frac{4}{3}\pi G\rho a^2x^2m$;
 $\mathbf{r} = a(t)\mathbf{x}$.

Cosmological Equivalence Principle



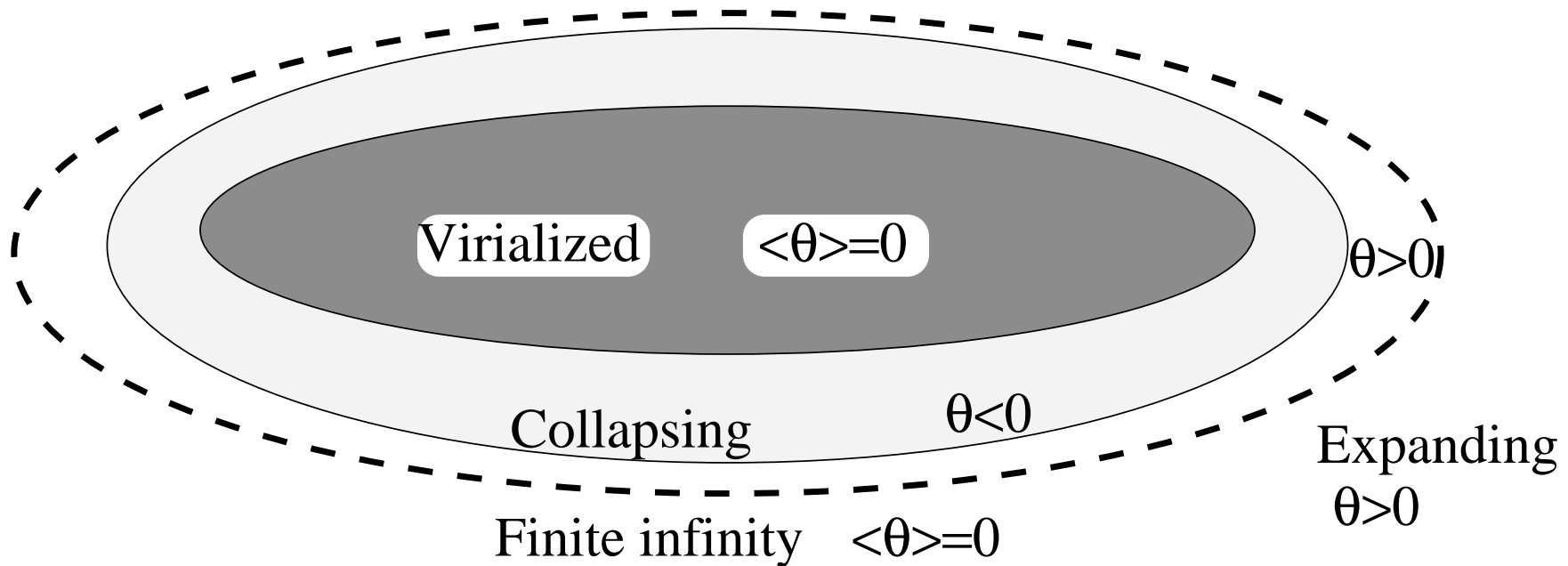
Thought experiment
equivalent situations:

- SR: observers volume decelerate at different rates
- GR: regions of different density have different volume deceleration (for same initial conditions)

Bound and unbound systems...

- Isotropic observers “at rest” within expanding space in voids may have clocks ticking at a rate $d\tau_v = \gamma(\tau_w, \mathbf{x})d\tau_w$ with respect to static observers in bound systems.
Volume average: $dt = \bar{\gamma}_w d\tau_w$, $\bar{\gamma}_w(\tau_w) = \langle -\xi^\mu n_\mu \rangle_{\mathcal{H}}$
- We are not restricted to $\gamma = 1 + \epsilon$, $\epsilon \ll 1$, as expected for typical variations of binding energy.
- Observable universe is assumed unbound.
- I find $\bar{\gamma} \simeq 1.38$ at present epoch from relative regional deceleration $\sim 10^{-10} \text{ms}^{-2}$ integrated over age of universe
- Where is infinity? In 1984 George Ellis suggested a notion of *finite infinity*: a region within which isolated systems, such as stars or galaxies, or galaxy clusters are approximately independent dynamical systems.

Finite infinity



- Define *finite infinity*, "*fi*" as boundary to minimal *connected* region within which *average expansion* vanishes $\langle \theta \rangle = 0$ or average curvature vanishes $\langle R \rangle = 0$.
- Shape of *fi* boundary irrelevant (minimal surface generally): could typically contain a galaxy cluster.

Cosmic rest frame

- Patch together CIFs for observers who see an isotropic CMB by taking surfaces of uniform volume expansion

$$\left\langle \frac{1}{\ell_r(\tau)} \frac{d\ell_r(\tau)}{d\tau} \right\rangle = \frac{1}{3} \langle \theta \rangle_1 = \frac{1}{3} \langle \theta \rangle_2 = \dots = \bar{H}(\tau)$$

- Average over regions in which (i) spatial curvature is zero or negative; (ii) space is expanding at the boundaries, at least marginally.
- Solves the Sandage–de Vaucouleurs paradox implicitly.
- Voids appear to expand faster; but their local clocks tick faster, locally measured expansion can still be uniform.
- Global average H_{av} on large scales with respect to *any one set of clocks* will differ from \bar{H}

Two/three scale model

$$\bar{a}^3 = f_{\text{wi}} a_{\text{w}}^3 + f_{\text{vi}} a_{\text{v}}^3$$

- Splits into void fraction with scale factor a_{v} and “wall” fraction with scalar factor a_{w} , combined average \bar{a} .
- Solve Buchert equations for volume averaged observer, with $f_{\text{v}}(t) = f_{\text{vi}} a_{\text{v}}^3 / \bar{a}^3$ (void volume fraction) and $k_{\text{v}} < 0$
- Bare cosmological parameters, $\bar{\Omega}_M + \bar{\Omega}_k + \bar{\Omega}_{\mathcal{Q}} = 1$, as inferred at volume average position, differ little from an open Friedmann model: $|\bar{\Omega}_{\mathcal{Q}}| \lesssim 0.04$
- But must account for cumulative growth in spatial curvature and gravitational energy gradients – *clock rate variance* – between galaxies and voids
- Perform light cone average. Conformally match radial null geodesics in average geometry and wall geometry with uniform local Hubble flow. Match volume factors.

Dressed cosmological parameters

- Conventional parameters for “wall observers” in galaxies: defined by assumption (no longer true) that others in entire observable universe have synchronous clocks and same local spatial curvature

$$\begin{aligned} ds_{\mathcal{F}_I}^2 &= -d\tau_w^2 + a_w^2(\tau_w) [d\eta_w^2 + \eta_w^2 d\Omega^2] \\ &= -d\tau_w^2 + \frac{\bar{a}^2}{\bar{\gamma}_w^2} [d\bar{\eta}^2 + r_w^2(\bar{\eta}, \tau_w) d\Omega^2] \end{aligned}$$

where $r_w \equiv \bar{\gamma}_w (1 - f_v)^{1/3} f_{wi}^{-1/3} \eta_w(\bar{\eta}, \tau_w)$, and volume-average conformal time $d\bar{\eta} = dt/\bar{a} = \bar{\gamma}_w d\tau_w/\bar{a}$.

- This leads to conventional dressed parameters *which do not sum to 1*, e.g.,

$$\Omega_M = \bar{\gamma}_w^3 \bar{\Omega}_M .$$

Tracker solution PRL 99, 251101

- General exact solution possesses a “tracker limit”

$$\bar{a} = \frac{\bar{a}_0 (3\bar{H}_0 t)^{2/3}}{2 + f_{v0}} \left[3f_{v0}\bar{H}_0 t + (1 - f_{v0})(2 + f_{v0}) \right]^{1/3}$$
$$f_v = \frac{3f_{v0}\bar{H}_0 t}{3f_{v0}\bar{H}_0 t + (1 - f_{v0})(2 + f_{v0})},$$

- Void fraction $f_v(t)$ determines many parameters:

$$\bar{\gamma}_w = 1 + \frac{1}{2}f_v = \frac{3}{2}\bar{H}t$$

$$\tau_w = \frac{2}{3}t + \frac{2(1 - f_{v0})(2 + f_{v0})}{27f_{v0}\bar{H}_0} \ln \left(1 + \frac{9f_{v0}\bar{H}_0 t}{2(1 - f_{v0})(2 + f_{v0})} \right)$$

$$\bar{\Omega}_M = \frac{4(1 - f_v)}{(2 + f_v)^2}$$

Apparent cosmic acceleration

- Volume average observer sees no apparent cosmic acceleration

$$\bar{q} = \frac{2(1 - f_v)^2}{(2 + f_v)^2}.$$

As $t \rightarrow \infty$, $f_v \rightarrow 1$ and $\bar{q} \rightarrow 0^+$.

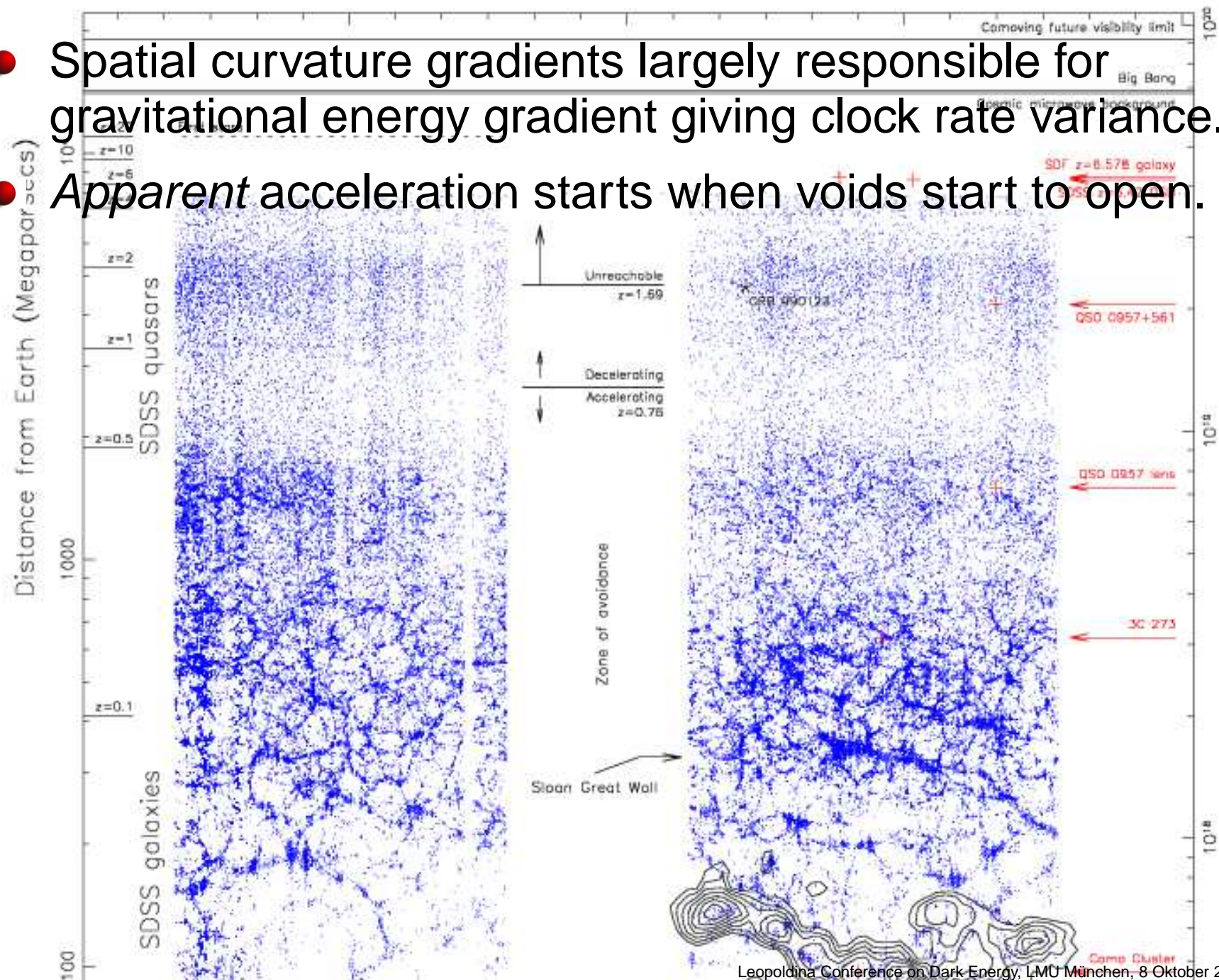
- A wall observer registers apparent cosmic acceleration

$$q = \frac{-(1 - f_v)(8f_v^3 + 39f_v^2 - 12f_v - 8)}{(4 + f_v + 4f_v^2)^2},$$

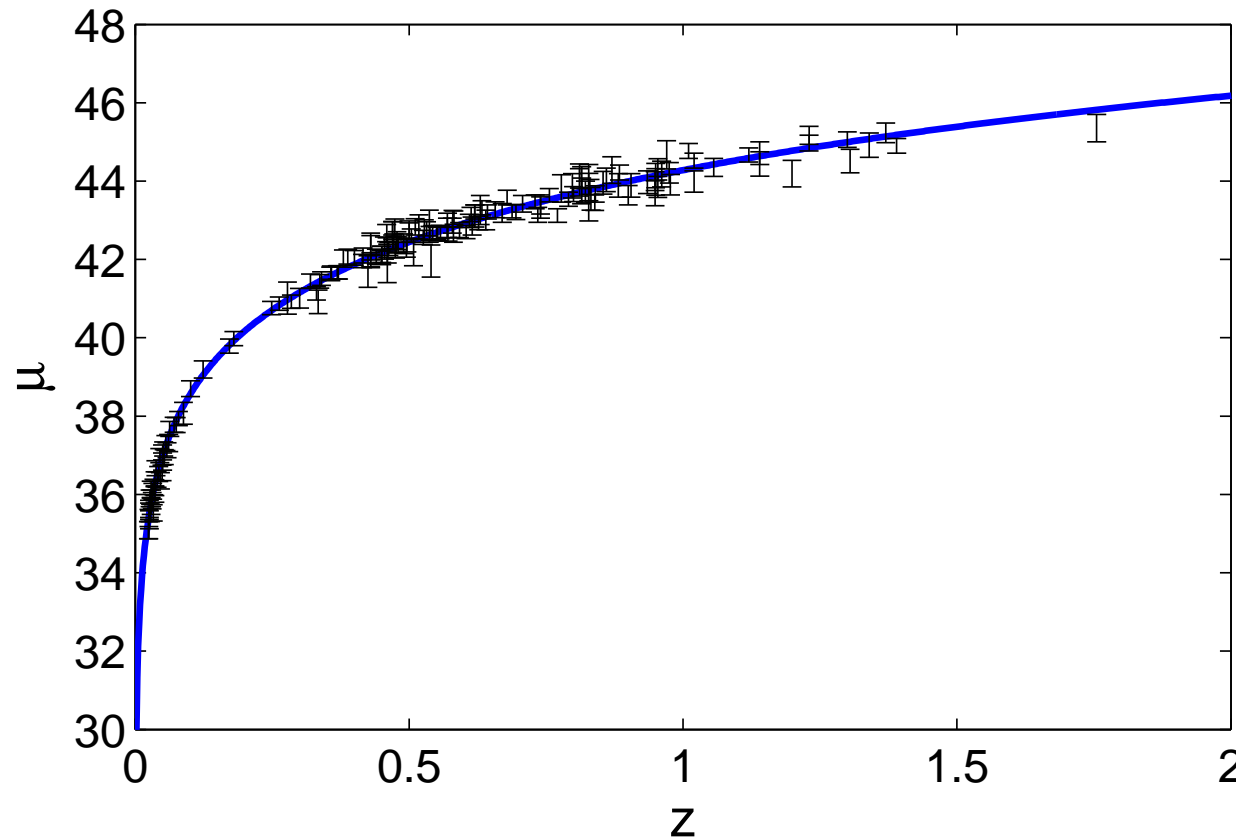
Effective deceleration parameter starts at $q \sim \frac{1}{2}$, for small f_v ; changes sign when $f_v = 0.58670773\dots$, and approaches $q \rightarrow 0^-$ at late times.

Cosmic coincidence problem solved

- Spatial curvature gradients largely responsible for gravitational energy gradient giving clock rate variance.
- *Apparent* acceleration starts when voids start to open.

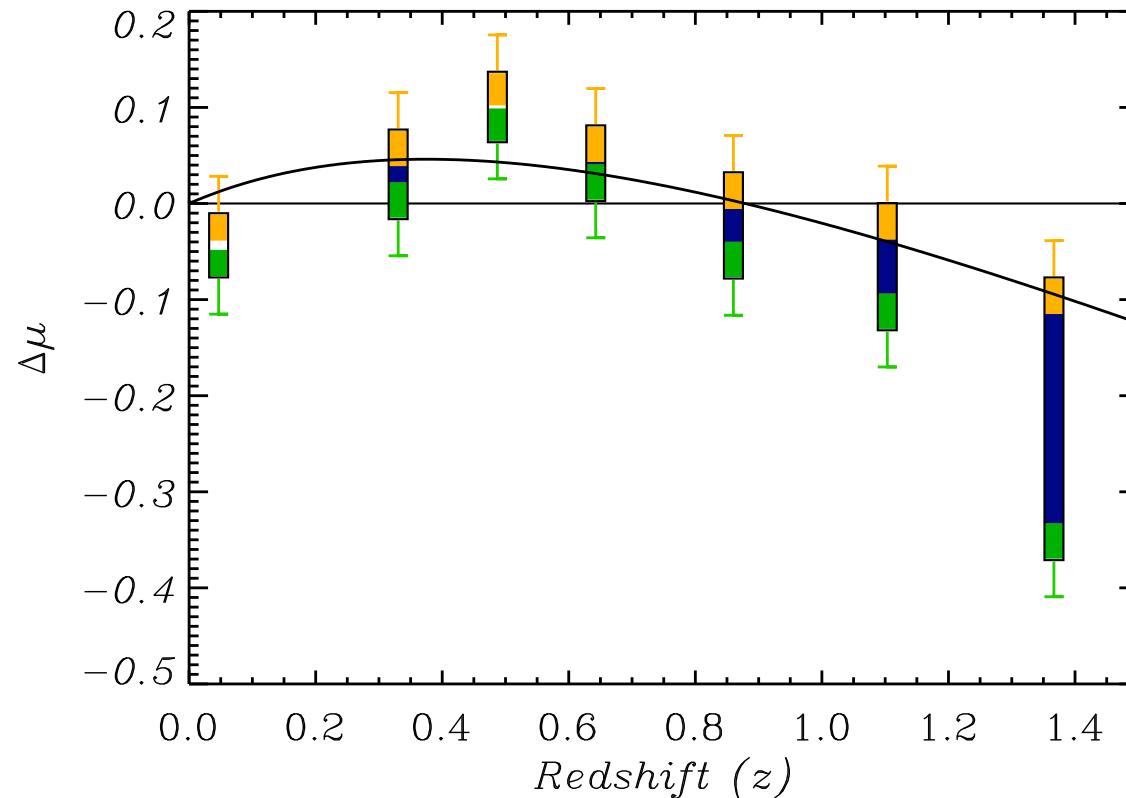


Test 1: Snela luminosity distances



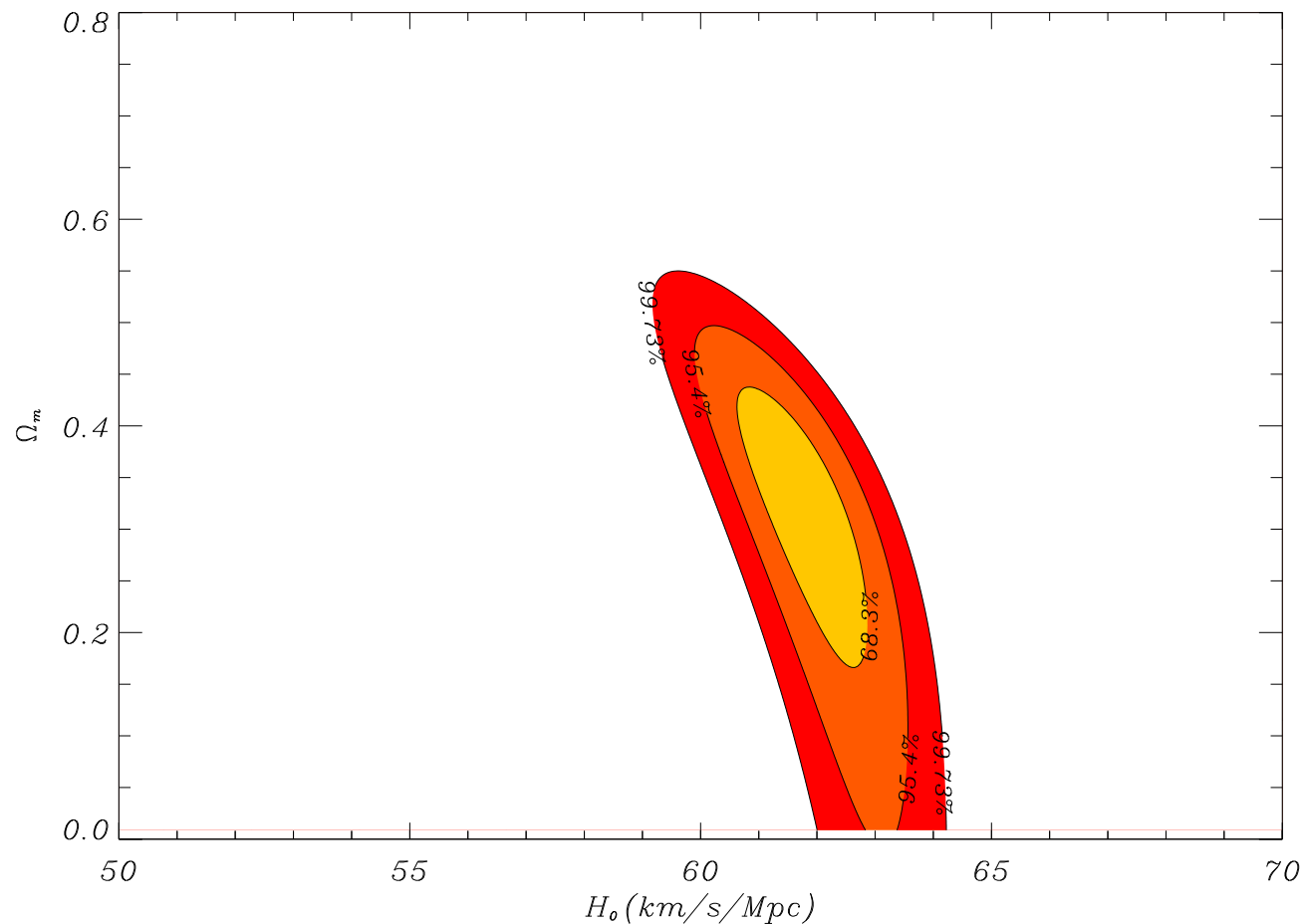
- Type Ia supernovae of Riess07 Gold data set fit with χ^2 per degree of freedom = 0.9
- With $55 \leq H_0 \leq 75 \text{ km sec}^{-1} \text{ Mpc}^{-1}$, $0.01 \leq \Omega_{M0} \leq 0.5$, find Bayes factor $\ln B = 0.27$ in favour or FB model (marginally): statistically indistinguishable from Λ CDM.

Test 1: Snela luminosity distances



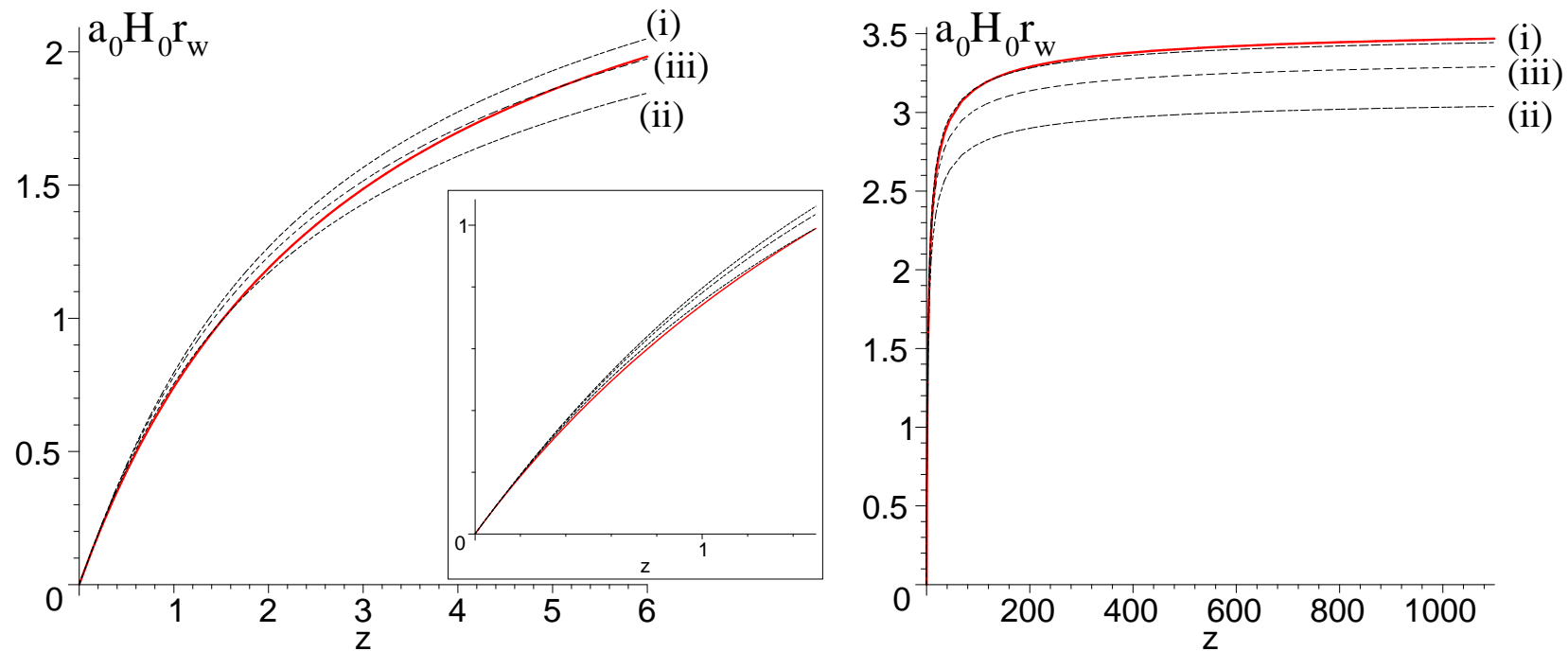
- Plot shows difference of model apparent magnitude and that of an empty Milne universe of same Hubble constant $H_0 = 61.73 \text{ km sec}^{-1} \text{ Mpc}^{-1}$. Note: residual depends on the expansion rate of the Milne universe subtracted (2σ limits on H_0 indicated by whiskers)

Test 1: Snela luminosity distances



Best-fit H_0 agrees with HST key team, Sandage et al.,
 $H_0 = 62.3 \pm 1.3$ (stat) ± 5.0 (syst) km sec⁻¹ Mpc⁻¹ [ApJ 653
(2006) 843].

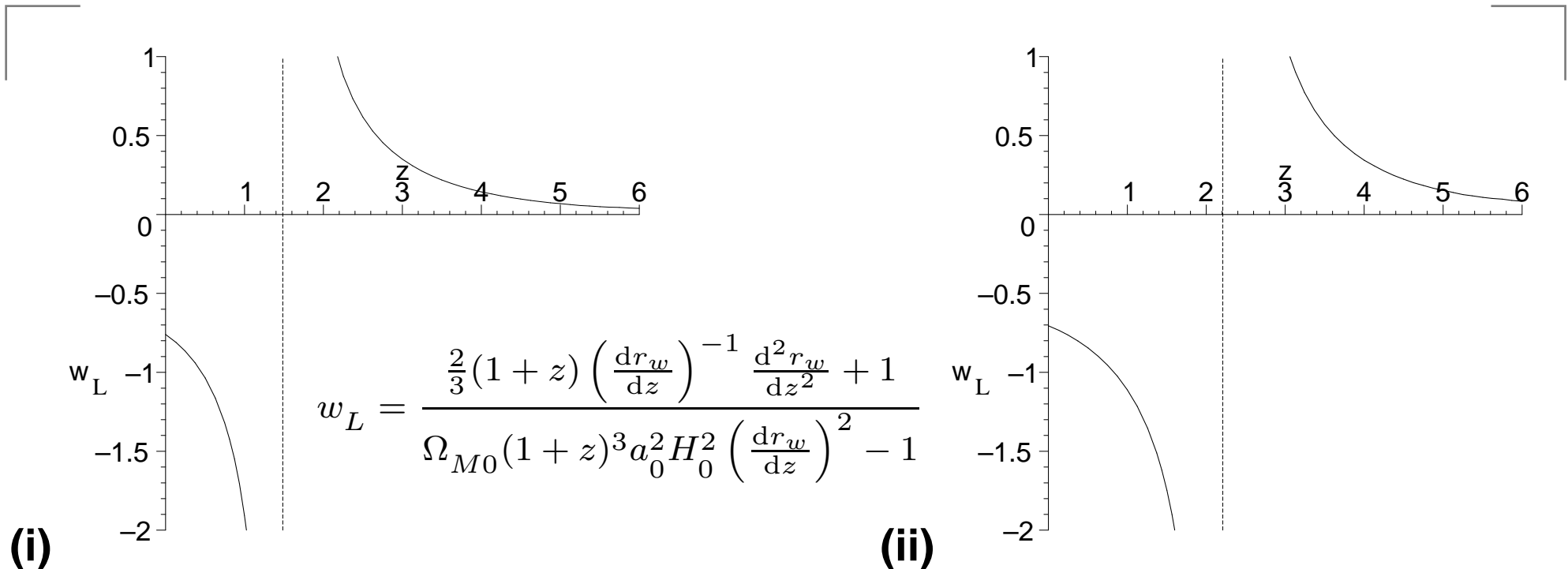
Dressed “comoving distance” $r_w(z)$



Best-fit FB model (**red line**) compared to 3 spatially flat Λ CDM models: **(i)** best-fit to WMAP5 only ($\Omega_\Lambda = 0.751$); **(ii)** best-fit to (Riess07) Snela only ($\Omega_\Lambda = 0.66$); **(iii)** joint WMAP5 + BAO + Snela fit ($\Omega_\Lambda = 0.721$)

- FB model closest to best-fit Λ CDM to *Snela only* result ($\Omega_{M0} = 0.34$) at *low redshift*, and to *WMAP5 only* result ($\Omega_{M0} = 0.249$) at *high redshift*

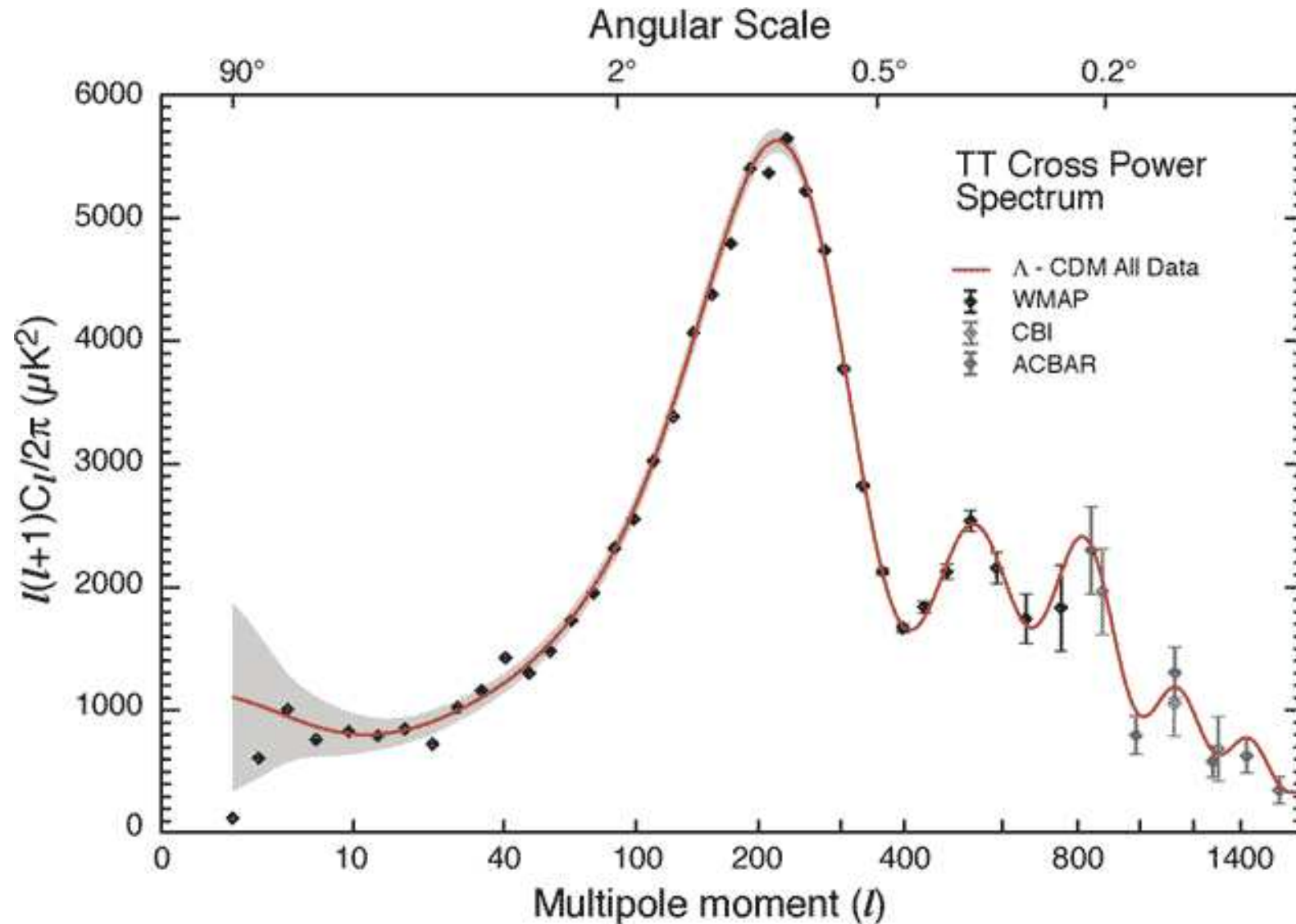
Equivalent “equation of state”?



A formal “dark energy equation of state” $w_L(z)$ for the best-fit FB model, $f_{v0} = 0.76$, calculated directly from $r_w(z)$: (i) $\Omega_{M0} = 0.33$; (ii) $\Omega_{M0} = 0.279$.

- Description by a “dark energy equation of state” makes no sense when there is no physics behind it; but average value $w_L \simeq -1$ for $z < 0.7$ makes empirical sense.

Test 2: Angular scale of CMB Doppler peaks

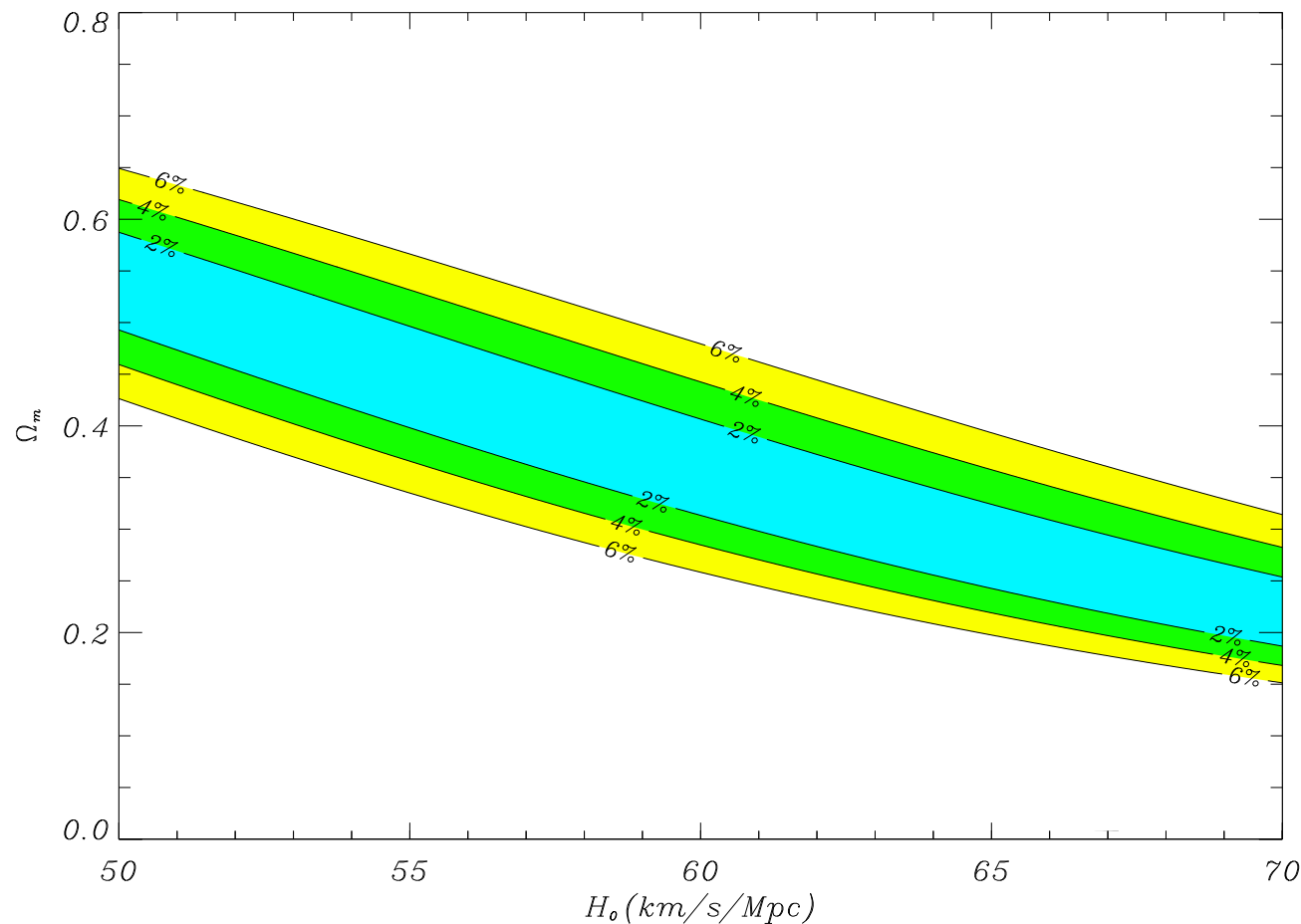


Power in CMB temperature anisotropies versus angular size of fluctuation on sky

Test 2: Angular scale of CMB Doppler peaks

- Angular scale is related to spatial curvature of FLRW models
- Relies on the simplifying assumption that spatial curvature is same everywhere
- In new approach spatial curvature is not the same everywhere
- Volume–average observer measures lower mean CMB temperature ($\bar{T}_0 \sim 1.98$ K, c.f. $T_0 \sim 2.73$ K in walls) and a smaller angular anisotropy scale
- Relative focussing between voids and walls
- Integrated Sachs–Wolfe effect needs recomputation
- Here just calculate angular–diameter distance of sound horizon

Test 2: Angular scale of CMB Doppler peaks

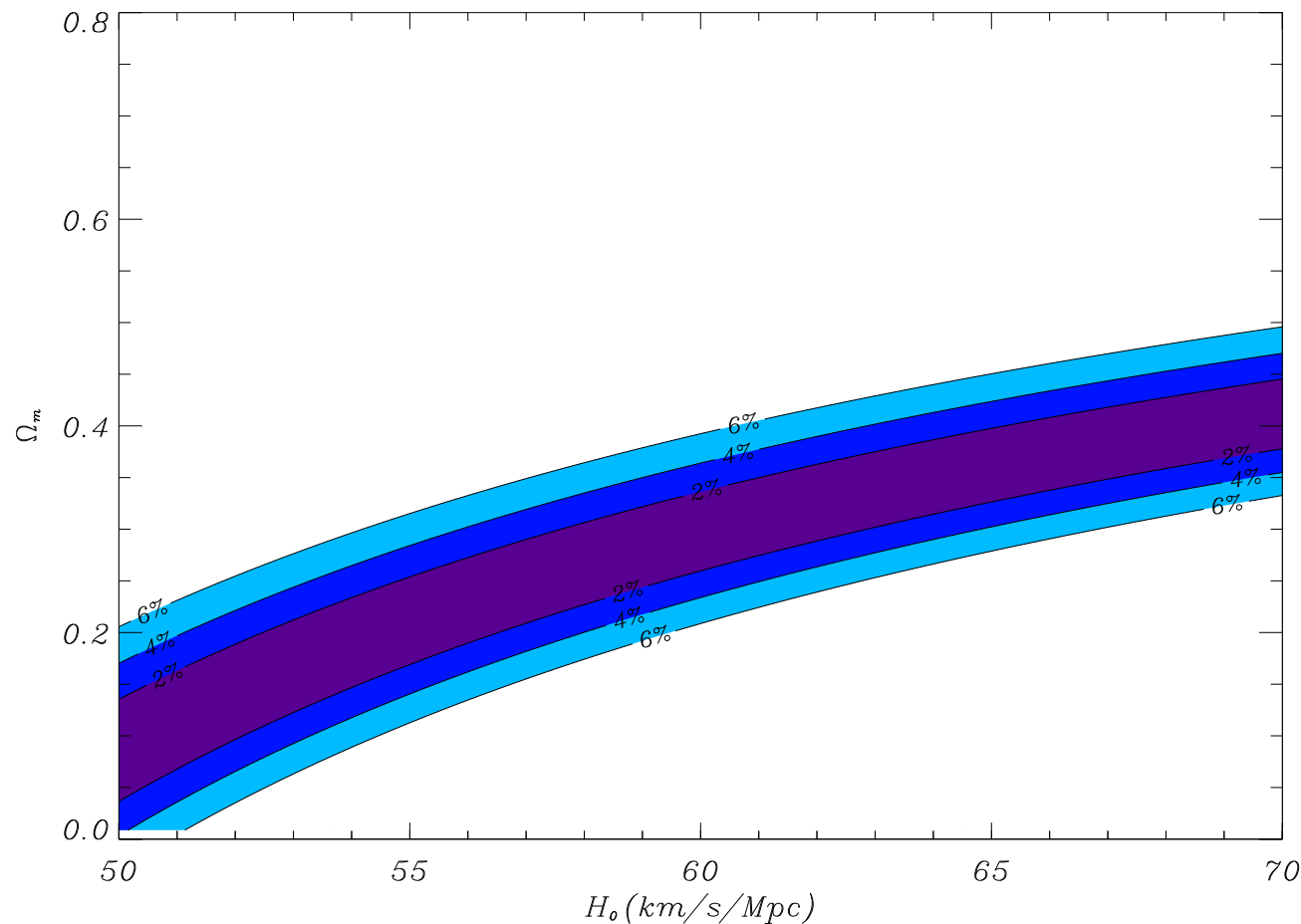


Parameters within the (Ω_m, H_0) plane which fit the angular scale of the sound horizon $\delta = 0.01$ rad deduced for WMAP, to within 2%, 4% and 6%.

Test 3: Baryon acoustic oscillation scale

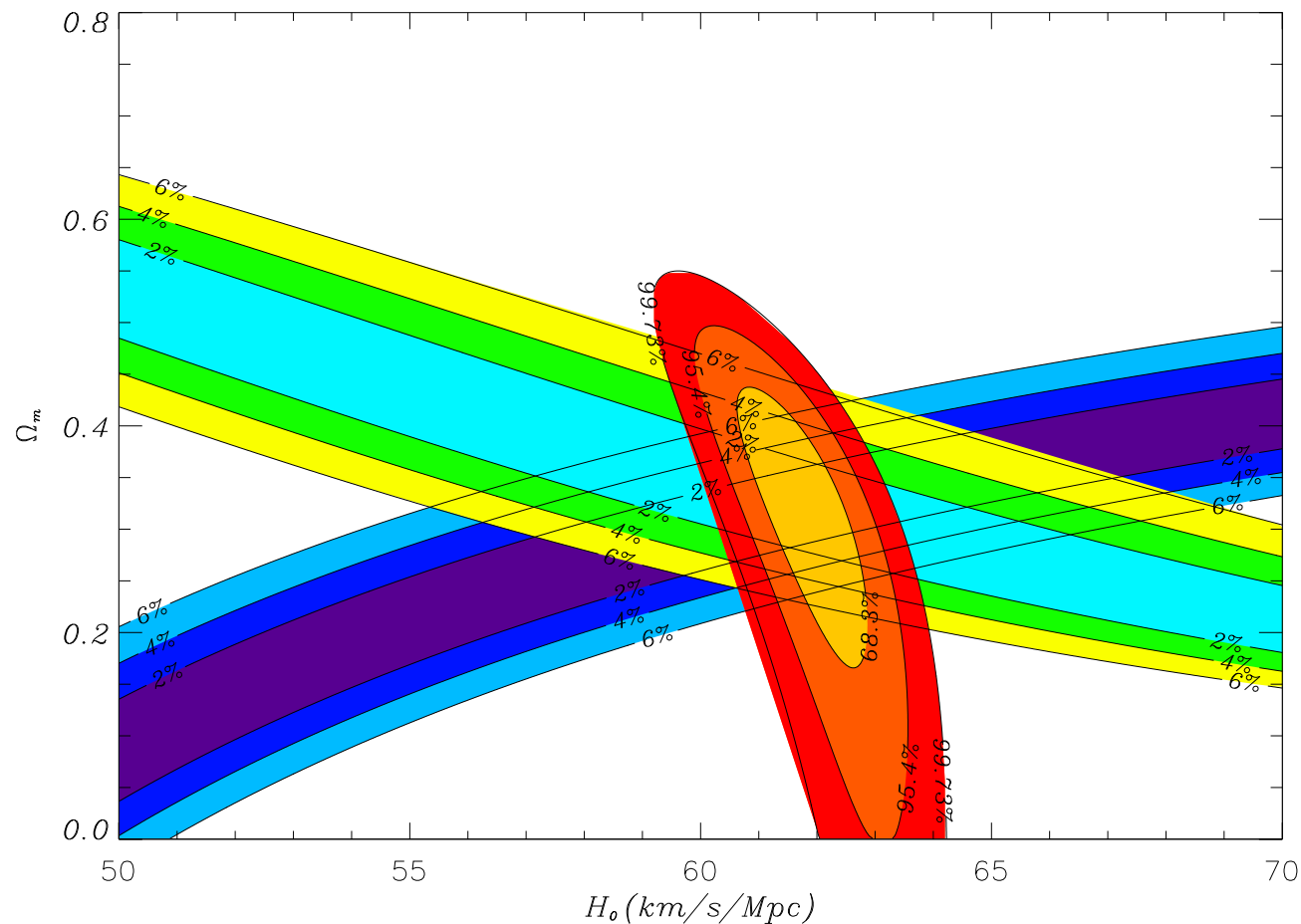
- In 2005 Cole et al. (2dF), and Eisenstein et al. (SDSS) detected the signature of the comoving baryon acoustic oscillation in galaxy clustering statistics
- Powerful independent probe of “dark energy”
- Here the effective dressed geometry should give an equivalent scale

Test 3: Baryon acoustic oscillation scale



Parameters within the (Ω_m, H_0) plane which fit the effective comoving baryon acoustic oscillation scale of $104h^{-1}$ Mpc, as seen in 2dF and SDSS.

Agreement of independent tests



Best-fit parameters: $H_0 = 61.7^{+1.2}_{-1.1} \text{ km sec}^{-1} \text{ Mpc}^{-1}$,
 $\Omega_m = 0.33^{+0.11}_{-0.16}$ (1σ errors for Snela only) [Leith, Ng & Wiltshire, ApJ 672 (2008) L91]

Resolution of Li abundance anomaly?

- Tests 2 & 3 shown earlier use the baryon-to-photon ratio $\eta_{B\gamma} = 4.6\text{--}5.6 \times 10^{-10}$ admitting concordance with lithium abundances favoured prior to WMAP in 2003
- Conventional dressed parameter $\Omega_{M0} = 0.33$ for wall observer means $\bar{\Omega}_{M0} = 0.127$ for the volume-average.
- Conventional theory predicts the *volume-average baryon fraction*. With old BBN favoured $\eta_{B\gamma}$:
 $\bar{\Omega}_{B0} \simeq 0.027\text{--}0.033$; but this translates to a conventional dressed baryon fraction parameter $\Omega_{B0} \simeq 0.072\text{--}0.088$
- The mass ratio of baryonic matter to non-baryonic dark matter is typically increased to 1:3
- Enough baryon drag to fit peak heights ratio

Alleviation of age problem

- Old structures seen at large redshifts are a problem for Λ CDM.
- Problem alleviated here; expansion age is increased, by an increasingly larger relative fraction at larger redshifts, e.g., for best-fit values
 Λ CDM $\tau = 0.85$ Gyr at $z = 6.42$, $\tau = 0.365$ Gyr at $z = 11$
FB $\tau = 1.14$ Gyr at $z = 6.42$, $\tau = 0.563$ Gyr at $z = 11$
- Present age of universe for best-fit is $\tau_0 \simeq 14.7$ Gyr for wall observer; $t_0 \simeq 18.6$ Gyr for volume-average observer.
- Suggests problems of under-emptiness of voids in Newtonian N-body simulations may be an issue of using volume-average time?? The simulations need to be carefully reconsidered.

Variance of Hubble flow

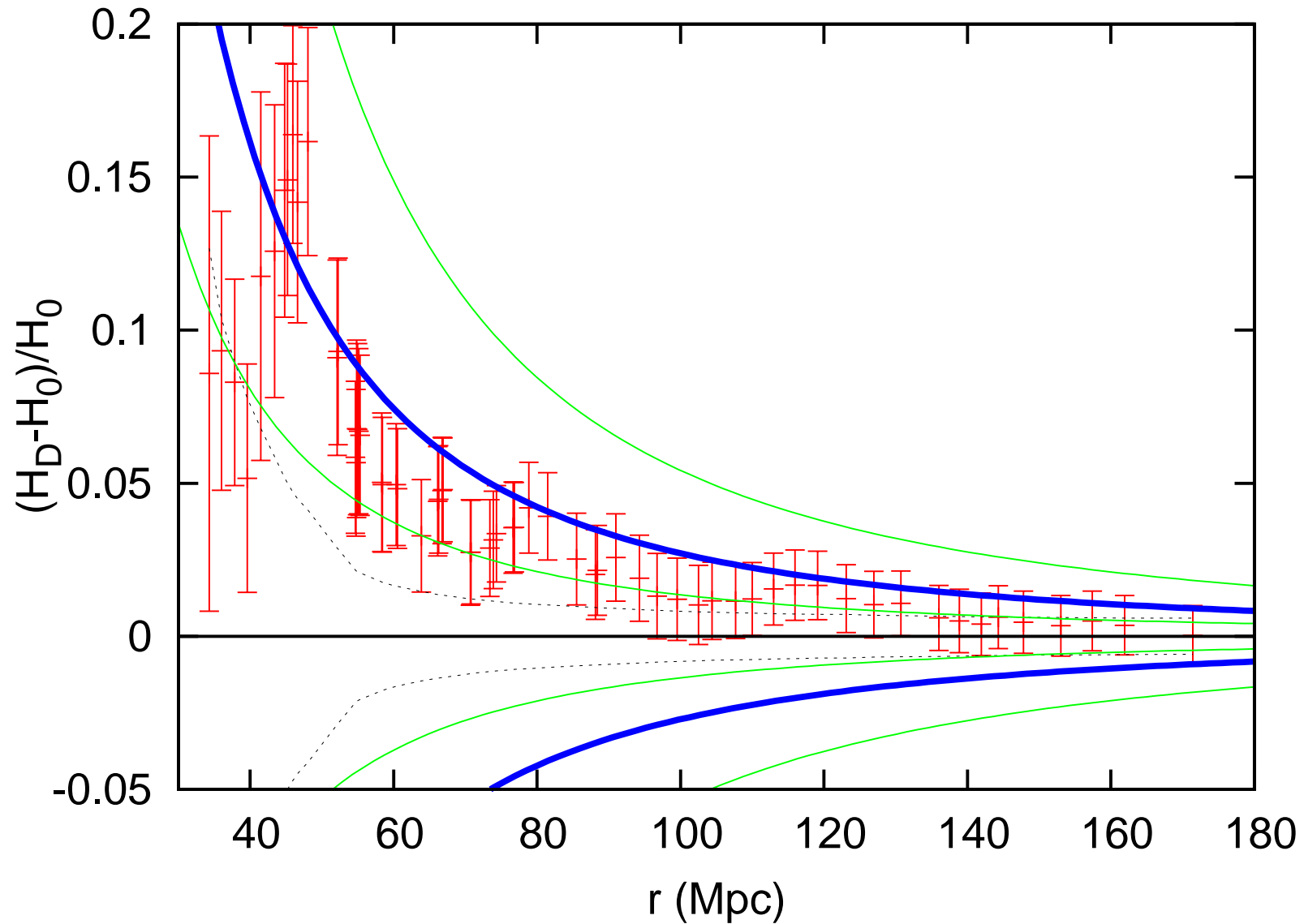
- Relative to “wall clocks” the global average Hubble parameter $H_{\text{av}} > \bar{H}$
- \bar{H} is nonetheless also the locally measurable Hubble parameter within walls
- TESTABLE PREDICTION:

$$H_{\text{av}} = \bar{\gamma}_{\text{w}} \bar{H} - \bar{\gamma}_{\text{w}}^{-1} \bar{\gamma}'_{\text{w}}$$

- With $H_0 = 62 \text{ km sec}^{-1} \text{ Mpc}^{-1}$, expect according to our measurements:
 $\bar{H}_0 = 48 \text{ km sec}^{-1} \text{ Mpc}^{-1}$ within ideal walls (e.g., around Virgo cluster?); and
 $\bar{H}_{\text{v}0} = 76 \text{ km sec}^{-1} \text{ Mpc}^{-1}$ across local voids (scale $\sim 45 \text{ Mpc}$)

Explanation for Hubble bubble

- As voids occupy largest volume of space expect to measure higher average Hubble constant locally until the global average relative volumes of walls and voids are sampled at scale of homogeneity; thus expect maximum H_0 value for isotropic average on scale of dominant void diameter, $30h^{-1}\text{Mpc}$, then decreasing till levelling out by $100h^{-1}\text{Mpc}$.
- Consistent with observed Hubble bubble feature (Jha, Riess, Kirshner ApJ 659, 122 (2007)), which is unexplained (and problem for) ΛCDM .
- Intrinsic variance in apparent Hubble flow exposes a local scale dependence which may partly explain difficulties astronomers have had in converging on a value for H_0 .



N. Li and D. Schwarz, arxiv:0710.5073v1–2

Best fit parameters

- Hubble constant $H_0 = 61.7_{-1.1}^{+1.2} \text{ km sec}^{-1} \text{ Mpc}^{-1}$
- present void volume fraction $f_{v0} = 0.76_{-0.09}^{+0.12}$
- bare density parameter $\bar{\Omega}_{M0} = 0.125_{-0.069}^{+0.060}$
- dressed density parameter $\Omega_{M0} = 0.33_{-0.16}^{+0.11}$
- non-baryonic dark matter / baryonic matter mass ratio
 $(\bar{\Omega}_{M0} - \bar{\Omega}_{B0}) / \bar{\Omega}_{B0} = 3.1_{-2.4}^{+2.5}$
- bare Hubble constant $\bar{H}_0 = 48.2_{-2.4}^{+2.0} \text{ km sec}^{-1} \text{ Mpc}^{-1}$
- mean lapse function $\bar{\gamma}_0 = 1.381_{-0.046}^{+0.061}$
- deceleration parameter $q_0 = -0.0428_{-0.0002}^{+0.0120}$
- wall age universe $\tau_0 = 14.7_{-0.5}^{+0.7} \text{ Gyr}$

Model comparison

	Λ CDM	FB scenario
Sn Ia luminosity distances	Yes	Yes
BAO scale (clustering)	Yes	Yes
Sound horizon scale (CMB)	Yes	Yes
Doppler peak fine structure	Yes	[still to calculate]
Integrated Sachs–Wolfe effect	Yes	[still to calculate]
Primordial ${}^7\text{Li}$ abundances	No	Yes?
CMB ellipticity	No	[Maybe]
CMB low multipole anomalies	No	[?Foreground void: Rees–Sciama dipole]
Hubble bubble	No	Yes
Nucleochronology dates of old globular clusters	Tension	Yes
X-ray cluster abundances	Marginal	Yes
Emptiness of voids	No	[Maybe]
Sandage-de Vaucouleurs paradox	No	Yes
Coincidence problem	No	Yes

Conclusion

- Apparent cosmic acceleration can be understood purely within general relativity; by (i) treating geometry of universe more realistically; (ii) understanding fundamental aspects of general relativity which have not been fully explored – *quasi-local gravitational energy*, of *gradients* in spatial curvature etc.
- The “fractal bubble” model passes three major independent tests which support Λ CDM and may resolve significant puzzles and anomalies.
- Every cosmological parameter requires subtle recalibration, but no “new” physics beyond dark matter: no Λ , no exotic scalars, no modifications to gravity.
- Questions raised – otherwise unanswered – should be addressed irrespective of phenomenological success.